



EUROfusion

WPJET1-CPR(18) 19498

E Joffrin et al.

Overview of the JET preparation for Deuterium-Tritium Operation

Preprint of Paper to be submitted for publication in Proceeding of
27th IAEA Fusion Energy Conference



This work has been carried out within the framework of the EUROfusion Consortium and has received funding from the Euratom research and training programme 2014-2018 under grant agreement No 633053. The views and opinions expressed herein do not necessarily reflect those of the European Commission.

This document is intended for publication in the open literature. It is made available on the clear understanding that it may not be further circulated and extracts or references may not be published prior to publication of the original when applicable, or without the consent of the Publications Officer, EUROfusion Programme Management Unit, Culham Science Centre, Abingdon, Oxon, OX14 3DB, UK or e-mail Publications.Officer@euro-fusion.org

Enquiries about Copyright and reproduction should be addressed to the Publications Officer, EUROfusion Programme Management Unit, Culham Science Centre, Abingdon, Oxon, OX14 3DB, UK or e-mail Publications.Officer@euro-fusion.org

The contents of this preprint and all other EUROfusion Preprints, Reports and Conference Papers are available to view online free at <http://www.euro-fusionscipub.org>. This site has full search facilities and e-mail alert options. In the JET specific papers the diagrams contained within the PDFs on this site are hyperlinked

Overview of the JET preparation for Deuterium-Tritium Operation with the ITER Like-Wall

EUROfusion Consortium JET, Culham Science Centre, Abingdon, OX14 3DB, UK

e-mail contact of first author: Emmanuel.Joffrin@jet.euro-fusion.org

E. Joffrin⁷, S. Abdualley³⁴, M. Abhangi⁴⁶, P. Abreu⁵², V. Afanasev⁵³, M. Afzal⁶, K.M. Aggarwal⁷⁸, T. Ahlgren⁹⁹, L. Aho-Mantila¹⁰⁷, N. Aiba⁶⁶, M. Airila¹⁰⁷, T. Alarcon⁷, R. Albanese¹⁰, D. Alegre⁵⁶, S. Aleiferis⁶⁸, E. Alessi⁴², P. Aleynikov⁵⁴, A. Alkseev⁶⁹, M. Allinson⁶, B. Alper⁶, E. Alves⁵², G. Ambrosino¹⁰, R. Ambrosino¹⁰, V. Amosov⁸⁸, E. Andersson Sundén¹⁸, R. Andrews⁶, M. Angelone⁸⁹, M. Anghel⁸⁵, C. Angioni⁶⁰, L. Appel⁶, C. Appelbee⁶, P. Arena²⁶, M. Ariola¹⁰, S. Arshad³⁷, J. Artaud⁷, W. Arter⁶, A. Ash⁶, N. Ashikawa⁶⁵, V. Aslanyan⁶², O. Asunta¹, O. Asztalos¹⁰⁸, F. Auriemma¹¹, Y. Austin⁶, L. Avotina¹⁰², M. Axton⁶, C. Ayres⁶, A. Baciero⁵⁶, D. Baião⁵², I. Balboa⁶, M. Balden⁶⁰, N. Balshaw⁶, V.K. Bandaru⁶⁰, J. Banks⁶, Y.F. Baranov⁶, C. Barcellona²⁶, T. Barnard⁶, M. Barnes⁷⁷, R. Barnsley⁵⁴, A. Baron Wiechec⁶, L. Barrera Orte²⁹, M. Baruzzo¹¹, V. Basiuk⁷, M. Bassan⁵⁴, R. Bastow⁷⁵, A. Batista⁵², P. Batistoni⁸⁹, L. Baumane¹⁰², B. Bauvir⁵⁴, L. Baylor⁷⁰, P.S. Beaumont⁶, M. Beckers³⁴, B. Beckett⁶, N. Bekris³⁰, M. Beldishevski⁶, K. Bell⁶, F. Belli⁸⁹, É. Belonohy^{6,60}, J. Benayas⁶, H. Bergsaker³⁸, J. Bernardo⁵², M. Bernert⁶⁰, M. Berry⁷⁵, L. Bertalot⁵⁴, C. Besiliu⁶, H. Betar⁴³, M. Beurskens⁶¹, J. Bielecki⁴⁷, T. Biewer⁷⁰, R. Bilato⁶⁰, O. Biletskyi⁴⁵, P. Bilková⁵⁰, F. Binda¹⁸, G. Birkenmeier⁶⁰, J.P.S. Bizarro⁵², C. Björkas⁹⁹, J. Blackburn⁷⁵, T.R. Blackman⁶, P. Blanchard²⁸, P. Blatchford⁶, V. Bobkov⁶⁰, A. Boboc⁶, O. Bogar¹⁶, P. Bohm⁵⁰, T. Bohm¹⁰⁶, I. Bolshakova⁵⁹, T. Bolzonella¹¹, N. Bonanomi^{60,97}, L. Boncagni⁸⁹, D. Bonfiglio¹¹, X. Bonnin⁵⁴, J. Boom⁶⁰, D. Borba^{30,52}, D. Borodin³⁴, I. Borodkina³⁴, C. Boulbe⁹⁵, C. Bourdelle⁷, M. Bowden⁶, C. Bowman¹⁰⁹, T. Boyce⁶, H. Boyer⁶, S.C. Bradnam⁶, V. Braic⁸⁷, R. Bravanec³⁶, B. Breizman¹⁰³, D. Brennan⁶, S. Breton⁷, A. Brett⁶, S. Brezinsek³⁴, M. Bright⁶, M. Brix⁶, W. Broeckx⁷⁹, M. Brombin¹¹, A. Broślawski⁶³, B. Brown⁶, D. Brunetti⁴², E. Bruno⁵⁴, J. Buch⁴⁶, J. Buchanan⁶, R. Buckingham⁶, M. Buckley⁶, M. Bucolo²⁶, R. Budny⁷³, H. Bufferand⁷, S. Buller²⁰, P. Bunting⁶, P. Buratti⁸⁹, A. Burckhart⁶⁰, G. Burroughes⁷⁵, A. Buscarino²⁶, A. Busse⁶, D. Butcher⁶, B. Butler⁶, I. Bykov³⁸, P. Cahyna⁵⁰, G. Calabrò¹⁰⁵, L. Calacci⁹³, D. Callaghan⁶, J. Callaghan⁶, I. Calvo⁵⁶, Y. Camenen³, P. Camp⁶, D.C. Campling⁶, B. Cannas¹⁵, A. Capat⁶, S. Carcangiu¹⁵, P. Card⁶, A. Cardinali⁸⁹, P. Carman⁶, D. Carnevale⁹³, M. Carr⁶, D. Carralero^{56,60}, L. Carraro¹¹, B.B. Carvalho⁵², I. Carvalho⁵², P. Carvalho⁵², D.D. Carvalho⁵², F.J. Casson⁶, C. Castaldo⁸⁹, N. Catarino⁵², F. Causa⁴², R. Cavazzana¹¹, K. Cave-Ayland⁶, M. Cavedon⁶⁰, M. Cecconello¹⁸, S. Ceccuzzi⁸⁹, E. Cecil⁷³, C.D. Challis⁶, D. Chandra⁴⁶, C.S. Chang⁷³, A. Chankin⁶⁰, I.T. Chapman⁶, B. Chapman²⁴, S.C. Chapman²⁴, M. Chernyshova⁴⁹, A. Chiariello¹⁰, G. Chitarin¹¹, P. Chmielewski⁴⁹, L. Chone¹, G. Ciraolo⁷, D. Ciric⁶, J. Citrin³³, F. Clairet⁷, M. Clark⁶, E. Clark⁷⁵, R. Clarkson⁶, R. Clay⁶, C. Clements⁶, J.P. Coad⁶, P. Coates⁶, A. Cobalt⁶, V. Coccoresse¹⁰, V. Cocilovo⁸⁹, R. Coelho⁵², J.W. Coenen³⁴, I. Coffey⁷⁸, L. Colas⁷, B. Colling⁶, S. Collins⁷⁵, D. Conka¹⁰², S. Conroy¹⁸, N. Conway⁶, D. Coombs⁶, S.R. Cooper⁶, C. Corradino²⁶, Y. Corre⁷, G. Corrigan⁶, D. Coster⁶⁰, T. Craciunescu⁸⁶, S. Cramp⁶, C. Crapper⁶, F. Crisanti⁸⁹, G. Croci⁹⁷, D. Croft⁶, K. Crombé¹⁴, N. Cruz⁵², G. Cseh¹⁰⁸, A. Cufar⁸⁰, A. Cullen⁶, P. Curson⁶, M. Curuia⁸⁵, A. Czarnecka⁴⁹, T. Czarski⁴⁹, I. Cziegler¹⁰⁹, H. Dabirikhah⁶, A. Dal Molin⁹⁷, P. Dalglish⁶, S. Dalley⁶, J. Dankowski⁴⁷, D. Darrow⁷³, P. David⁶⁰, A. Davies⁶, W. Davis^{54,73}, K. Dawson⁶, I. Day⁶, C. Day⁵⁵, M. De Bock⁵⁴, A. de Castro⁵⁶, G. De Dominici⁷, E. de la Cal⁵⁶, E. de la Luna⁵⁶, G. De Masi¹¹, G. De Temmerman⁵⁴, G. De Tommasi¹⁰, P. de Vries⁵⁴, J. Deane⁶, R. Dejarnac⁵⁰, D. Del Sarto⁴³, E. Delabie⁷⁰, V. Demerdzhiev⁶, A. Dempsey²⁷, N. den Harder³³, R.O. Dendy⁶, J. Denis⁷, P. Denner³⁴, S. Devaux⁴³, P. Devynck⁷, F. Di Maio⁵⁴, A. Di Siena⁶⁰, C. Di Troia⁸⁹, D.

Dickinson¹⁰⁹, P. Dinca⁸⁶, T. Dittmar³⁴, J. Dobrashian⁶, H. Doerk⁶⁰, R.P. Doerner⁸, F. Domptail⁶, T. Donné²⁹, S.E. Dorling⁶, D. Douai⁷, S. Dowson⁶, A. Drenik^{60,80}, M. Dreval⁴⁵, P. Drewelow⁶¹, P. Drews³⁴, Ph. Duckworth⁵⁴, R. Dumont⁷, P. Dumortier⁵⁷, D. Dunai¹⁰⁸, M. Dunne⁶⁰, I. Đuran⁵⁰, F. Durodié⁵⁷, P. Dutta⁴⁶, B.P. Duval²⁸, R. Dux⁶⁰, K. Dylst⁷⁹, P.V. Edappala⁴⁶, A.M. Edwards⁶, J.S. Edwards⁶, Th. Eich⁶⁰, N. Eidietis³⁹, A. Eksaeva³⁴, R. Ellis⁶, G. Ellwood⁶, C. Elsmore⁶, S. Emery⁶, M. Enachescu⁸⁴, G. Ericsson¹⁸, J. Eriksson¹⁸, F. Eriksson²⁵, L.G. Eriksson³¹, S. Ertmer³⁴, S. Esquembri⁹², A.L. Esquisabel⁷⁰, H.G. Esser³⁴, G. Ewart⁶, E. Fable⁶⁰, D. Fagan⁷⁵, M. Faitsch⁶⁰, D. Falie⁸⁶, A. Fanni¹⁵, A. Farahani⁶, A. Fasoli²⁸, B. Faugeras⁹⁵, S. Fazinić⁷⁶, F. Felici^{28,33}, R.C. Felton⁶, S. Feng⁹⁷, A. Fernades⁵², H. Fernandes⁵², J. Ferreira⁵², D.R. Ferreira⁵², G. Ferro⁹³, J.A. Fessey⁶, O. Ficker⁵⁰, A. Field⁶, S. Fietz⁶⁰, L. Figini⁴², J. Figueiredo³⁰, A. Figueiredo⁵², J. Figueiredo⁵², N. Fil⁶², P. Finburg⁶, U. Fischer⁵⁵, L. Fittill⁶, M. Fitzgerald⁶, D. Flammini⁸⁹, J. Flanagan⁶, K. Flinders⁶, S. Foley⁶, N. Fonnese⁹³, J.M. Fontdecaba⁵⁶, A. Formisano¹⁰, L. Forsythe⁷⁵, L. Fortuna²⁶, E. Fransson²⁵, M. Frasca²⁶, L. Frassinetti³⁸, M. Freisinger³⁴, R. Fresa¹⁰, R. Fridström³⁸, D. Frigione⁸⁹, V. Fuchs⁵⁰, V. Fusco⁸⁹, S. Futatani⁹⁴, K. Gál^{29,60}, D. Galassi², K. Gałazka⁴⁹, S. Galeani⁹³, D. Gallart⁵, R. Galvão⁹, Y. Gao³⁴, J. Garcia⁷, A. Garcia-Carrasco³⁸, M. García-Muñoz⁹¹, M. Gardener⁶, L. Garzotti⁶, J. Gaspar³, P. Gaudio⁹³, D. Gear⁶, T. Gebhart⁷⁰, S. Gee⁶, B. Geiger⁶⁰, M. Gelfusa⁹³, R. George⁶, S. Gerasimov⁶, G. Gervasini⁴², M. Gethins⁶, Z. Ghani⁶, M. Ghaté⁴⁶, M. Gherendi⁸⁶, F. Ghezzi⁴², J.C. Giacalone⁷, L. Giacomelli⁴², G. Giacometti³, K. Gibson¹⁰⁹, T. Giegerich⁵⁵, L. Gil⁵², M.R. Gilbert⁶, D. Gin⁵³, E. Giovannozzi⁸⁹, C. Giroud⁶, S. Glöggler⁶⁰, J. Goff⁶, P. Gohil³⁹, V. Goloborod'ko¹⁰⁰, V. Goloborodko^{45,100}, R. Gomes⁵², B. Gonçalves⁵², M. Goniche⁷, A. Goodyear⁶, G. Gorini⁹⁷, T. Görler⁶⁰, R. Goulding⁷³, A. Goussarov⁷⁹, B. Graham⁶, J.P. Graves²⁸, H. Greuner⁶⁰, B. Grierson⁷³, J. Griffiths⁶, S. Griph⁶, D. Grist⁶, M. Groth¹, R. Grove⁷⁰, M. Gruca⁴⁹, D. Guard⁶, C. Guérard²⁹, C. Guillemaut⁵², R. Guirlet⁷, S. Gulati⁶, C. Gurl⁶, A. Gutierrez-Milla⁵, H. H Utoh⁶⁶, L. Hackett⁶, S. Hacquin^{7,30}, R. Hager⁷³, A. Hakola¹⁰⁷, M. Halitovs¹⁰², S. Hall⁶, S. Hallworth-Cook⁶, C. Ham⁶, M. Hamed⁷, N. Hamilton⁷⁵, C. Hamlyn-Harris⁶, K. Hammond⁶, G. Hancu⁷⁵, J. Harrison⁶, D. Harting⁶, F. Hasenbeck³⁴, Y. Hatano¹⁰⁴, D.R. Hatch¹⁰³, T. Haupt⁶, J. Hawes⁶, N.C. Hawkes⁶, J. Hawkins⁶, P. Hawkins⁷⁵, S. Hazel⁷⁵, P. Heesterman⁶, K. Heinola⁹⁹, C. Hellesen¹⁸, T. Hellsten³⁸, W. Helou⁷, O. Hemming⁶, T.C. Hender⁶, S.S. Henderson⁶, S.S. Henderson¹⁷, M. Henderson⁵⁴, R. Henriques⁵², D. Hepple⁶, J. Herfindal⁷⁰, G. Hermon⁷⁵, C. Hidalgo⁵⁶, W. Higginson⁶, E.G. Highcock²⁰, J. Hillesheim⁶, D. Hillis⁷⁰, K. Hizanidis⁶⁷, A. Hjalmarsson¹⁸, A. Ho³³, J. Hobirk⁶⁰, C.H.A. Hogben⁶, G.M.D. Hogewij³³, A. Hollingsworth⁶, S. Hollis⁶, M. Hölzl⁶⁰, J.-J. Honore⁷⁵, M. Hook⁶, D. Hopley⁶, J. Horáček⁵⁰, G. Hornung¹⁴, A. Horton⁶, L.D. Horton³¹, L. Horvath¹⁰⁹, S.P. Hotchin⁶, R. Howell⁷⁵, A. Hubbard⁶², A. Huber³⁴, V. Huber³⁴, T.M. Huddleston⁶, M. Hughes⁶, J. Hughes⁶², G.T.A. Huijsmans⁵⁴, P. Huynh⁷, A. Hynes⁶, i. Igaune¹⁰², D. Iglesias⁶, N. Imazawa⁶⁶, M. Imrišek⁵⁰, M. Incelli⁸⁹, P. Innocente¹¹, I. Ivanova-Stanik⁴⁹, E. Ivings⁶, S. Jachmich^{30,57}, A. Jackson⁶, T. Jackson⁶, P. Jacquet⁶, J. Jansons¹⁰², F. Jaulmes³³, S. Jednoróg⁴⁹, I. Jenkins⁶, I. Jepu⁸⁶, T. Johnson³⁸, R. Johnson⁶, J. Johnston⁶, L. Joita⁶, J. Joly⁷, E. Jonasson⁷⁵, T. Jones⁶, C. Jones⁶, L. Jones⁶, G. Jones⁶, N. Jones⁷⁵, M. Juvonen⁶, K. K Hoshino⁶⁶, A. Kallenbach⁶⁰, M. Kalsey⁶, T. Kaltiaisenaho¹⁰⁷, K. Kamiya⁶⁶, J. Kaniewski⁶, A. Kantor⁷⁵, A. Kappatou⁶⁰, J. Karhunen¹, D. Karkinsky⁶, M. Kaufman⁷⁰, G. Kaveney⁶, Y. Kazakov⁵⁷, V. Kazantzidis⁶⁷, D.L. Keeling⁶, F.P. Keenan⁷⁸, M. Kempenaars⁶, O. Kent⁶, J. Kent⁷⁵, K. Keogh⁷⁵, E. Khilkevich⁵³, H.-T. Kim⁶, H.T. Kim³⁰, R. King⁶, D. King⁶, D.J. Kinna⁶, V. Kiptily⁶, A. Kirk⁶, K. Kirov⁶, A. Kirschner³⁴, G. Kizane¹⁰², M. Klas¹⁶, C. Klepper⁷⁰, A. Klix⁵⁵, M. Knight⁶, P. Knight⁶, S. Knipe⁶, S. Knott⁹⁶, T. Kobuchi⁶⁶, F. Köchl⁸³, G. Kocsis¹⁰⁸, I. Kodeli⁸⁰, F. Koechl⁶, D. Kogut⁷, S. Koivuranta¹⁰⁷, Y. Kolesnichenko⁴⁵, Z. Kollo⁶, Y. Kominis⁶⁷, M. Köppen³⁴, S. Korolczuk⁶³, B. Kos⁸⁰, H.R. Koslowski³⁴, M. Kotschenreuther¹⁰³, M. Koubiti³, R. Kovaldins¹⁰², O. Kovanda⁵⁰, E. Kowalska-Strzęciwilk⁴⁹, A. Krasilnikov⁸⁸, V. Krasilnikov⁸⁸, N. Krawczyk⁴⁹, M. Kresina⁷, K. Krieger⁶⁰, A. Krivska⁵⁷, U. Kruezi⁵⁴, I. Książek⁴⁸, A. Kukushkin⁶⁹, A. Kundu⁴⁶, T. Kurki-Suonio¹, S. Kwak⁶¹, O.J.

Kwon¹², L. Laguardia⁴², A. Lahtinen⁹⁹, A. Laing⁶, P. Lalousis³⁵, N. Lam⁶, C. Lamb⁷⁵, H.T. Lambert³⁴, P.T. Lang⁶⁰, S. Lanthaler²⁸, E. Lascas Neto²⁸, E. Łaszyńska⁴⁹, R. Lawless⁶, K.D. Lawson⁶, A. Lazaros⁶⁷, E. Lazzaro⁴², R. Leach⁷⁵, G. Learoyd⁶, S. Leerink¹, X. Lefebvre⁶, H.J. Leggate²⁷, J. Lehmann⁶, M. Lehnen⁵⁴, P. Leichauer⁶, D. Leichtle³⁷, F. Leipold⁵⁴, I. Lengar⁸⁰, M. Lennholm³¹, B. Lepiavko⁴⁵, J. Leppänen¹⁰⁷, E. Lerche⁵⁷, A. Lescinskis¹⁰², B. Lescinskis¹⁰², S. Lesnoj⁶, M. Leyland¹⁰⁹, W. Leysen⁷⁹, Y. Li³⁴, L. Li³⁴, Y. Liang³⁴, J. Likonen¹⁰⁷, J. Linke³⁴, Ch. Linsmeier³⁴, B. Lipschultz¹⁰⁹, X. Litaudon^{7,30}, G. Liu⁵⁴, B. Lloyd⁶, V.P. Lo Schiavo¹⁰, T. Loarer⁷, A. Loarte⁵⁴, B. Lomanowski¹, P.J. Lomas⁶, J. Lönnroth^{1,30}, J.M. López⁹², R. Lorenzini¹¹, U. Losada⁵⁶, M. Loughlin⁵⁴, C. Lowry³¹, T. Luce⁵⁴, R. Lucock⁶, A. Lukin⁷¹, C. Luna⁴, M. Lungaroni⁹³, C.P. Lungu⁸⁶, M. Lungu⁸⁶, A. Lunniss¹⁰⁹, T. Lunt⁶⁰, I. Lupelli⁶, V. Lutsenko⁴⁵, A. Lysoivan⁵⁷, P. Macheta⁶, E. Macusova⁵⁰, B. Magesh⁴⁶, C. Maggi⁶, R. Maggiora⁷², S. Mahesan⁶, H. Maier⁶⁰, J. Mailloux⁶, R. Maingi⁷³, R. Makwana⁴⁶, A. Malaquias⁵², K. Malinowski⁴⁹, A. Malizia⁹³, P. Manas^{3,60}, G. Manduchi¹¹, M.E. Manso⁵², P. Mantica⁴², M. Mantsinen⁴¹, A. Manzanares⁹⁰, Ph. Maquet⁵⁴, Y. Marandet³, N. Marcenko⁸⁸, C. Marchetto⁴², O. Marchuk³⁴, N. Marconato¹¹, A. Mariani⁴², M. Marin³³, M. Marinelli⁹³, M. Marinucci⁸⁹, T. Marković⁵⁰, D. Marocco⁸⁹, L. Marot²², J. Marsh⁶, A. Martin⁶, A. Martín de Aguilera⁵⁶, J.R. Martín-Solís¹³, R. Martone¹⁰, Y. Martynova³⁴, S. Maruyama⁵⁴, M. Maslov⁶, S. Matejcik¹⁶, M. Mattei¹⁰, G.F. Matthews⁶, D. Matveev³⁴, E. Matveeva⁵⁰, A. Mauriya⁵², F. Maviglia¹⁰, T. May-Smith⁶, M. Mayer⁶⁰, M.L. Mayoral²⁹, D. Mazon⁷, C. Mazzotta⁸⁹, R. McAdams⁶, P.J. McCarthy⁹⁶, K.G. McClements⁶, O. McCormack¹¹, P.A. McCullen⁶, D. McDonald²⁹, M. McHardy⁶, R. McKean⁶, J. McKehon⁶, L. McNamee⁶, C. Meadowcroft⁶, A. Meakins⁶, S. Medley⁶, S. Meigh⁶, A.G. Meigs⁶, G. Meisl⁶⁰, S. Meiter⁷⁰, S. Meitner⁷⁰, L. Meneses⁵², S. Menmuir⁶, K. Mergia⁶⁸, A. Merle²⁸, P. Merriman⁶, Ph. Mertens³⁴, S. Meshchaninov⁸⁸, A. Messiaen⁵⁷, H. Meyer⁶, R. Michling⁵⁴, D. Milanese⁷², F. Militello⁶, E. Militello-Asp⁶, A. Milocco⁹⁷, G. Miloshevsky⁷⁴, F. Mink⁶⁰, S. Minucci¹⁰⁵, I. Miron⁸⁶, S. Mistry⁶, Y. Miyoshi⁶⁶, J. Mlynář⁵⁰, V. Moiseenko⁴⁵, P. Monaghan⁶, I. Monakhov⁶, S. Moon³⁸, R. Mooney⁶, S. Moradi³², J. Morales⁷, J. Moran⁶, S. Mordijck³⁹, L. Moreira⁶, F. Moro⁸⁹, J. Morris⁷⁵, L. Moser²², S. Mosher⁷⁰, D. Moulton⁶, T. Mrowetz⁶, A. Muir⁶, M. Muraglia³, A. Murari^{11,30}, A. Muraro⁴², S. Murphy⁶, P. Muscat⁶, N. Muthusonai⁶, C. Myers⁷³, N. N Asakura⁶⁶, B. N'Konga⁹⁵, F. Nabais⁵², R. Naish⁶, J. Naish⁶, T. Nakano⁶⁶, F. Napoli⁸⁹, E. Nardon⁷, V. Naulin⁸², M.F.F. Nave⁵², I. Nedzelskiy⁵², G. Nemtsev⁸⁸, V. Nesenevich⁵³, F. Nespoli²⁸, A. Neto³⁷, R. Neu⁶⁰, V.S. Neverov⁶⁹, M. Newman⁶, S. Ng⁶, M. Nicassio⁶, A.H. Nielsen⁸², D. Nina⁵², D. Nishijima⁹⁸, C. Noble⁶, C.R. Nobs⁶, M. Nocente⁹⁷, D. Nodwell⁷⁵, K. Nordlund⁹⁹, H. Nordman²⁵, R. Normanton⁷⁵, J.M. Noterdaeme⁶⁰, S. Nowak⁴², I. Nunes⁵², T. O'Gorman⁶, M. O'Mullane¹⁹, M. Oberkofler⁶⁰, M. Oberparleiter²⁵, T. Odupitan⁶, M.T. Ogawa⁶⁶, M. Okabayashi⁷³, H. Oliver⁴⁴, R. Olney⁶, L. Omoregie⁶, J. Ongena⁵⁷, F. Orsitto¹⁰, J. Orszagh¹⁶, T. Osborne³⁹, R. Otin⁶, A. Owen⁶, T. Owen⁷⁵, R. Paccagnella¹¹, L.W. Packer⁶, E. Pajuste¹⁰², S. Pamela⁶, S. Panja⁴⁶, P. Papp¹⁶, G. Papp⁶⁰, V. Parail⁶, C. Pardanaud³, F. Parra Diaz⁷⁷, A. Parsloe⁶, N. Parsons⁶, M. Parsons⁷⁰, R. Pasqualotto¹¹, M. Passeri⁹³, A. Patel⁶, S. Pathak⁴⁶, H. Patten²⁸, A. Pau¹⁵, G. Pautasso⁶⁰, R. Pavlichenko⁴⁵, A. Pavone⁶¹, E. Pawelec⁴⁸, C. Paz Soldan³⁹, A. Peackoc³¹, S.-P. Pehkonen¹⁰⁷, E. Peluso⁹³, C. Penot⁵⁴, J. Penzo⁶, K. Pepperell⁶, R. Pereira⁵², E. Perelli Cippo⁴², C. Perez von Thun^{30,34}, V. Pericoli⁴⁹, S. Peruzzo¹¹, M. Peterka⁵⁰, P. Petersson³⁸, G. Petravich¹⁰⁸, A. Petre⁸⁴, V. Petržilka⁵⁰, V. Philipps³⁴, L. Pigatto¹¹, M. Pillon⁸⁹, S. Pinches⁵⁴, G. Pintsuk³⁴, P. Piovesan¹¹, W. Pires de Sa⁵¹, A. Pires dos Reis⁵¹, L. Piron^{6,11}, C. Piron¹¹, A. Pironti¹⁰, F. Pisano¹⁵, R. Pitts⁵⁴, V. Plyusnin⁵², F.M. Poli⁷³, N. Pomaro¹¹, O.G. Pompilian⁸⁶, P. Pool⁶, S. Popovichev⁶, M. Poradziński⁴⁹, M.T. Porfiri⁸⁹, C. Porosnicu⁸⁶, M. Porton⁶, G. Possnert¹⁸, S. Potzel⁶⁰, G. Poulipoulis¹⁰¹, T. Powell⁷⁵, V. Prajapati⁴⁶, R. Prakash⁴⁶, I. Predebon¹¹, G. Prestopino⁹³, D. Price⁶, M. Price⁶, R. Price⁷⁵, D. Primetzhof¹⁸, P. Prior⁶, G. Pucella⁸⁹, P. Puglia^{28,51}, M.E. Puiatti¹¹, K. Purahoo⁶, I. Pusztai²⁰, Th. Pütterich⁶⁰, E. Rachlew²¹, M. Rack³⁴, R. Ragona⁵⁷, M. Rainford⁶, P. Raj⁵⁵, A. Rakha⁵, G. Ramogida⁸⁹, S. Ranjan⁴⁶, C.J. Rapson⁶⁰, D. Rasmussen⁷⁰, J.J. Rasmussen⁸², K. Rathod⁴⁶, G.

Rattá⁵⁶, S. Ratynskaia⁸¹, G. Ravera⁸⁹, M. Rebai⁴², A. Reed⁷⁵, D. Réfy¹⁰⁸, J. Regaña²⁹, M. Reich⁶⁰, N. Reid⁷⁵, F. Reimold³⁴, M. Reinhart²⁹, M. Reinke⁷⁰, D. Reiser³⁴, D. Rendell⁶, C. Reux⁷, S.D.A. Reyes Cortes⁵², S. Reynolds⁶, D. Ricci⁴², M. Richiusa⁶, D. Rigamonti⁴², F.G. Rimini⁶, J. Risner⁷⁰, M. Riva⁸⁹, J. Rivero-Rodriguez⁹¹, C. Roach⁶, R. Robins⁶, S. Robinson⁶, D. Robson⁶, R. Rodionov⁸⁸, P. Rodrigues⁵², J. Rodriguez⁷⁵, V. Rohde⁶⁰, M. Romanelli⁶, F. Romanelli⁹³, J. Romazanov³⁴, S. Rowe⁶, M. Rubel³⁸, G. Rubinacci¹⁰, G. Rubino¹⁰⁵, L. Ruchko⁵¹, C. Ruset⁸⁶, J. Rzakiewicz⁶³, S. Saarelma⁶, R. Sabot⁷, X. Sáez⁵, E. Safi⁹⁹, A. Sahlberg¹⁸, G. Saibene³⁷, M. Saleem⁷⁵, M. Salewski⁸², A. Salmi¹⁰⁷, R. Salmon⁶, F. Salzedas⁵², U. Sann³⁴, D. Sandiford⁶, P. Santa⁴⁶, M.I.K. Santala¹, B. Santos⁵², A. Santucci⁸⁹, F. Sartori³⁷, R. Sartori³⁷, O. Sauter²⁸, R. Scannell⁶, F. Schluck³⁴, T. Schlummer³⁴, K. Schmid⁶⁰, S. Schmuck^{6,42}, K. Schöpf¹⁰⁰, J. Schweinzer⁶⁰, D. Schwörer²⁷, S.D. Scott⁷³, G. Sergienko³⁴, M. Sertoli⁶⁰, A. Shabbir¹⁴, S.E. Sharapov⁶, A. Shaw⁶, H. Sheikh⁶, A. Shepherd⁶, A. Shevelev⁵³, D. Shiraki⁷⁰, A. Shumack³³, G. Sias¹⁵, M. Sibbald⁶, B. Sieglin⁶⁰, S. Silburn⁶, J. Silva⁶, A. Silva⁵², C. Silva⁵², D. Silvagni⁶⁰, P. Simmons⁷⁵, J. Simpson^{1,6}, A. Sinha⁴⁶, S.K. Sipilä¹, A.C.C. Sips³¹, P. Sirén¹⁰⁷, A. Sirinelli⁵⁴, H. Sjöstrand¹⁸, M. Skiba¹⁸, R. Skilton⁷⁵, V. Skvara⁵⁰, B. Slade⁶, R. Smith⁶, P. Smith⁶, S.F. Smith¹⁰⁹, L. Snoj⁸⁰, S. Soare⁸⁵, E.R. Solano⁵⁶, A. Somers²⁷, C. Sommariva²⁸, P. Sonato¹¹, M. Sos⁵⁰, J. Sousa⁵², C. Sozzi⁴², S. Spagnolo¹¹, P. Sparapani¹⁵, T. Spelzini⁶, F. Spineanu⁸⁶, D. Sprada⁶, S. Sridhar⁷, G. Stables⁶, J. Stallard⁶, I. Stamatelatos⁶⁸, M.F. Stamp⁶, C. Stan-Sion⁸⁴, Z. Stancar⁸⁰, P. Staniec⁶, G. Stankūnas⁵⁸, M. Stano¹⁶, C. Stavrou¹, E. Stefanikova³⁸, I. Stepanov⁵⁷, A.V. Stephen⁶, M. Stephen⁴⁶, J. Stephens⁶, B. Stevens⁶, J. Stober⁶⁰, C. Stokes²³, J. Strachan⁷³, P. Strand²⁵, H.R. Strauss⁴⁰, P. Ström³⁸, W. Studholme⁶, F. Subba⁷², E. Suchkov¹⁶, H.P. Summers¹⁹, H. Sun⁶⁰, N. Sutton⁷⁵, J. Svensson⁶¹, D. Sytnikov⁴⁵, T. Szabolics¹⁰⁸, G. Szepesi⁶, T. T Suzuki⁶⁶, F. Tabarés⁵⁶, T. Tadić⁷⁶, B. Tal⁶⁰, B. Tál¹⁰⁸, T. Tala¹⁰⁷, C. Taliercio¹¹, A. Tallargio⁶, K. Tanaka⁶⁵, W. Tang⁷³, M. Tardocchi⁴², R. Tatali², D. Taylor⁶, D. Tegnered²⁵, G. Telesca^{14,49}, N. Teplova⁵³, A. Teplukhina²⁸, D. Terranova¹¹, C. Terry⁷⁵, D. Testa²⁸, E. Tholerus³⁸, J. Thomas⁷⁵, V.K. Thompson⁶, A. Thornton⁶, W. Tierens⁶⁰, I. Tiseanu⁸⁶, H. Tojo⁶⁶, M. Tokitani⁶⁴, P. Tolas⁸¹, M. Tomeš⁵⁰, P. Trimble⁶, M. Tripsky⁵⁷, M. Tsalias³³, P. Tsavalas⁶⁸, D. Tskhakaya⁸³, D. Tskhakaya jun⁸³, I. Turner⁶, M.M. Turner²⁷, M. Turnyanskiy²⁹, G. Tvalashvili⁶, M. Tyshchenko⁴⁵, A. Uccello⁴², J. Uljanovs¹, H. Urano⁶⁶, A. Urban⁶³, G. Urbanczyk⁷, I. Uytendhouwen⁷⁹, A. Vadgama⁶, D. Valcarcel⁶, R. Vale⁶, M. Valentinuzzi⁷, K. Valerii⁴⁵, M. Valisa¹¹, P. Vallejos Olivares³⁸, M. Valovic⁶, D. Van Eester⁵⁷, W. Van Renterghem⁷⁹, G.J. van Rooij³³, J. Varje¹, S. Vartanian⁷, K. Vasava⁴⁶, T. Vasilopoulou⁶⁸, M. Vecsei¹⁰⁸, J. Vega⁵⁶, S. Ventre¹⁰, G. Verdoolaege⁵⁷, C. Verona⁹³, G. Verona Rinati⁹³, E. Veshchev⁵⁴, N. Vianello¹¹, J. Vicente⁵², E. Viezzer⁹¹, S. Villari⁸⁹, F. Villone¹⁰, M. Vincent⁶, P. Vincenzi¹¹, I. Vinyar⁷¹, B. Viola⁸⁹, A. Vitins¹⁰², Z. Vizvary⁶, M. Vlad⁸⁶, I. Voitsekhovitch^{6,29}, D. Voltolina¹¹, U. von Toussaint⁶⁰, P. Vondráček⁵⁰, M. Vukšić⁷⁶, B. Wakeling⁶, C. Waldon⁶, N. Walkden⁶, R. Walker⁶, M. Walker⁷⁵, M. Walsh⁵⁴, N. Wang⁶, E. Wang³⁴, N. Wang³⁴, S. Warder⁶, R. Warren⁶, J. Waterhouse⁶, C. Watts⁵⁴, T. Wauters⁵⁷, M. Webb⁶, A. Weckmann³⁸, J. Weiland²⁰, M. Weiland⁶⁰, H. Weisen²⁸, M. Weiszflog¹⁸, P. Welch⁷⁵, A. West⁶, M. Wheatley⁶, S. Wheeler⁷⁵, A.M. Whitehead⁶, D. Whittaker⁶, A.M. Widdowson⁶, S. Wiesen³⁴, G. Wilkie²⁰, J. Williams⁶, D. Willoughby⁶, J. Wilson⁶, I. Wilson⁶, H.R. Wilson¹⁰⁹, M. Wischmeier⁶⁰, A. Withycombe⁶, D. Witts⁶, E. Wolfrum⁶⁰, R. Wood⁶, R. Woodley⁷⁵, C. Woodley⁶, S. Wray⁶, J.C. Wright⁶², P. Wright⁷⁵, S. Wukitch⁶², A. Wynn¹⁰⁹, L. Xiang⁸², T. Xu⁶, Y. Xue⁶, D. Yadikin²⁵, Y. Yakovenko⁴⁵, W. Yanling³⁴, V. Yavorskij¹⁰⁰, I. Young⁶, R. Young⁶, D. Young⁶, J. Zacks⁶, R. Zagorski⁴⁹, F.S. Zaitsev¹⁶, L. Zakharov⁹⁹, R. Zanino⁷², A. Zarins¹⁰², R. Zarins¹⁰², D. Zarzoso Fernandez³, K.D. Zastrow⁶, M. Zerbini⁸⁹, W. Zhang⁶⁰, Y. Zhou³⁸, E. Zilli¹¹, A. Zocco⁶¹, V. Zoita⁸⁶, S. Zoletnik¹⁰⁸, W. Zwingmann⁵², I. Zychor⁶³,

- ¹Aalto University, P.O.Box 14100, FIN-00076 Aalto, Finland,
- ²Aix-Marseille Univ., CNRS, M2P2, UMR 7340, 13013 Marseille, France,
- ³Aix-Marseille Univ., CNRS, PIIM, UMR 7345, 13013 Marseille, France,
- ⁴Arizona State University, Tempe, United States,
- ⁵Barcelona Supercomputing Center, Barcelona, Spain,
- ⁶CCFE, Culham Science Centre, Abingdon, Oxon, OX14 3DB, UK,
- ⁷CEA, IRFM, F-13108 Saint Paul Lez Durance, France,
- ⁸Center for Energy Research, University of California at San Diego, La Jolla, CA 92093, USA,
- ⁹Centro Brasileiro de Pesquisas Físicas, Rua Xavier Sigaud, 160, Rio de Janeiro CEP 22290-180, Brazil,
- ¹⁰Consorzio CREATE, Via Claudio 21, 80125 Napoli, Italy,
- ¹¹Consorzio RFX, corso Stati Uniti 4, 35127 Padova, Italy,
- ¹²Daegu University, Jillyang, Gyeongsan, Gyeongbuk 712-174, Republic of Korea,
- ¹³Departamento de Física, Universidad Carlos III de Madrid, 28911 Leganés, Madrid, Spain,
- ¹⁴Department of Applied Physics UG (Ghent University) St-Pietersnieuwstraat 41 B-9000 Ghent, Belgium,
- ¹⁵Department of Electrical and Electronic Engineering, University of Cagliari, Piazza d'Armi 09123 Cagliari, Italy,
- ¹⁶Department of Experimental Physics, Faculty of Mathematics, Physics and Informatics Comenius University Mlynska dolina F2, 84248 Bratislava, Slovak Republic,
- ¹⁷Department of Physics and Applied Physics, University of Strathclyde, Glasgow, G4 ONG, UK,
- ¹⁸Department of Physics and Astronomy, Uppsala University, SE-75120 Uppsala, Sweden,
- ¹⁹Department of Physics SUPA, University of Strathclyde, Glasgow, G4 ONG, UK,
- ²⁰Department of Physics, Chalmers University of Technology, SE-41296 Gothenburg, Sweden,
- ²¹Department of Physics, SCI, KTH, SE-10691 Stockholm, Sweden,
- ²²Department of Physics, University of Basel, Switzerland,

- ²³Department of Physics, University of Bath, Bath, BA2 7AY, UK,
- ²⁴Department of Physics, University of Warwick, Coventry, CV4 7AL, UK,
- ²⁵Department of Space, Earth and Environment, Chalmers University of Technology, SE-41296 Gothenburg, Sweden,
- ²⁶Dipartimento di Ingegneria Elettrica Elettronica e Informatica-Università degli Studi di Catania, 95125 Catania, Italy,
- ²⁷Dublin City University (DCU), Ireland,
- ²⁸Ecole Polytechnique Fédérale de Lausanne (EPFL), Swiss Plasma Center (SPC), CH-1015 Lausanne, Switzerland,
- ²⁹EUROfusion Programme Management Unit, Boltzmannstr. 2, 85748 Garching, Germany,
- ³⁰EUROfusion Programme Management Unit, Culham Science Centre, Culham, OX14 3DB, UK,
- ³¹European Commission, B-1049 Brussels, Belgium,
- ³²Fluid and Plasma Dynamics, ULB - Campus Plaine - CP 231 Boulevard du Triomphe, 1050 Bruxelles, Belgium,
- ³³FOM Institute DIFFER, Eindhoven, the Netherlands,
- ³⁴Forschungszentrum Jülich GmbH, Institut für Energie- und Klimaforschung - Plasmaphysik, 52425 Jülich, Germany,
- ³⁵Foundation for Research and Technology - Hellas, N. Plastira 100, Vassilika Vouton 70013, Heraklion, Crete, Greece,
- ³⁶Fourth State Research, 503 Lockhart Dr, Austin, Texas, United States,
- ³⁷Fusion for Energy Joint Undertaking, Josep Pl. 2, Torres Diagonal Litoral B3, 08019, Barcelona, Spain,
- ³⁸Fusion Plasma Physics, EES, KTH, SE-10044 Stockholm, Sweden,
- ³⁹General Atomics, P.O.Box 85608, San Diego, CA 92186-5608, California, USA,
- ⁴⁰HRS Fusion, West Orange NJ, USA,
- ⁴¹ICREA and Barcelona Supercomputing Center, Barcelona, Spain,
- ⁴²IFP-CNR, via R. Cozzi 53, 20125 Milano, Italy,
- ⁴³Institut Jean Lamour, UMR 7198, CNRS-Université de Lorraine, 54500 Vandoeuvre-lès-Nancy, France,

- ⁴⁴Institute for Fusion Studies, University of Texas at Austin, Austin, Texas, USA,
- ⁴⁵Institute for Nuclear Research, Prospekt Nauky 47, Kyiv 03680, Ukraine,
- ⁴⁶Institute for Plasma Research, Bhat, Gandhinagar - 382 428, Gujarat State, India,
- ⁴⁷Institute of Nuclear Physics, Radzikowskiego 152, 31-342 Kraków, Poland,
- ⁴⁸Institute of Physics, Opole University, Oleska 48, 45-052 Opole, Poland,
- ⁴⁹Institute of Plasma Physics and Laser Microfusion, Hery 23, 01-497 Warsaw, Poland,
- ⁵⁰Institute of Plasma Physics of the CAS, Za Slovankou 1782/3, 182 00 Praha 8, Czech Republic,
- ⁵¹Instituto de Física - Universidade de São Paulo Rua do Matão Travessa R Nr.187 CEP 05508-090 Cidade Universitária, São Paulo, Brasil,
- ⁵²Instituto de Plasmas e Fusão Nuclear, Instituto Superior Técnico, Universidade de Lisboa, Portugal,
- ⁵³Ioffe Physico-Technical Institute, 26 Politekhnicheskaya, St Petersburg 194021, Russian Federation,
- ⁵⁴ITER Organization, Route de Vinon, CS 90 046, 13067 Saint Paul Lez Durance, France,
- ⁵⁵Karlsruhe Institute of Technology, P.O.Box 3640, D-76021 Karlsruhe, Germany,
- ⁵⁶Laboratorio Nacional de Fusión, CIEMAT, Madrid, Spain,
- ⁵⁷Laboratory for Plasma Physics Koninklijke Militaire School - Ecole Royale Militaire Renaissancelaan 30 Avenue de la Renaissance B-1000, Brussels, Belgium,
- ⁵⁸Lithuanian energy institute, Breslaujos g. 3, LT-44403, Kaunas, Lithuania,
- ⁵⁹Magnetic Sensor Laboratory, Lviv Polytechnic National University, Lviv, Ukraine,
- ⁶⁰Max-Planck-Institut für Plasmaphysik, D-85748 Garching, Germany,
- ⁶¹Max-Planck-Institut für Plasmaphysik, Teilinstitut Greifswald, D-17491 Greifswald, Germany,
- ⁶²MIT Plasma Science and Fusion Centre, Cambridge, MA 02139, Massachusetts, USA,
- ⁶³National Centre for Nuclear Research(NCBJ),05-400 Otwock-Świerk, Poland,
- ⁶⁴National Institute for Fusion Science, Oroshi, Toki, Gifu 509-5292, Japan,
- ⁶⁵National Institute for Fusion Science, Toki, 509-5292, Japan,

- ⁶⁶National Institutes for Quantum and Radiological Science and Technology, Naka, Ibaraki 311-0193, Japan,
- ⁶⁷National Technical University of Athens, Iroon Politechniou 9, 157 73 Zografou, Athens, Greece,
- ⁶⁸NCSR "Demokritos" 153 10, Agia Paraskevi Attikis, Greece,
- ⁶⁹NRC Kurchatov Institute, 1 Kurchatov Square, Moscow 123182, Russian Federation,
- ⁷⁰Oak Ridge National Laboratory, Oak Ridge, TN 37831-6169, Tennessee, USA,
- ⁷¹PELIN LLC, 27a, Gzhatskaya Ulitsa, Saint Petersburg, 195220, Russian Federation,
- ⁷²Politecnico di Torino, Corso Duca degli Abruzzi 24, I-10129 Torino, Italy,
- ⁷³Princeton Plasma Physics Laboratory, James Forrestal Campus, Princeton, NJ 08543, New Jersey, USA,
- ⁷⁴Purdue University, 610 Purdue Mall, West Lafayette, IN 47907, USA,
- ⁷⁵RACE, Culham Science Centre, Abingdon, Oxon, OX14 3DB, UK,
- ⁷⁶Ruđer Bošković Institute, Bijenička 54, 10000 Zagreb, Croatia,
- ⁷⁷Rudolf Peierls Centre for Theoretical Physics, University of Oxford, Oxford OX1 3PU, UK,
- ⁷⁸School of Mathematics and Physics, Queen's University, Belfast, BT7 1NN, UK,
- ⁷⁹SCK-CEN, Nuclear Research Centre, 2400 Mol, Belgium,
- ⁸⁰Slovenian Fusion Association (SFA), Jozef Stefan Institute, Jamova 39, SI-1000 Ljubljana, Slovenia,
- ⁸¹Space and Plasma Physics, EES, KTH SE-100 44 Stockholm, Sweden,
- ⁸²Technical University of Denmark, Department of Physics, Bldg 309, DK-2800 Kgs Lyngby, Denmark,
- ⁸³Technische Universität Wien, Fusion@Österreichische Akademie der Wissenschaften (ÖAW), Austria,
- ⁸⁴The "Horia Hulubei" National Institute for Physics and Nuclear Engineering, Magurele-Bucharest, Romania,
- ⁸⁵The National Institute for Cryogenics and Isotopic Technology, Ramnicu Valcea, Romania,
- ⁸⁶The National Institute for Laser, Plasma and Radiation Physics, Magurele-Bucharest, Romania,

- ⁸⁷The National Institute for Optoelectronics, Magurele-Bucharest, Romania,
- ⁸⁸Troitsk Insitute of Innovating and Thermonuclear Research (TRINITI), Troitsk 142190, Moscow Region, Russian Federation,
- ⁸⁹Unità Tecnica Fusione - ENEA C. R. Frascati - via E. Fermi 45, 00044 Frascati (Roma), Italy,
- ⁹⁰Universidad Complutense de Madrid, Madrid, Spain,
- ⁹¹Universidad de Sevilla, Sevilla, Spain,
- ⁹²Universidad Politécnica de Madrid, Grupo I2A2, Madrid, Spain,
- ⁹³Università di Roma Tor Vergata, Via del Politecnico 1, Roma, Italy,
- ⁹⁴Universitat Politècnica de Catalunya, Barcelona, Spain,
- ⁹⁵Université Cote d'Azur, CNRS, Inria, LJAD, Parc Valrose, 06108 Nice Cedex 02, France,
- ⁹⁶University College Cork (UCC), Ireland,
- ⁹⁷University Milano-Bicocca, piazza della Scienza 3, 20126 Milano, Italy,
- ⁹⁸University of California, 1111 Franklin St., Oakland, CA 94607, USA,
- ⁹⁹University of Helsinki, P.O. Box 43, FI-00014 University of Helsinki, Finland,
- ¹⁰⁰University of Innsbruck, Fusion@Österreichische Akademie der Wissenschaften (ÖAW), Austria,
- ¹⁰¹University of Ioannina, Panepistimioupoli Ioanninon, P.O.Box 1186, 45110 Ioannina, Greece,
- ¹⁰²University of Latvia, 19 Raina Blvd., Riga, LV 1586, Latvia,
- ¹⁰³University of Texas at Austin, Institute for Fusion Studies, Austin, TX 78712, Texas, USA,
- ¹⁰⁴University of Toyama, Toyama, 930-8555, Japan,
- ¹⁰⁵University of Tuscia, DEIM, Via del Paradiso 47, 01100 Viterbo, Italy,
- ¹⁰⁶University of Wisconsin-Madison, Madison, Wisconsin 53706, USA,
- ¹⁰⁷VTT Technical Research Centre of Finland, P.O.Box 1000, FIN-02044 VTT, Finland,
- ¹⁰⁸Wigner Research Centre for Physics, P.O.B. 49, H - 1525 Budapest, Hungary,
- ¹⁰⁹York Plasma Institute, University of York, Heslington, York, YO10 5DD, UK,

Abstract:

For the past years, the JET scientific Programme [1] is engaged in a multi-campaigns effort including experiments in D, H and T, leading up in 2020 to the first experiments with 50%/50% D-T mixtures since 1997 and the first ever D-T plasmas with the ITER mix of plasma-facing component materials. For this purpose, a concerted physics and technology programme was launched in view of preparing the DT campaign (DTE2). This paper addresses the key elements developed by the JET programme directly contributing to the DT preparation. This intense preparation includes the review of the physics basis for the D-T operational scenarios, including the fusion power predictions through first principle and integrated modelling and the impact of isotopes in the operation and physics of D-T plasmas (thermal and particle transport, H-mode access, Be & W erosion, fuel recovery, etc.). This effort also requires improving several aspects of plasma operation for DTE2, such as real time control schemes, heat load control, disruption avoidance and mitigation system (including the installation of a new shattered pellet injector), novel ion cyclotron resonance heating (ICRH) schemes (such as the three-ions scheme), new diagnostics (neutrons camera and spectrometer, active Alfvén Eigen-modes antennas, neutral gauges, radiation hard imaging systems ...) and the calibration of the JET neutron diagnostics at 14 MeV for the accurate fusion power measurement. The active preparation of JET for the 2020 DT campaign provides an incomparable source of information and basis for the future DT operation of ITER and it is also foreseen that a large number of key physics issues will be addressed in support of burning plasmas.

1. INTRODUCTION

Since 2016, the JET scientific Programme is engaged in a multi-campaigns effort including experiments in D, H and T, leading up in 2020 to the first experiments with 50%/50% D-T mixtures since 1997 (DTE1 campaign [2, 3]), where 16MW of fusion power was achieved transiently and 4MW in steady state, and the first ever D-T plasmas with the ITER mix of plasma-facing component materials [4,5,6]. This effort is also driven by the EUROfusion research roadmap to secure the success of future ITER operation via specific preparation and experiments, including D-T operation of JET [7].

For this purpose, a concerted physics and technology programme was launched in view of preparing the second JET DT campaign (DTE2) [8]. This overview paper addresses the key elements developed by the JET programme directly contributing to the DT preparation. JET is a unique device in the sense that it has been designed from the start as a DT fusion tokamak with the aim to study plasma behavior in conditions and dimensions approaching those required in a fusion reactor and therefore has the capability to study the physics of alpha power. JET is equipped with a tritium plant and is capable of confining efficiently the alpha particles in the plasma (90% of alphas confined for plasma current above 2.5MA) thanks to its size and the plasma current it can reach (up to 5MA in the present configuration).

In addition, since DTE1 in 1997, the original carbon wall of JET has been changed to an ITER-like wall with a tungsten divertor and Beryllium first wall the total input power upgraded to 40MW and the set of diagnostics dramatically improved. Our goal is to reach 15MW of fusion power in stationary conditions in this environment.

This intense preparation for DT includes the review of the physics basis for the D-T plasma scenarios including the fusion power predictions through first principle and integrated modelling and the impact of isotopes on the operation and physics of D-T plasmas (thermal and particle transport, H-mode access, Be & W erosion, fuel recovery, etc). This also requires improving several aspect of plasma operation for DTE2, such as real time control schemes,

heat load control, disruption avoidance and mitigation system (including the installation of the new shattered pellet injector), dedicated ion cyclotron resonance heating (ICRH) schemes, new diagnostics (neutrons camera and spectrometer, active Alfvén Eigen modes (AE) antennas, neutral gauges, radiation hard imaging systems, etc...), new tritium injection valves, and the calibration of the JET neutron monitors at 14 MeV. The DT phase projects to reach a total of 40MW of input power (resonance ion cyclotron and neutral beam combined), a budget of up to $\sim 700\text{g}$ of reprocessed tritium gas and $1.7 \cdot 10^{21}$ 14MeV neutrons (compared to 35g and 3.10^{20} neutrons respectively for DTE1) [9].

The preparation for the DT campaign is reviewed in this paper along three main sections:

- 1- *The scenario development and the prediction for fusion power* are essential for optimizing the operational tools and reach the target of 15MW of fusion power for about 5s. This includes a continuous effort on modelling for predicting the fusion power in the DT phase, specific scenario development for the studies of alpha particle physics and the role of ICRH in fusion performance.
- 2- *The isotope physics* is being studied within a suite of campaigns in hydrogen, deuterium and full tritium as an indispensable preparation for DT to assess the effect of the isotopes mass on core and pedestal confinement, H-mode power threshold, particle transport, and plasma wall interactions.
- 3- *Operational preparation* includes a large set of items such as new diagnostics and tools for alpha physics studies, 14MeV neutron calibration, operational safety and procedures when using tritium in a fusion machine.

The preparation for the DT phase is led in an integrated way and also requires the scientific community to develop a unique platform for the study of isotope and fusion power in a first wall environment that is as close possible to the future ITER wall. The impact on the ITER research plan [10] of these developments will also be discussed within each section of this paper

2. SCENARIO DEVELOPMENT AND PREDICTION FOR DT

2.1 Analysis of scenario development for DT

In view of the preparation for DT, developing the physics basis for the integrated scenarios is paramount in order to achieve the fusion power target of 15MW for 5s [11] and for ensuring clear observation of alpha-particle effects and allow their detailed study. Two complementary lines of research are followed for developing scenarios suitable for sustained high D-T fusion power over 5s [12]: the baseline scenario (with $\beta_N \sim 1.8$ and $q_{95} \sim 3$) [13] and the hybrid scenario (β_N up to 3 and $q_{95} \sim 4$) [14] (fig 1). The baseline scenario focuses mainly on type I ELMy H mode at high current and toroidal field operation with a relaxed current profile whereas the hybrid experiments addresses operation at high β_N with a shaped current profile and q_0 close to or above unity. Both are aiming at achieving stationary conditions for 5s.

In 2016 during and just after the last 2016 Fusion energy Conference, encouraging results were achieved for the baseline scenario at 3 MA/2.8 T with injected power of ~ 28 MW of NBI and ~ 5 MW of ICRH (fig 1 right picture). $\sim 3.10^{16}$ neutrons/s could be achieved for >5 energy confinement times (~ 1.5 s) making an equivalent fusion power of ~ 7 MW as computed with TRANSP [15] shared by 40% beam-target and 60% thermal-thermal fusion power.

Here, the equivalent fusion power for a DT pulse is computed by TRANSP as in [16] assuming equal power from neutral beam in deuterium and tritium in the deuterium scenario considered. In addition, in the JET calculation, the neutral beam fractional energies are taken into account and deuterium and tritium concentration are forced to be the same (50% Deuterium-tritium mix). In these calculation, no credit is taken for the alpha heating produced and possible favourable isotopic effects on confinement.

These equivalent fusion power performances have been achieved by lowering the injected gas rate at high power, thus accessing lower collisionality in the core and achieving high rotation at the H-mode pedestal. Lower particle throughput has been achieved by means of a combination of gas and ELM pacing pellets injection, which resulted in moderate high Z impurity accumulation with a better confinement than with gas fuelling alone. Together with the increase of input power (up to 33MW), lower collisionality helped in decoupling the ion and electron channels in the core and higher T_i/T_e also induced a positive feedback on the stabilisation of ion temperature gradient turbulence. The positive feedback was stronger at high rotation, which was enabled by low gas injection [17]. The operation with the baseline scenario also confirmed that ICRH power up to 5MW, aided by an optimised coupling to the plasma with appropriate edge fuelling is essential to control the accumulation of high Z impurity in the plasma core (more details below in section 2.4)

Similar results in terms of neutron yield and equivalent fusion power were obtained at reduced plasma current but higher normalised beta in the hybrid scenario (2.2-2.5 MA/2.8-2.9 T) (fig 1 left picture). In this scenario, real time control of the ELM frequency with gas injection has been introduced to help in flushing tungsten from the edge. ICRH core deposition also helps in controlling the electron density peaking which could in turn lead to W accumulation [18]. Particularly for the hybrid scenario, heavy impurity accumulation is driven by neoclassical convection enhanced by poloidal asymmetries, and is highly sensitive to the main ion density and temperature peaking. Multi-channel predictive modelling of the high performance hybrid scenario ($T_i > T_e$) reproduces well central tungsten and nickel accumulation 1.5s after the H-mode onset, as observed in experiment [19]. The high-Z impurities are controlled by application of high power density ICRH power near the plasma axis by enhancing turbulent diffusion [20]. As reported previously [21], tearing modes can also impact the discharge performance, therefore control was attempted by q profile tailoring. At $\beta_N=2.4$ (feed-back controlled using NBI power) m=1 MHD activity and tearing modes can be avoided for 3.5 s using q profile tailoring by means of beam timing and current ramps and will be further optimized the experimental campaigns in 2019. Further analysis using quasi-linear codes have shown that low density conditions in the hybrid scenario are an advantage for boosting the neutron rate generation [22]. In the hybrid scenario, enhanced fusion power is explained by the higher penetration of the NBI beams to the plasma core and a reduced Ion Temperature Gradient (ITG) turbulence by fast ions when electromagnetic effects are taken into account [23, 24].

In both scenarios, strike point sweeping of 3.5cm on the divertor tile is used and was proven efficient to mitigate the power peak heat load with $P_{IN}=30$ MW for 5s [25]. Neon seeding has also been attempted as an additional method to mitigate the divertor heat load. Although efficient in reducing the temperature of the divertor target plates, the neon had the detrimental effect of increasing the central density thus reducing the central temperature and resulting in a non-negligible penalty on the fusion yield [26]. Strike point sweeping is at present the main method to handle high exhaust power, but the use of neon seeding cannot be ruled out to reach the target of 5s and is being considered in tritium plasmas if tungsten sputtering by tritium becomes intolerable [27]. Nitrogen seeding cannot be used in the JET tritium campaigns because the JET gas handling system does not handle tritiated ammonia and could also contaminate the uranium beds of the JET tritium plant.

To fulfil the mission of alpha physics in the DT phase, a third plasma scenario has been developed in view of the DT phase [28]. In next-step devices, including ITER, the impact of alpha-driven TAEs on the redistribution of fast ions causing a degradation of the plasma performance and losses to the first wall remains to be quantified. It is therefore essential to prepare scenarios aimed at observing α -driven TAEs in a future JET DT campaign. The main challenge for this type of study is to overcome the strong Landau damping from the neutral beams [29]. Discharges at low density, large core temperatures associated with the presence

of ITBs (Internal Transport Barriers) and good energetic ion confinement (i.e. $I_p > 2.5\text{MA}$) have been performed in plasma with elevated q profile (q_{\min} from 1.5 to 3) in JET (fig 2a). As in TFTR [30], the afterglow scheme has been developed, consisting in switching-off the auxiliary NBI power abruptly, and relying on the faster decay of the fast NBI ions compared to the fusion alphas to observe the α -driven TAEs in the afterglow phase where alpha heating will be therefore dominating transiently. In tests of this strategy in deuterium plasmas, the presence of MeV ions driven by ICRH power has resulted in the experimental observation of $n=4,5,6$ TAEs (fig 2b), also predicted by stability calculation at $\rho \sim 0.4$ using the MISHKA code [31] (fig 3). Extrapolating this plasma to DT shows that the obtained $\beta_{T\alpha}$ achieved should be comparable, or even slightly larger, than what was achieved in similar successful TFTR experiments. This DT prediction has been used for stability calculation using MISHKA and HAGIS [32] and core TAE with toroidal mode numbers $n=4,5,6$ have been found thus matching those observed in the DD ICRH version of the pulse. The computation also predicts an alpha drive of typically 0.3% which is comparable to what was found in previous work [29]. These results give confidence that alpha driven TAE could be observed in the future DT phase of JET. In addition, the damping rates of Alfvén Eigen modes are also measured experimentally using the newly refurbished Alfvén Eigen-modes Active Diagnostic (AEAD) [33]. Good agreement has been obtained between simulations of damping rate with the GTC code and the measurements for weakly damped AEs probed by the antennas [34]. It should be noted that the impact of instabilities such as fishbone has also been studied using the HAGIS code and shows that no more than 1% of alpha losses would be incurred by fishbones in JET D-T experiments [35] This is particularly important for the hybrid scenario, which is often subject to significant fishbone 1/1 activity in the plasma core.

2.2 Scenario termination and disruptivity

For all the future DT scenarios presented above, the control of impurity content is particularly demanding at the transition from the H-mode to the L-mode and during plasma termination [36] and will also be critical in ITER [37]. Dedicated experiments were conducted in JET, specifically designed to examine the evolution of plasma parameters during the H-mode termination phase, the conditions under which W accumulation develops and how it can be controlled with external actuators that are known to affect impurity transport, such as central electron heating or active ELM control (through pellet ELM pacing and vertical kicks) [38]. The experiments show that maintaining ELM control during the exit from the H-mode phase is a very effective way to avoid W accumulation and to achieve a smooth and well controlled termination. Integrated fully predictive core-edge-SOL transport modelling studies applying discrete models for the description of ELMs have been performed for the entire transition from stationary H-mode until the time when the plasma returns to L-mode, focusing on the W transport behaviour. Simulations have shown that the existing models appropriately reproduce the plasma profile evolution in the core, edge and SOL as well as W accumulation trends in the termination phase of JET H-mode discharges as function of the applied ICRH and ELM control schemes [39]. These studies have prompted the development of specific disruption avoidance schemes [40] and controllers using operation scenario modelling for handling the termination phases and minimise the risk of disruption at the H-L transition [41].

2.3 Fusion predictions for DT scenarios

DT performance predictions are particularly challenging since they require taking into account a large number of non-linear effects, including the heating deposition (both NBI and ICRH), fuelling sources, isotopic effects, fast particles, rotation and pedestal-core synergy.

For the calculation of the fusion power prediction, the initial task has been the development of the interpretative analyses of the reference discharges in Deuterium plasmas.

These are modelled using quasi-linear models (such as TGLF [42] or QUALIKIZ [43]). The impact of the different heating sources and mix on the performance are studied by suite of codes such as JINTRAC [44] and CRONOS [45].

On this basis, DT prediction have been carried out using the full JET power (NBI=32MW and ICRH=8MW) and CRONOS (TGLF) for both hybrid and baseline scenarios [46] using the pedestal scaling in $m_i^{-0.2}$ given in [47] (fig 4). The inclusion of the ExB shearing rate can generate an isotope effect which improves confinement, as also obtained from gyro-kinetic simulations [48]. For the third scenario dedicated to the study of the alpha physics which exhibit an internal transport barrier (section 2.2), the transport simulations used the CDBM model [49]. Additional predictions to DT have been using the Bohm-Gyro-Bohm model for describing the core turbulence and the core-edge self-consistent interplay using EUROped [50]. According to these predictions made on the basis of the pulses shown on figure 1 and 2a, 12 to 15MW of fusion power can be produced at full input power for both the baseline and hybrid scenario [46]. In addition, the predictions are showing that the electron heating in the plasma core by fusion alphas (~2-3 MW) will be comparable to electron heating by NBI, allowing a demonstration of alpha heating and other alpha particle effects (e.g. impurity screening, ITG stabilisation). The scenario for alpha physics studies (see above) could approach a fusion peak power of 10MW (stars on figure 4). A separate prediction made with JETTO (QUALIKIZ), [51] did not use a pedestal scaling dependence upon the isotope mass but used stiff electron temperature gradient (ETG) transport in the core. Despite these different physics hypotheses, both predictions are giving very close estimate of the fusion power for both hybrid and baseline scenario.

The equivalent fusion power has been also computed over a large database of different scenarios by a simpler fusion power simulator that approximates the trajectory of the JET NBI system as a single pencil in the plasma equatorial plane for the beam deposition calculation, and determines the average fast ion slowing down rate ignoring finite thermal ion temperature effects. The orbit averaged radial displacement due to the first orbit is estimated, after which orbit effects and radial transport are ignored. Fusion reaction rates are calculated for beam-target and thermal reactions, but beam-beam reactions and the effects of ICRF acceleration of fast ions (see section 2.4 below) are not included. As for the TRANSP calculation described in section 2.1, this calculation takes no credit on the additional alpha heating nor on possible favourable isotope effect. It has also been validated by TRANSP using the two pulses shown on figure 1, and represented on figure 5 by the circled points in black. This figure allows a direct comparison with the fusion power achieved in DTE1 in various scenarios. Despite the lower total energy content but with more input power (up to 40MW) and less dilution, the fusion power calculations are predicting that JET has already reached larger equivalent fusion power than the achieved fusion power in DTE1 for the baseline scenario (red circles and triangles respectively). In addition, the stationary scenarios (hybrid and baseline) are predicted to produce equivalent fusion power comparable to the transient fusion power achieved in hot-ion H-mode in DTE1 [2] (orange open triangles) despite having significant less total energy content. This difference can be explained both by the increased power (22MW of input power in DTE1 instead of 40MW planned for DTE2) and the lower dilution with the metallic wall (typically $Z_{eff}=1.5$ with the metallic wall instead of 2.5 with the carbon wall in DTE1).

Recently, predictive capability for D-T plasmas has been further enhanced by exploiting the European Transport Simulator (ETS) workflow developed by EUROfusion. ETS allows to solve the transport equations in plasmas including all hydrogenic species at the same time. The NBI and ICRH source modules in ETS have been thoroughly benchmarked with TRANSP and validated against JET data. [52]. Using this tool, the analysis and modeling of Neoclassical Tearing Modes (NTMs) onset and their effect on heavy impurity transport has been performed during the main scenario [53].

2.4 The role of ICRH power in the development of DT scenario.

The ICRH power plays a particular role in the preparation for DT and in the fusion power optimization. With the requirement of low gas puffing to preserve edge/pedestal confinement, ICRH heating is crucial for ensuring the control of high-Z impurity in the core of D-T scenarios. Fundamental hydrogen minority heating ($N=1$ H coinciding with $N=2$ D resonance) ICRH was extensively used in the baseline and hybrid experiments aiming at preparing DT scenarios. The $N=1$ H/ $N=2$ D ICRH scenario has been shown to maximize central electron heating thus providing peaked temperature profiles (figure 6, top picture) and increased turbulent transport to prevent W accumulation [54, 55, 56]. ICRH coupling is aided by dedicated gas injection in the main chamber magnetically connected with the ICRH antennas [57]. Although ICRH is the main tool used to prevent W accumulation in the centre of JET plasmas, application of ICRH can lead to an overall increase of the plasma impurity content, and in particular in JET-ILW, tungsten (W) and nickel (Ni). This is in general attributed to enhanced Plasma Wall Interactions (PWI) and sputtering of the Plasma Facing Components (PFC) when applying ICRH. In JET less high-Z impurities influx is observed with the so-called ITER-like Antenna (ILA) [58] in comparison with the other antennas (A2 antennas). The enhanced plasma wall interaction and the recipes to suppress their effects have been studied experimentally as well using antenna and RF-sheath models applied to JET [59, 60], such as the SSWICH code [61] and conclude that plasma facing components on the antenna (such as the septum) could be unfavourable for minimizing the RF-sheath and the impurity production.

By lowering the H concentration on the $n=1$ H/ $n=2$ D ICRH scheme, $N=2$ deuterium absorption in neutral beams ions becomes dominant and the D-D fusion rate is enhanced with respect to the higher H concentration cases [62]. Modelling has shown that this effect is essential to reproduce the neutron rate and will contribute to boost the D-D neutron yield in the range of 10-25% in the hybrid scenario by channelling the maximum ICRF power to Deuterium. However, this effect may not play such a large role in DT since DT fusion cross sections peaks at lower energy (~ 100 keV) than the DD cross section. The calculations are showing that the fusion enhancement due to ICRH in DT is about 5% in the hybrid scenario and negligible in the baseline scenario [63].

$N=1$ ^3He minority heating scheme has also been tested (figure 6 bottom picture) for enhancing ion heating (in D plasmas), although more experiment would be needed to conclude on its efficiency to prevent W accumulation. However, high ion heating power density near the axis can maximise the beneficial effect of ICRH against tungsten central peaking and is predicted to provide a stronger effect than electron heating. [51]. In TT and DT plasmas, ion temperature peaking is predicted to be further amplified by the decrease in the ion-electron collision coupling, giving a positive isotope scaling for both stored energy and impurity screening. Efficient ion and electron heating was also obtained when using a combined heating scheme where both H and ^3He minorities are simultaneously heated in the plasma centre by splitting the ICRH power between two frequencies [64].

Fundamental D minority heating in T-rich plasmas led to very efficient fusion performance in DTE1 ICRH-only plasmas [65]. Modelling indicates that this heating scheme applied to T-rich H-mode plasmas with neutral beam ions has a strong potential for achieving high DT fusion yield and preliminary conceptual studies for using it in DTE2 are ongoing. Further experiments are planned in the next campaign with deuterium and tritium plasmas to validate the ICRH scenarios foreseen for DTE2 from the point of view of fusion performance and prevention of core W accumulation.

Energetic species, such as injected NBI ions and fusion products, can also play the role of the 'third' species and resonate between the two cyclotron layers of the main ions, because of the Doppler shift in their resonance position. Indeed, effective ICRH heating of H-D mixed

plasmas using D-NBI ions as resonant species was recently demonstrated on JET [66]. Figure 7a shows an overview of JET pulse #91256 (2.9T/2MA, H-D \approx 85%-15%), where the neutron rate was increased by a factor of 10 to 15 when 2.5MW of ICRH power ($f = 25$ MHz, dipole phasing) was added to 3.5MW of D-NBI. The presence of a population of energetic D ions with energies $\sim 1\text{--}2$ MeV during the combined ICRH+NBI phase was confirmed by neutron spectrometry (TOFOR) and γ -ray measurements. A consistent simulation of the TOFOR measurements for this advanced heating scenario was done with the TRANSP and SCENIC codes [67] (figure 7b). This novel three-ion ICRH scenario is an efficient technique for heating plasma mixtures. As a proof-of-principle experiments on the tokamaks JET and Alcator C-Mod, a very small amount of ^3He ions absorbed nearly all launched radiofrequency (RF) power and provided an efficient heating of the background H-D mixed plasma [68]. In D-T plasmas, intrinsic ^9Be impurity ions will inherently absorb part of the ICRH power and lead to bulk ion heating. Depending on the natural Be concentration in the discharge, the D-T plasma mix could be adjusted to align the ion-ion hybrid layer closer to the Be cyclotron resonance layer to enhance the Be absorption (3-ion scheme). Due to their larger atomic mass, ICRH-heated ^9Be impurities will transfer an even larger fraction of RF power via Coulomb collisions to bulk D and T ions than the lighter ^3He minority [69]. Off-axis fundamental ICRH heating of T-NBI ions from $\langle E_T \rangle \approx 70$ keV (injection energy 118 keV) to $\langle E_T \rangle \approx 200$ keV can become a promising scenario to maximize the D-T fusion reactivity and steady-state Q -value in future DTE2 experiments.

3. DT ISOTOPE PHYSICS IN SUPPORT OF DT

The impact of isotope on confinement has also been assessed in details in the analysis of the hydrogen and deuterium campaigns [70] in both L-mode and H-mode and the future tritium campaign will uniquely complement the exploration of the isotope physics in fusion plasmas in support of the DT mix phase.

3.1 L-H-threshold dependence with isotope and divertor geometry

L-H power threshold ($P_{\text{sep}} = P_{\text{loss}} - P_{\text{rad}}$) studies have made use of about 200 pulses in JET-ILW spanning a range of plasma magnetic geometries, density and toroidal magnetic field values, hydrogen isotopes, ion species mixtures, effects from impurity seeding, and differences in heating and momentum sources [71]. It is notable that vertical target (VT) and corner/corner (CC) divertor geometry have about the same L-H power threshold, which is roughly a factor of two larger than in the horizontal target (HT) case possibly because of the different X-points height that are modifying the neutral circulation in the scrape-off [72]. In VT and C/C configurations, results are consistent with the other experiments finding that P_{LH} in Hydrogen is about twice that in Deuterium ($1/m_i$ dependence), for the high-density branch [73] (fig 8). This would confirm that in future full tritium and DT pulses the power threshold for the transition could be decreased by the same scaling as illustrated on figure 8, as it was also observed in with JET in carbon during the DTE1 campaign in 1998 [74].

For horizontal target data, access to the low-density branch is reached in both H and D, and significant differences depending on the heating methods (NBI or ICRH) have been identified. With the same shape, toroidal magnetic field, and plasma current, an isotope dependence is found for the density minimum of P_{LH} . Also in hydrogen, $P_{\text{L-H}}$ is much higher with NBI than ICRH, while there is little difference in D, which is similar to DIII-D results on the effect of torque [75]. The existence of a minimum density has also been investigated as in ASDEX Upgrade [76]. The power coupled to the ions computed from the JINTRAC suite of code [44] does not show a linear dependency with respect to density in the low density branch [77].

$P_{\text{L-H}}$ was also studied in mixed species plasmas, yielding unexpected results. It was found that most of the variation of P_{LH} in H-D mixtures was for H concentration less than 20% or

more than 80%, with little variation in between. Helium-4 fuelling into H plasmas was also performed, resulting in a ~25% reduction of the threshold with up to about 10% helium concentration. This reduction in L-H threshold in H-He mixtures may have beneficial application for the non-active phase of ITER operations.

The formation of the edge transport barrier (ETB) with isotope has been investigated with the four-field drift-fluid model HESEL [78,79]. The model can capture the observed isotope mass dependence of P_{LH} ($P_{LH} \propto A^{-1.4}$) linked with a faster development of the shear flow for higher mass but does not reproduce the behaviour of P_{LH} dependence with the H-D mixtures mentioned above.

3.2 Confinement and transport properties with the ion mass in L-mode and H-mode.

The results in L-mode are of importance because of its impact on ITER access condition to the burning phase and the current ramp-up in ITER. Confinement studies have found a weak dependence of $\tau_{E,th}$ on isotope mass demonstrated ($\sim A^{0.15 \pm 0.05}$) not following Gyro-Bohm transport ($\sim A^{-0.5}$) (fig 9). This near independence with the ion mass has been confirmed experimentally using NBI heated L-mode plasmas in both H and D, with matched profiles of the dimensionless plasma parameters, ρ^* , v^* , β and q in the plasma core confinement region [80]. Predictive core transport modelling with JETTO-TGLF [81,44] of the H and D identity pair is in good agreement with experiment for both isotopes, although H over-predicted in non-identity discharges. The stiff core heat transport, typical of JET-ILW NBI heated L-modes, overcomes the Gyro-Bohm scaling of gradient-driven TGLF, explaining the lack of isotope mass dependence in the core confinement region. The effect of ExB shearing on the predicted heat and particle transport channels is negligible for these low beta and low momentum input plasmas [70]; Other effects, linked with the isotope mass, may also break the gyro-Bohm scaling such as the impact of fast ions generated by the neutral beam on micro-instabilities [82, 83, 84], the multi-scale effect i.e. the interplay between ion and electron micro-instabilities of different scales [85], or collisionality [46].

In type I ELMy H-mode, establishing the same densities in hydrogen as in deuterium could not be achieved. This resulted in hydrogen densities up to a factor 2 below their deuterium counterparts showing a more significant reduction of particle confinement than in L-mode [80]. This behaviour was not observed in JET-C (carbon wall), nor in JT60-U [86] suggesting that the different wall recycling conditions [87] may play a role on the pedestal. Using kinetic measurements, a global scaling of confinement with the ion mass, $\tau_E \propto A^{0.4 \pm 0.1}$ in type I ELMy H-modes is derived assuming $T_i = T_e$ [80] (fig 10). The $A^{0.4}$ dependence of $\tau_{th,e}$ on the isotope mass is robust against which combination of variables is chosen in the regression of the data [88]. The mass dependence is stronger than the standard ITER scaling IPB98(y,2) ($\tau_E \propto A^{0.19}$) [89] and opposite in exponent sign to the Gyro-Bohm scaling. If confirmed with tritium, this strong scaling has favourable implications for ITER and for future D-T mixture plasmas. Particle confinement and angular momentum confinement are also found to scale strongly with ion mass ($\tau_\phi, \tau_p, \propto A^{0.5}$) [88].

As in L-mode, H-modes has resilient profiles with $8.5 < R/L_{Ti} < 10.5$ at mi-radius, T_i profile data show that T_i varies between $0.9T_e$ and $1.4T_e$ depending on density and power, and ITG's are dominant in the plasma core, as found by linear gyro-kinetic calculations with the GENE code [88]. The combination of larger pedestal pressure in deuterium and similar core R/L_T for both H and D suggests that, for this type of ELMy H-modes dataset, the favourable isotope effect on confinement originates in the pedestal and propagates to the plasma core via a near-constant critical temperature gradient. However, the effect of the isotope species on the pedestal is not yet understood and is under investigation.

3.3 Particle transport with isotope: isotope mixing effect.

In JET experiments, experiments with mixed isotopes, have provided a new insight on particle transport properties. In a multi-ion plasma, additional complexity is indeed introduced where ion profiles may respond to transients at a different timescale than the electrons. Two pulses with similar kinetic profiles T_e , T_i , n_e but different isotope compositions ($n_H/(n_H+n_D)$ ~ 0.86 and 0.33 respectively) have been produced [90]. Plasma isotope ratio H/D is controlled by changing the H/D ratio by gas dosing. Both pulses are type I ELMy H-modes with equal (8MW) deuterium NBI heating producing strong core fuelling of pure deuterium. Despite the strong core deuterium fuelling by the neutral beams in both discharges, it is observed that the hydrogen profile remains as peaked as the deuterium profile even without the core particle source (fig 11), implying the presence of fast transport of the individual ion components. Such a behaviour has also been observed in JET using pellet injection instead of gas injection to control the H/D ratio and in trace tritium experiments [91] and in other devices in the past [92,93]. In the ion Temperature Gradient (ITG) dominated turbulent regime (the case in most existing tokamaks and future devices) non-linear and quasilinear simulations show that both diffusion coefficient ratio D_i/D_e and pinch velocity ratio V_i/V_e can become much larger than unity for normalised gradient length larger than ~ 6 [94], hence changes in isotope composition can occur at much shorter time scales than changes in electron density profiles. The experiments described above have been modelled with the JETTO integrated modelling suite [44] with QuaLiKiz for turbulent transport and T_e , T_i and n_e predictions are in good agreement with experiments [95,96] and support the statement that sources have a negligible impact on density peaking for multi-ion plasmas. In addition, electron density peaking has also been observed when Neon is used as a seeding gas and is explained by the same physics process [97]. This new effect is likely to have an important consequence for DT fuelling and He ash removal in the JET DT phase and later in ITER [98].

For single ion plasma, on the other hand, core density profile peaking and electron particle transport have also been studied by performing several dimensionless collisionality (ν^*) scans in various plasma operation scenarios on JET with gas puff modulation technique using the high resolution diagnostics, allowing to distinguish between the NBI source and inward pinch in contributing to density peaking [99]. The NBI particle source is contributing 50-60% to the peaking in plasmas where $T_e/T_i \sim 1$ and at $\nu^* = 0.1-0.5$ (averaged between $r/a = 0.3-0.8$) independent of ν^* . These dimensionless ν^* scans give the most appropriate data for model validation. TGLF simulations are in good agreement with the experimental results with respect to the role of NBI particle source versus inward pinch in density peaking. GENE predicts flat or hollow density profiles for JET H-mode plasmas, thus giving higher weight on NBI fuelling than experimentally observed. For $T_e/T_i \sim 1.5$ and low β_N H-mode conditions, both TGLF and GENE predicts correctly the peaked n_e profile. Overall the various modelling results gives good confidence that these models can be used to predicting density peaking in lower ν^* plasmas.

3.4 Pedestal and ELM stability dependence with isotope mass.

Fusion power is affected by the pedestal performance through its structure and its stability. This is in large part caused by the electron temperature and density pedestals relative radial positions that tend to vary with plasma conditions as reported in JET [100] and other device [101]. In addition, JET-ILW tends to have a larger relative shift compared to JET-C, suggesting a possible role of the plasma facing materials in affecting the density profile location [102]. Stability analysis using EUROPED model [103] shows an improvement of the pedestal stability, when the relative shift is reduced. This has been mainly ascribed to the increase of the edge bootstrap current [104]. For the optimization of the DT scenario performance, the dependence of the pedestal structure on isotope mass is therefore an essential ingredient. A recent estimate of the impact of the density position on a ITER

baseline scenario shows that the maximum reduction in the pedestal height is 10% while the reduction in the fusion power is between 10% and 40% depending on the assumptions for core transport model used [105].

Recent JET-ILW studies [70] have identified several differences in the pedestal structure in hydrogen and deuterium plasmas. The transition from L-mode to type III ELMs and type III ELMs to type I ELM is strongly impacted by the isotope mass, in agreement with previous results in the JET-C. Also, identical discharges at the same input power and gas injection rate show a lower pedestal pressure in hydrogen than in deuterium. This is primarily due to the lower density (and higher particle transport in hydrogen) as observed in L-mode (section 3.1 above). The ELM frequency is also lower in D. Larger ELM frequency in H than D may indeed contribute to lower density and thus lower pedestal confinement [106]. Neutral penetration alone does not explain the differences between H and D pedestal density alone and transport must be invoked as well. Neutrals can indeed impact on stability calculation [105], however, the direct isotope effect on linear Peeling-Ballooning stability is small. Stability calculation are also strongly sensitive to the separatrix temperature, which may not be identical in H and D.

The transport mechanisms that produce energy and particle transport in the pedestal have been investigated through gyro-kinetic simulations, and analysis of the relative size of heat and particle fluxes in JET-ILW pedestals. Recent work [107] has attempted to constrain the interpretation of pedestal instabilities on the basis of contributions to different transport channels. Simulations of JET-ILW [44, 108] find that particle transport in the pedestal is small compared to heat transport, which may be consistent with kinetic ballooning mode marginal stability moderating the density profile, but that other mechanisms may be needed to explain heat transport. Non-linear gyro-kinetic simulations of high current (3 MA) baseline scenario JET-ILW plasmas find that electron temperature gradient turbulence as well as ion-scale turbulence driven by ion temperature gradient modes can be significant contributions to heat transport in the pedestal [109].

3.5 Dependence of plasma wall interaction with isotope mass.

After the ITER-like wall has been installed, C levels reduced by more than one order of magnitude. This low carbon level has not changed, despite high power operation on W-coated CFC tiles in the divertor, suggesting no significant change to the tungsten tile coatings. The future DT experiment are therefore not likely to experience strong tritiated hydro-carbon production contrary to DTE1 [110].

The erosion of W by plasma (H/D/T) and intrinsic (Be) and extrinsic impurity (such as seeded gas, N, Neon or Argon) sputtering determines both the lifetime of divertor components and the impact on plasma performance as it governs the W impurity source, thus, affects W influx to the confined plasma region determined also by other factors, for instance screening. Sputtering yields and fluxes as function of impact energy are strongly dependent on the plasma conditions (such as edge temperature) and should also include the interplay between intra- and inter-ELM. The isotope effect on Be erosion by hydrogenic ions is so far relatively unstudied since most of dedicated JET experiments were in D plasma. However, binary-collision approximation SDTrimSP simulations [111] allow predicting the potential isotope dependences for the physical sputtering (fig. 12) and thus give an insight of how sputtering may impact the integration of the main DT scenarios with the W/Be wall. In particular, it is apparent that in D-T and T-T, tritium impact on W sputtering will become larger inter ELMs and in L-mode. Intra-ELM erosion by tritium can increase by an order of magnitude with respect to hydrogen, and may amplify the level of tungsten sputtering by Be impurity [112,113].

Erosion and deposition were studied in the JET divertor during the JET-ILW campaigns using marker tiles analysed before and after the campaigns. The erosion/deposition patterns

show profound changes compared to JET-C. The main sink for impurities is at the inner divertor, where the most deposition and fuel retention occurs [114]. The total material deposition rate in the divertor decreased by a factor of 4-9 compared to the deposition rate of carbon in JET-C and the deuterium retention in the divertor decreased globally by a factor of 10-20. Tritium retention is expected to follow the same pattern as confirmed by the analysis of tritium from D-D reaction in plasma facing components [115]. Deposits on the inner divertor consist predominantly of beryllium (typically 5.0×10^{19} at/cm² of Be to 0.3×10^{19} at/cm² of D) [116] with 5-20% of carbon and oxygen, respectively and small amounts of Ni and W [117]. Charge exchange neutral erosion is found to be the main source of nickel, whereas erosion of divertor plasma facing components is the main source of tungsten. The tungsten sources come mainly from the outer divertor (10^{23} atoms per campaign), the top of the outer divertor, (10^{21} - 10^{22} atoms) and coating imperfections and fatigue cracking. The deposition rates in the inner and outer divertor corners and in remote areas decreased substantially and deposition is often only observed in valleys of the rough tile surfaces. These findings are consistent for all three ILW campaigns despite the considerable differences in strike point positions and heating powers [118].

In JET with the ITER-Like Wall experiment, beryllium erosion originating from the main chamber is measured by filtered BeII light emission from the JET wide angle camera and emission spectroscopy observing BeII (527 nm) viewing the inner wall. Beryllium erosion has been successfully modelled by ERO2.0 using the full 3D geometry of the inner wall, the SOL flows of the main plasma from EDGE2D-EIRENE and the plasma experimental kinetic profiles. The results have been validated successfully by comparison to the synthetic reconstruction of the camera image [119]. In addition, the spectroscopy viewing the inner wall has also provided material for the code validation for different operational phases: limiter, L-mode, H-mode, including the charge exchange neutral contribution and the drifts. Chemically assisted physical sputtering (releasing Be-D molecules) is a source of uncertainty in this modelling which is not yet fully taken into account because of the lack of sputtering yield data. However, the good agreement with the respective observations confirms the validity of the model and qualitatively confirm the significance of Be migration into the divertor, in particular to the top of the inner baffle (fig 13), which is in line with post-mortem analysis results and provides a firm basis for ITER predictions [112].

In JET-ILW and in other metallic devices, it has been shown that W gross erosion in the divertor averaged over the entire campaign is governed by the intra-ELM phase of the H-mode [120]. Gross tungsten erosion is observed in-situ by the particle flux from the tungsten I line emission and the measurements modelled using ERO [121]. It has been found that in-between ELMs deuterium ions have too low energy for tungsten sputtering and thus only beryllium impurity ions lead to tungsten erosion. In contrast, during intra-ELM phases the sputtering is dominated by deuterium ions by ~70% (over Be sputtering ~30%), which are assumed to have high impact energy due to the ELMs. With tritium, the effect of ELMs on erosion is likely to increase and lead to more tungsten erosion by ELMs (fig 14). Tungsten emission on the divertor plates is also observed using two interference filters of different bandwidths centered on the WI emission on cameras [122]. Thanks to this new technique, the intra-ELM erosion is quantified (7×10^{18} W atoms/ELM) in typical 3MA/2.9T H-mode pulses with 21MW of input power. This shows that inter-ELM tungsten atom fluxes have a strong in/out divertor asymmetry by typically a factor of 18, thus making the outer divertor tiles net W erosion areas. In the case of the inner divertor, net W sputtering is dominated by the intra-ELM erosion. W gross erosion measured by spectroscopy is indeed comparable in both divertor legs in the intra-ELM phase. Spectroscopy amounts to a W gross erosion of 40-60g and post-mortem analysis a W net erosion of 2.4-4.8g, thus corresponding to a re-deposition fraction of more than 90% averaged over both phases and averaged over the JET campaigns [113]. W sputtering fluence per ELM was modelled in deuterium and hydrogen plasma [123].

Depending on the plasma pedestal temperature W sputtering flux is lower in hydrogen than in deuterium plasmas due to the lower Be ion flux to the target because of the lower main chamber Be source and higher sputtering threshold for protons than for deuterons. This is consistent with quartz micro-balance measurements in the divertor showing increased deposit by a factor of 1.5 to 2 in deuterium with respect to hydrogen [124].

4. TECHNICAL AND OPERATIONAL PREPARATION FOR THE DT PHASE

4.1 Key upgrades for the development of DT scenarios.

In the last years, important upgrades have been made on the JET NBI system in preparation for the DT phase. In 2014, following the successful installation of new beam injectors [125], the full energy ion dump (FEID) curved end plate power handling capabilities and fatigue life have been reviewed, redesigned and then replaced in the last shutdown. The new curved end plate design includes a threefold increase in water flow, giving improved cooling performance, additionally high strains caused by the heat flux is relieved by the introduction of slotting, thus improving fatigue life. These new plates give confidence that the design goal of more than 34MW of injected deuterium beam power can be achieved. In addition, for tritium operation, prediction of beam power fractions has been carried out [126] suggesting that more than 2MW of tritium beam power can be achieved with the new injectors, thus giving up to 16MW in deuterium and 17MW in tritium for each beam box. This will therefore provide a significant increase in input power (by ~ 12 MW) with respect to the 1997 DTE1 campaign.

For tritium fueling, a total of five Tritium Introduction Modules (TIMs) have been installed at JET, one on top of the machine, one at the mid-plane and three in the divertor region [127]. The design was heavily inspired by the unique tritium valve used in DTE1. Gas flow rates from the old valve were insufficient to meet the demands for the future TT or DT campaigns, so the new modules incorporate high flow piezo-valves with local tritium storage. Since no human intervention or serviceability to the TIMs is foreseen during the tritium experimental campaigns, their design have incorporated redundancy, reliability, secure operation and conform to JET specific design, inspection, testing and safety case requirements.

New feedback control algorithms have been designed for the ITER-Like-Antenna (ILA) for phase control and second stage matching thus extending the range of the operation to lower (29MHz) and higher (51MHz) frequencies than previously achieved and allowing more flexible and reliable operation. Operation with coupled power levels up to 2.8MW and voltages up to 40kV was achieved in 2016 [128]

For the complete exploitation of the JET ITER-like Wall and to take full benefit from deuterium-tritium experiments on JET, a set of diagnostic system refurbishments and upgrades [129] have been deployed. JET can now rely on a comprehensive set of techniques to measure the neutron yield, neutron spectra and fast particles and is now equipped with vertical and horizontal lines of sight for neutron spectrometry, allowing the separation of RF, NBI and thermal contributions as shown in figure 7b of section 2.4. Various gamma ray spectrometers are expected to provide unique data to codes for simulating the contributions of the trapped and passing fast particles. The redistribution of the alphas and fast ions, by various instabilities, can be measured with the gamma ray cameras. A new scintillator probe and refurbished Faraday cups will be able to relate the fast particle losses to the MHD instabilities. Also, given the importance of ion temperature measurements for DT campaigns, the charge exchange instrumentation has been upgraded to improve the throughput and allow the simultaneous measurement of impurity and fuel-ion charge exchange. Synthetic diagnostics are being added to the most routinely used codes, for easier comparison to diagnostics data and to facilitate the interpretation of plasmas in the coming experimental campaigns, in support of the experiments execution. From a technological perspective, the

planned DT campaign will provide a unique opportunity to test ITER relevant diagnostic technologies, by performing measurements in hostile environments since the expected neutron flux at the first wall (10^{16}n/cm^2) is comparable to that of ITER behind the blanket. [130]

4.2 Real time control in support of DT operation

With the objective of increasing DT scenario robustness and duration, several new plasma control schemes have been designed and will be tested and optimised in deuterium plasmas in 2019.

During tritium and DT campaign both Tritium inventory and neutron budget are severely limited (see below). Hence, real-time detectors have been developed for an early termination of underperforming discharges so as to reduce consumption. Based on normalised confinement and neutron rate indicators, these so-called “dud” detectors can be used to trigger an alarm and a safe plasma termination. [131]. In DT campaigns, the isotope control relies on visible spectroscopy for inferring the isotopic ratio at the plasma edge in real time and specific gas flow control schemes have been designed to manage the injection with the new five tritium gas injection modules and ensure the requested fuel mix (fig 15).

For ensuring that the target of 40MW of input power is sustained for 5s without damaging the first wall, an extensive real-time protection system monitoring the surface temperature of PFCs has been developed [132]. The imaging system comprises four wide-angle views, four tangential divertor views, and two top views of the divertor. In that way, it covers 66% of the first wall and up to 43% of the divertor. This is supervised by a powerful software package analysing the origin of heat loads events and a hot spot management tool allowing the set-up of specific thresholds and assertion times depending on the type of event. Importantly, since D-T operation will cause failure of camera electronics located in the torus hall [133], the images of two key camera systems have been relayed outside of the JET biological shield.

ELM frequency control using pellet injection or gas dosing, has been used together with β_N control for stabilizing the scenario discharges [12]. Plasma protection from radiation peaking to prevent the discharge from collapsing by excessive high Z radiation is also commonly used and detachment control is also being explored in support of scenarios with radiative layers [134].

4.3 Fuel recovery and pumping

For the purpose of fuel recovery and cleaning, reliable isotope ratio measurements are essential for assessing the level of tritium in the vessel. The isotope ratio obtained from Residual Gas Analysis (RGA) signals agrees in general with optical spectroscopy measurement in the divertor [135] and provides an isotopic ratio accuracy of 0.5%. Determination of recycled hydrogen isotopologues released in the W divertor (H_2 , HD, D_2) and from the main chamber wall (the same and Be-D, Be-T, Be-H) allows assessing the local isotope ratio with accuracy below 0.5% [136].

Also, a new strategy for reducing T wall inventory below 1% has been prepared [137]. This procedure will be first validated in hydrogen to reduce the inventory of deuterium below 1% prior to the tritium campaign in order to avoid excessive D-T neutrons in this phase. It also involves one week of vacuum vessel baking at 320°C, combined with isotopic exchange by hydrogen glow discharges and ion cyclotron wall cleaning (ICWC) [138]. This procedure mimics closely the ITER wall cleaning planned strategy. In addition, the baking phase will be followed by plasma pulses in the optimized configurations to access the deposits at the divertor baffles.

The impact of isotope mass and divertor configuration on the divertor conditions and neutral pressures has been addressed by modelling for various isotopes [139]. The results show that a change from hydrogen to deuterium as main fuel decreases the neutral mean free

path, leading to higher neutral density in the divertor. A continuation of this study into tritium is expected to yield a further increase of roll-over densities at detachment when compared to deuterium and hydrogen plasmas. This effect indicates that the isotope ratio control may also be different in H-D and D-T.

4.4 Calibration of 14MeV neutron detectors

For DT operation, calibration of the JET neutron monitors at 14 MeV neutron energy has been performed using a calibrated 14-MeV neutron generator deployed inside the vacuum vessel by the remote handling system. The neutron generator (fig 16) was equipped with two previously calibrated diamond detectors and activation foils [140] which continuously monitored the neutron emission rate during the calibration. The monitoring activation foils were retrieved at the end of each day for decay γ -ray counting, and replaced by fresh ones. About 76 hours of irradiation, in 9 days, were needed to scan the neutron generator through 73 different poloidal and toroidal positions and calibrate both neutron yield monitoring systems available at JET, namely the ^{235}U fission chambers and the activation system using the $^{93}\text{Nb}(n,2n)^{92\text{m}}\text{Nb}$ and $^{27}\text{Al}(n,\alpha)^{24}\text{Na}$ activation reactions.

Neutronics calculations have been performed using MCNP code [141] and a detailed model of JET to derive the response of the neutron detectors to DT plasma neutrons from the measured response to the generator neutrons, also taking into account the anisotropy of the neutron generator the presence of the massive remote handling system and other calibration circumstances. These calculations have used a comprehensive geometrical description of the neutron generator and a validated neutron source routine producing neutron energy-angle distribution for the neutrons emitted by the neutron generator.

The calibration factors for a DT plasma have been determined to within $\pm 6\%$ - 8% total uncertainty for the activation system and to within $\pm 5\%$ for the fission chambers. The difference between the fission chamber response to DD neutrons and that to DT neutrons is within the uncertainties of the derived responses. The same conclusion can be derived for TT neutrons, which are emitted in an almost intermediate energy range, between 1 MeV and 9 MeV. This result has important consequences for operations with DT plasma mixtures with varying concentrations of T: in fact, fission chambers, having a flat response, always measure the total neutron yield (DD + TT + DT) whereas the 14 MeV neutrons can in any case be discriminated by the activation system using the $^{93}\text{Nb}(n,2n)$ reaction which has a threshold at about 10 MeV. The calibration factors of the two independent systems will be compared and validated during full DT operations (producing almost only 14 MeV neutrons) when both fission chambers and activation system will measure independently the same DT yield.

The experience gained and the lessons learnt are particularly valuable for the 14 MeV neutron calibrations in ITER.

4.5 Operation with tritium

There are a number of additional technical requirements (compared to deuterium and hydrogen) to operate JET with tritium gas including high DT neutron flux and neutron activation. Unlike usual gases supplied by gas bottles (e.g. deuterium, hydrogen, neon, argon), tritium is stored in uranium beds and will be supplied by the tritium plant to one or two neutral beam boxes and five new tritium introduction modules (see section 2.1). The torus hall atmosphere will be under depression to limit the spread of tritium in case of accidental tritium release. Access will be restricted to key operational areas of the JET building and tightened access restrictions are also applied to computer networks. During operation, the divertor and neutral beam cryo-pumps will be regenerated after every operational day.

In preparation for the future TT and DTE2 campaigns at JET, an 8-week technical rehearsal of the procedures and systems to be used in tritium operation (without the use of

tritium) has been performed [137]. The rehearsal demonstrated that JET is still capable of carrying out tritium experiments safely. In addition, an extra 3 weeks of plasma operation rehearsal with and without power are planned before tritium is introduced in the machine. Tritium accountancy is an essential part of the tritium operation and the DT safety case limits the tritium on the torus and neutral beam cryo-panels to 11g (44 bar.l). The expected tritium gas is approved 2 weeks before the operation and the gas inventory monitored routinely by dedicated software. Since JET has also a limited 14MeV budget of $2 \cdot 10^{21}$ neutrons, in order to limit the activation of the vessel, the neutron budget is also carefully monitored. Predictions of neutron activation have been validated by past periodic measurements of the dose rate inside the JET vessel and by sophisticated shutdown dose rate codes [138].

5. OUTLOOK

The active preparation of JET for the DT campaign provides an incomparable source of information for the physics basis of DT plasmas and the future operation of ITER. JET experimental and analysis campaigns have strengthened the preparation of the next tritium and D-T campaigns. Both baseline and hybrid operational regimes have produced promising results (more than 7MW of equivalent fusion power), while remaining compatible with the ILW. In support of the scenario a huge step forward in the understanding of the isotope physics has been achieved. This would not have been the case without a strong focus towards DT. In addition, the operational preparation for the DT phase has now reached a high quality level, tackling issues that had not been addressed in non-tritium compatible device such as neutron calibration, tritium handling, isotope control, wall erosion, etc.

JET is uniquely placed to provide a robust base for burning plasmas during its future DT phase. The shutdown required to prepare for deuterium-tritium operations is now complete and despite recent delays, JET is ready to move towards a new domain with the new DT phase. The main campaign elements until the end of 2020 are:

- 1- A deuterium campaign for the preparation of high performance scenarios and studies of disruption and run-away election mitigation using a new shatter-pellet injection system presently installed on JET as part of an ITER, EURATOM an US-DOE agreement;
- 2- Hydrogen and 100% tritium campaigns combined with reference pulses in deuterium for the study of isotope effects;
- 3- A deuterium-tritium campaign with the aim of producing 15 MW of fusion power for 5s in stationary conditions in order to study the key physics aspects, such as those related to alpha particles, that a change from deuterium to a deuterium-tritium plasma may involve.

The JET programme, will provide invaluable information for the non-activated phase of ITER, alpha particle and isotopic effects in plasma scenarios, particle and heat transport, retention, and wall cleaning. In parallel to the preparation of the DT phase, JET is developing a strong program in support of the ITER research plan for disruption mitigation with the new shattered pellet injector and radiative layer studies. All these efforts will ensure a maximized output to ITER operation from its first plasma to the DT phase.

ACKNOWLEDGEMENTS

This work has been carried out within the framework of the EUROfusion Consortium and has received funding from the Euratom research and training programme 2014-2018 and 2019-2020 under grant agreement No 633053. The views and opinions expressed herein do not necessarily reflect those of the European Commission.

REFERENCES

- [1]: J. Pamela, et al., Fusion Engineering and Design 82 (2007) 590–602
- [2]: M. Keilhacker Nuc Fus 39 (1999) 209
- [3]: J. Jacquinot Nuc Fus 39 (1999) 235
- [4]: G.F. Matthews et al Phys. Scr. T145 (2011) 014001;
- [5]: L. Horton et al, Fusion Engineering and Design 109–111 (2016) 925;
- [6]: X. Litaudon et al, Nucl. Fusion 57 (2017) 102001.
- [7]: <https://www.euro-fusion.org/eurofusion/roadmap>
- [8]: L. Horton et al., 12th International Symposium on Fusion Nuclear Technology (ISFNT), Jeju Island, Korea
- [9]: H. Weisen et al., AIP Conference Proceedings 1612, 77 (2014); doi 10.1063/1.489428
- [10]: ITER Research Plan, 2018 <https://www.iter.org/newsline/-/3121>
- [11]: I. Nunes et al., Plasma Phys. Control. Fusion 58 014034 (2016)
- [12]: L. Garzotti, et al, 27th Fusion energy conference, 22-27th October 2018 Gandhinagar, India, IAEA-CN-EX/3-6 to be published in Nuclear Fusion
- [13]: A.C.C Sips accepted in Nucl Fus 2018
- [14]: C.D. Challis et al 2015 Nucl. Fusion 55 053031
- [15]: TRANSP group in Princeton Plasma Physics Laboratory TRANSP Home page <http://w3.pppl.gov/transp/>
- [16]: R. Budny et al, Nuc. Fus 32 (1992) 429
- [17]: H. T. Kim et al., Nuclear Fusion, 58 (2018) 036020
- [18]: M. Mantsinen et al., 44th European Physical Society Conference on Plasma Physics (EPS), Belfast, Northern Ireland, 26th June 2017, O3.110
- [19]: S. Breton et al 2018 Nucl. Fusion 58 096003
- [20]: F. Casson, 27th Fusion energy conference, 22-27th October 2018 Gandhinagar, India, IAEA-CN-TH/3-2 to be published in Nuclear Fusion.
- [21]: T.C. Hender et al 2016 Nucl. Fusion 56 066002
- [22]: J Garcia et al 2017 Plasma Phys. Control. Fusion 59 014023
- [23]: J. Garcia et al 2015 Nucl. Fusion 55 053007;
- [24]: Citrin J. et al 2013 Phys. Rev. Lett. 111 155001;
- [25]: S A Silburn et al 2017 Phys. Scr. 2017 014040
- [26]: C. D. Challis, et al., 44th EPS Conference on Plasma Physics, 2017, P2.153
- [27]: C. Giroud, 27th Fusion energy conference, 22-27th October 2018 Gandhinagar, India, IAEA-CN-EX/3-4 to be published in Nuclear Fusion.
- [28]: R. Dumont, et al., 2018 Nucl. Fusion 58 082005
- [29]: S. Sharapov et al. FST 52(4) 989-1022
- [30]: R. Nazikian, Phys. Rev. Lett., 78:2976-2979, Apr 1997
- [31]: A. B. Mikhailovskii et al. Plasma Phys. Rep., 23:844, 1997
- [32]: Pinches, (1998). Computer Physics Communications, 111(1–3), 133–149
- [33]: P. Puglia et al., Nucl. Fusion 56, 112020 (2016)
- [34]: S. Sharapov et al., 27th Fusion energy conference, 22-27th October 2018 Gandhinagar, India, IAEA-CN-EX/P1-28 to be published in Nuclear Fusion
- [35]: M. Fitzgerald et al. to be published in Nuc Fus 2018
- [36]: J. Hobirk et al 2018 Nucl. Fusion 58 076027
- [37]: P. de Vries 2018 Nucl. Fusion 58 026019
- [38]: E. de la Luna et al., 27th Fusion energy conference, 22-27th October 2018 Gandhinagar, India, IAEA-CN-EX/2-1 to be published in Nuclear Fusion.
- [39]: F. Koechl et al., 2017 Nucl. Fusion 57 086023
- [40]: C. Sozzi et al, 27th Fusion energy conference, 22-27th October 2018 Gandhinagar, India, IAEA-CN-EX/P1-22 to be published in Nuclear Fusion;

- [41]: Teplukhina A. et al 2017 Plasma Phys. Control. Fusion 59 124004
- [42]: M. Staebler et al., 2005 Phys. Plasmas 12 102508
- [43]: C Bourdelle et al 2016 Plasma Phys. Control. Fusion 58 01403
- [44]: Romanelli M et al 2014 Plasma Fusion Res. 9 3403023
- [45]: J.F. Artaud et al 2010 Nucl. Fusion 50 043001
- [46]: J. Garcia et al., 27th Fusion energy conference, 22-27th October 2018 Ghandinagar, India, IAEA-CN-TH/3-1 to be published in Nuclear Fusion.
- [47]: Cordey J G 2003 Nucl. Fusion 43 670–4
- [48]: J. Garcia et al 2017 Nucl. Fusion 57 014007
- [49]: Honda M. et al 2006 Nucl. Fusion 46 580
- [50]: S Saarelma et al 2018 Plasma Phys. Control. Fusion 60 014042
- [51]: F. Casson, et al., 27th Fusion energy conference, 22-27th October 2018 Ghandinagar, India, IAEA-CN-TH/3-2 to be published in Nuclear Fusion.
- [52]: P. Strand et al., 27th Fusion energy conference, 22-27th October 2018 Ghandinagar, India, IAEA-CN-TH/P6-14 to be published in Nuclear Fusion.
- [53]: S. Nowak et al., 27th Fusion energy conference, 22-27th October 2018 Ghandinagar, India, IAEA-CN-TH/P6-26 to be published in Nuclear Fusion.
- [54, 55, 56]: C. Angioni et al., Nucl. Fusion 57 056015 (2017) ;
- [55]: E. Lerche et al., Nuclear Fusion 56 (2016) 3 036022;
- [56]: M. Goniche et al., Plasma Phys. Control. Fusion 59 (2017) 055001
- [57]: P. Jacquet et al., Nucl. Fusion 56 046001 (2016)
- [58]: F. Durodie, Fus. Eng. Des. 84 (2009) 279–283
- [59]: L. Colas et al. 45th European Physical Society Conference on Plasma Physics (EPS), Prague Czech Republic, 2nd July 2018, O4.101;
- [60]: V. Bobkov, PSI 2018 submitted to Nuclear Materials and Energy
- [61]: L. Colas et al., Phys. Plasmas 19 092505 (2012)
- [62]: D. Gallart et al 2018 Nucl. Fusion 58 106037
- [63]: K. Kirov et al., ‘Studies on the fast ions synergistic effects in JET high performance pulses’, submitted to Nucl. Fusion
- [64]: D. Van Eester et al., Proc. 26th IAEA-FEC Conference, Kyoto, EX/P6-10
- [65]: D.F.H. Start et al 1999 Nucl. Fusion 39 321
- [66]: J. Ongena et al., EPJ Web Conf. 157, 02006 (2017)
- [67]: Ye. Kazakov et al., 27th Fusion energy conference, 22-27th October 2018 Ghandinagar, India, IAEA-CN- to be published in nuclear fusion.
- [68]: Ye.O. Kazakov et al., Nature Physics 13, 973-978 (2017)
- [69]: Ye.O. Kazakov et al., Phys. Plasmas 22, 082511 (2015)
- [70]: C F Maggi et al 2018 Plasma Phys. Control. Fusion 60 014045
- [71]: J. Hillesheim, et al., 27th Fusion energy conference, 22-27th October 2018 Ghandinagar, India, IAEA-CN-EX/4-1 to be published in Nuclear Fusion.
- [72]: D. Moulton et al 2018 Nucl. Fusion 58 09602
- [73]: Y. Martin et al., Conference Series 123 (2008) 012033
- [74]: E Righi et al 1998 Plasma Phys. Control. Fusion 40 721
- [75]: P. Gohil et al., Nucl. Fusion 50, 064011 (2010)
- [76]: F. Ryter et al 2016 Plasma Phys. Control. Fusion 58 014007
- [77]: E. Solano, 23rd Joint EU-US Transport Task Force Meeting, Sevilla 2018
- [78]: A.H. Nielsen et al., Phys. Lett. A 79 (2015) 3097;
- [79]: J. J. Rassmussen, et al., 23st EU-US Transport Task Force Meeting, Seville, Spain, 11-14 September 2018
- [80]: C. F. Maggi et al, Plasma Physics and Controlled Fusion 60 (2018) 014045
- [81]: M. Romanelli, et al., Plasma and Fusion Research 9 (2014) 3403023],
- [82]: J. Garcia Physics of Plasmas 25, 055902 (2018);

- [83]: N. Bonanomi et al 2018 Nucl. Fusion 58 056025;
- [84]: A. Di Siena et al., EU-US TTF, Sevilla, 10th Sept 2018
- [85]: N. Bonanomi et al 2018 Nucl. Fusion 58 026028
- [86]: H. Urano et al, Nuclear Fusion 52 (2012) 114021
- [87]: Huber A et al. 2013 Journal Nucl Mater. 438 S139
- [88]: H. Weisen et al, 27th Fusion energy conference, 22-27th October 2018 Ghandinagar, India, IAEA-CN-EX/P1-4 to be published in Nuclear Fusion. [89]: ITER Physics Basis, Nuclear Fusion 39 (1999) 2175
- [90]: M. Maslov et al., Nucl. Fusion 58 (2018) 076022
- [91]: K-D Zastrow et al., Plasma Phys. Control. Fusion 46 B255 (2004)
- [92]: P. C. Efthimion et al., Phys. Rev. Lett. 75 1 (1995) ;
- [93] : M. R. Wade et al., Phys. Plasmas, 2 2357 (1995)
- [94]: C. Bourdelle, Nucl. Fusion 58 (2018) 076029
- [95]: J. Citrin et al., Plasma Phys. Control. Fusion 59 124005 (2017);
- [96] : M. Marin et al., EPS Prague 2018, O2.10
- [97]: D. Frigione, 27th Fusion energy conference, 22-27th October 2018 Ghandinagar, India, IAEA-CN-EX/P1-3 to be published in Nuclear Fusion.
- [98]: A. Kappatou, 27th Fusion energy conference, 22-27th October 2018 Ghandinagar, India, IAEA-CN-EX/P8-1 to be published in Nuclear Fusion.
- [99]: T. Tala et al., 27th Fusion energy conference, 22-27th October 2018 Ghandinagar, India, IAEA-CN-EX/4-4 to be published in Nuclear Fusion.
- [100]: L. Frassinetti et al., Plasma Phys. Control. Fusion 59 014014 (2017)
- [101]: M. Dunne et al., Plasma Phys. Control. Fusion 59 (2017) 025010
- [102]: E. Stefanikova et al 2018 Nucl. Fusion 58 056010
- [103]: S. Saarelma et al., Phys. Plasmas 22 056115 (2015)
- [104]: S. Saarelma et al., accepted to Plasma Phys. Control. Fusion, 2018
- [105]: L. Frassinetti et al., 27th Fusion energy conference, 22-27th October 2018 Ghandinagar, India, IAEA-CN-EX/P8-22 to be published in Nuclear Fusion.
- [106]: L. Horvath et al., to be published in Nucl. Fusion 2018
- [107]: M.T. Kotschenreuter et al., 27th Fusion energy conference, 22-27th October 2018 Ghandinagar, India, IAEA-CN-TH/2-2 to be published in Nuclear Fusion.
- [108]: L. Horvath et al., PPCF (60) 2018 086023
- [109] : D. Hatch et al., to be submitted to Nuclear fusion.
- [110]: S. Brezinsek et al., Journal of Nuclear Materials 463 (2015) 11–21
- [111]: W. Eckstein, Top. Appl. Phys. 110 (2007), 33–187
- [112]: D. Borodin et al., 27th Fusion energy conference, 22-27th October 2018 Ghandinagar, India, IAEA-CN-EX/P1-14 to be published in Nuclear Fusion.
- [113]: S. Brezinsek et al., 27th Fusion energy conference, 22-27th October 2018 Ghandinagar, India, IAEA-CN-EX/9-4 to be published in Nuclear Fusion.
- [114]: K. Heinola et al., Phys. Scr. T170 (2017) 014063
- [115]: E. Pajuste et al., submitted to Journal of Nuclear Material, 2018
- [116]: N. Catarino, et al., Nucl. Mater. Energy. 12 (2017) 559–563
- [117]: A. Widdowson, submitted to Journal of Nuclear Material, 2018
- [118]: M. Mayer et al., 27th Fusion energy conference, 22-27th October 2018 Ghandinagar, India, IAEA-CN-TH/P1-15 to be published in Nuclear Fusion.
- [119]: J. Romazanov et al., Phys. Scr. T170, 014018 (2017)
- [120]: N. Den Harder et al. Nucl. Fusion 56 (2016) 026014
- [121]: A. Kirschner et al., submitted to Journal of Nuclear Material, 2018
- [122]: A. Huber submitted to Journal of Nuclear Material, 2018
- [123]: I. Borodkina et al. Phys. Scr. T170 (2017) 014065 (6pp)
- [124]: G. Sergienko et al., 16th International Conference on Plasma-Facing Materials and

Components for Fusion Applications, 2018, Neuss/Duesseldorf, Germany

[125]: D. Ciric et al., Fusion Engineering and Design, Volume 86, Issues 6–8, October 2011, Pages 509-512

[126]: R.F. King, et al., Fusion Engineering Design 83 (2008) 1553-1558

[127]: I. Carvalho et al., Fusion Engineering and Design, Volume 124, November 2017, Pages 841-845

[128]: P. Dumortier et al., 22nd Topical Conference on Radio Frequency Power in Plasmas, 2017 Aix-en-Provence, France

[129]: J. Figueiredo, et al., 27th Fusion energy conference, 22-27th October 2018 Ghandinagar, India, IAEA-CN-FIP/P7-42 to be published in Nuclear Fusion.

[130]: A. Murari et al., International Conference on Research and Application of Plasmas, Warsaw, Poland, September 18-22, 2017

[131]: L. Piron et al., 30th Symposium on Fusion Technology (SOFT 2018), Sicilia, Italy, 17th September, 2018

[132]: A. Huber et al 2018 Nucl. Fusion 58 106021

[133]: V. Huber, et al., 27th Fusion energy conference, 22-27th October 2018 Ghandinagar, India, IAEA-CN-TH/P1-31 to be published in Nuclear Fusion.

[134]: C. Guillemaut et al., 2017 Plasma Phys. Control. Fusion 59 045001

[135]: A. Drenik, 2017 Phys. Scr. 2017 014021

[136]: C. Klepper, et al., 27th Fusion energy conference, 22-27th October 2018 Ghandinagar, India, IAEA-CN-EX/P1-29 to be published in Nuclear Fusion.

[137]: I. Borodkina et al., 45th EPS Conference on Plasma Physics O2.106, Prague, 2018

[138]: T. Wauters et al., Journal of Nuclear Materials 463 (2015) 1104–1108

[139]: J. Uljanovs et al. Nuclear Materials and Energy 12 (2017) 791–797

[140]: P. Batistoni et al 2018 Nucl. Fusion 58 106016

[141] : P. Batistoni et al 2018 Nucl. Fusion 58 026012

[141]: Goorley J.T. et al 2013 Initial MCNP6 Release Overview MCNP6 version 1.0 Report LA-UR-13-22934 (Los Alamos, NM: Los Alamos National Laboratory)

[142]: E. Belonohy et al. 29th Symposium on Fusion Technology (SOFT 2016), Prague, Czech Republic, 5th September 2016.

[143]: R. Villari et al., 29th Symposium on Fusion Technology (SOFT 2016), Prague, Czech Republic, 5th September 2016.

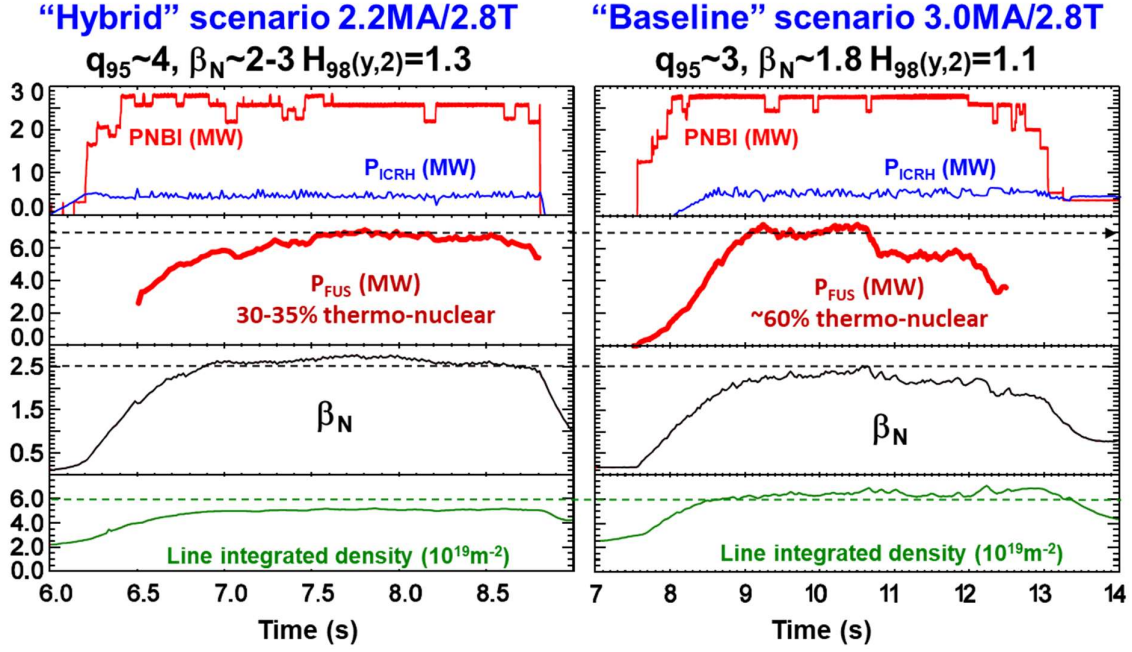


Figure 1: Time traces of the two main scenarios for achieving the target of 15MW of fusion power for 5s [12]. In the second box from the top, the equivalent fusion power is calculated as explained in the text using TRANSP. Note the difference in density β_N , density and thermo-nuclear component between the two discharges. The beam-beam neutron component is negligible in both cases (<few %).

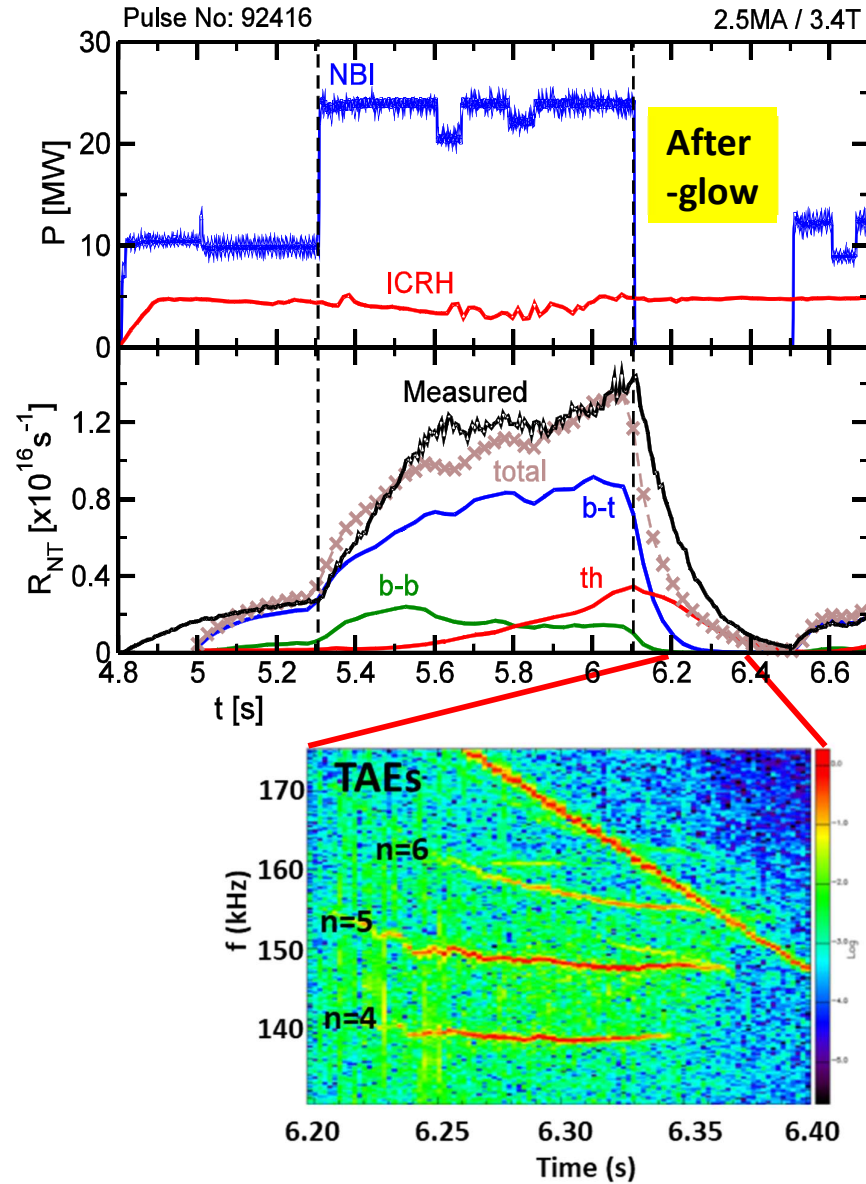


Figure 2a (top): Time traces of the “after-glow” scenario for alpha particle studies [28] featuring an abrupt interruption of the neutral beam injection. Box (b) shows the measured neutron rate (black line), and the TRANSP computed neutron rate components: beam-beam (b-b), thermal (th), beam-target (bt) and total neutron rate.

Figure 2b (bottom): spectrogram of TAE modes measured by the magnetics in the after-glow phase. In DD, the ICRH power has been kept to “illuminate” the fast particle driven modes. In DT, the presence of alpha particle is predicted to produce the same type of modes (see figure 2c and the text).

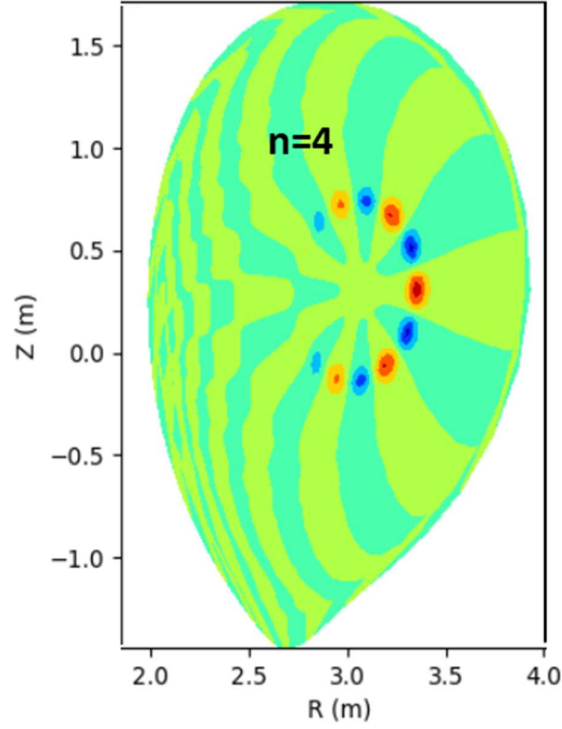


Figure 3: Toroidal Alfvén Eigenmodes stability (TAE) calculation with the MISHKA code applied on the DT prediction of the pulse of figure 2a. Core TAE with toroidal mode number $n=4$ (here represented), 5 and 6 have been found matching those observed with the DD pulse with ICRH (figure 2b).

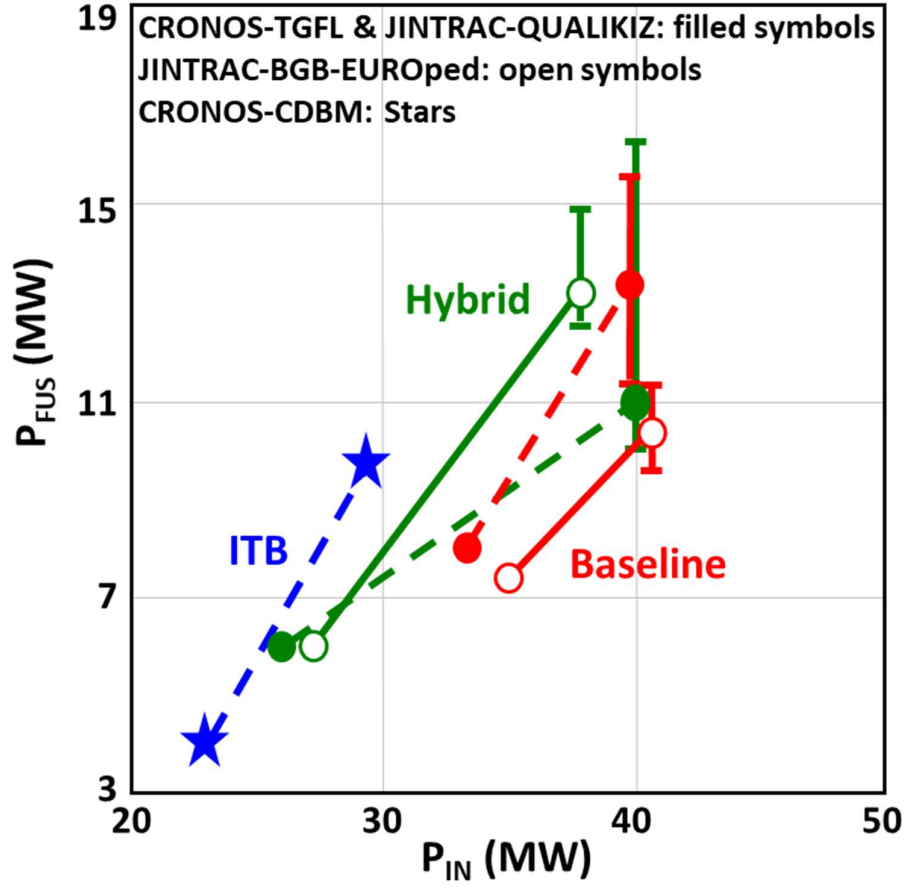


Figure 4: DT fusion power predictions from the discharges shown in figure 1 and 2a to maximum input power (32MW of neutral beam injection and 8MW with ICRH) with the JINTRAC and CRONOS codes. Error bars account for the different bootstrap current models total currents and isotope effects as described in the text (see also [46]).

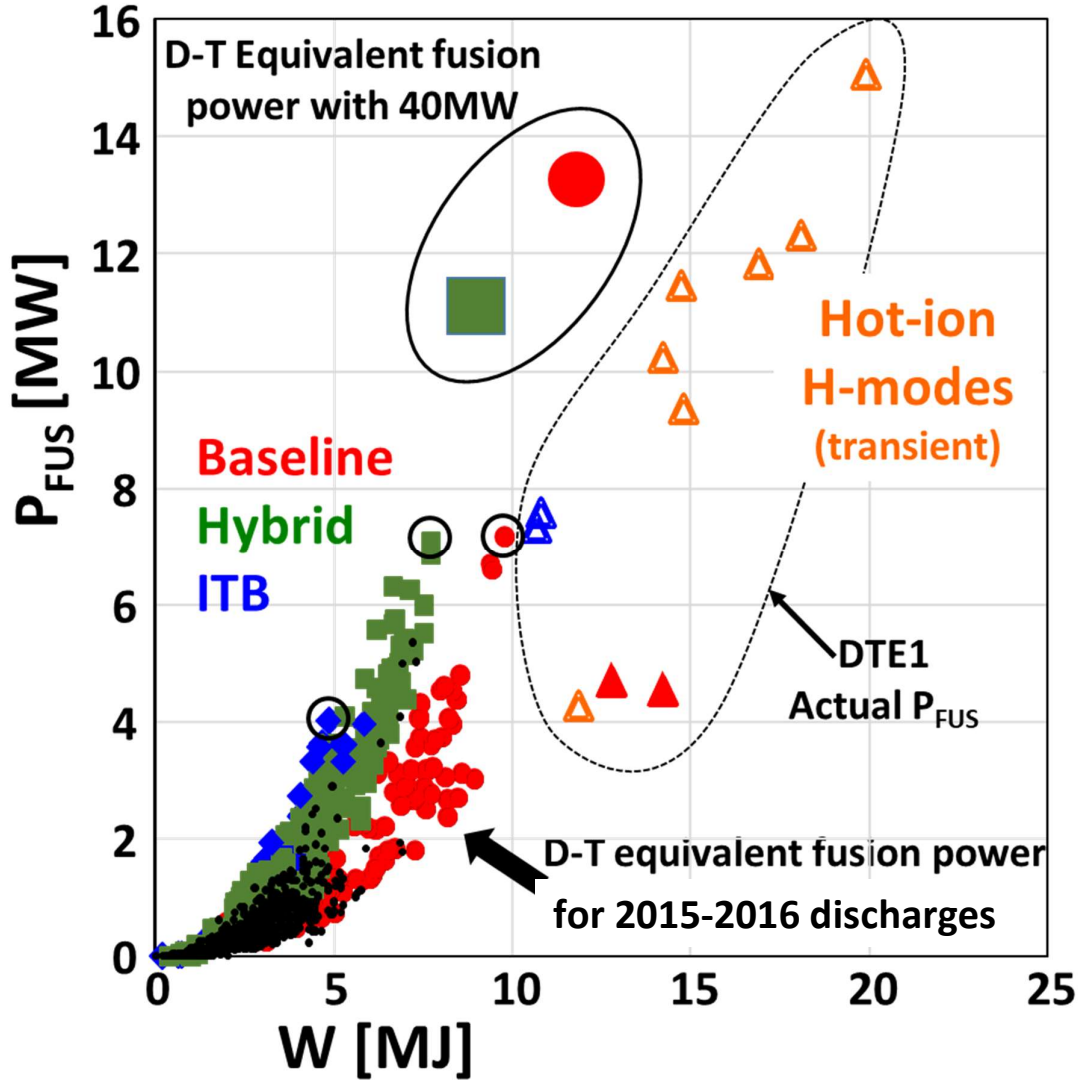


Figure 5: Equivalent fusion power predictions using the fusion power simulator (see text) for the database of DD pulses run with the ITER-like wall overlaid with the actual fusion power achieved in DTE1 (triangles encircled by the dotted line). The colors are red for the baseline stationary scenario, green for the stationary hybrid scenario and blue for the ITB scenario. The orange triangle are representing the transient hot-ion H-mode achieved in DTE1 [2] which is no more part of the reference scenario for DTE2. The large green square and red circle are equivalent fusion power computed from the DD predicitive runs with full power (40MW total) for both hybrid (CRONOS+TGLF see [22]) and baseline (JETTO+QUALIKIZ see [55]) scenarios of figure 1. The scenarios of figure 1 and 2a are also identified by the 3 black small circles on this figure.

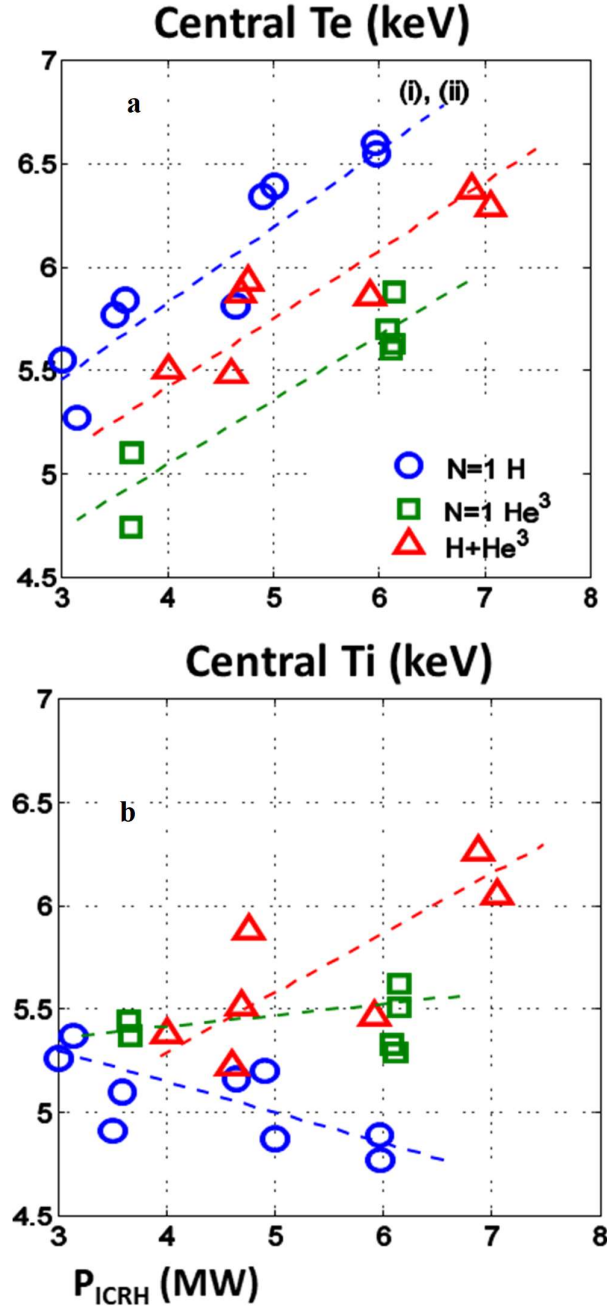


Figure 6: Comparison the effect on plasma electron and ion temperature of the different ICRH schemes N=1 H minority and N=1 ³He minority and combined: Te (top picture) and Ti (bottom picture) from charge exchange spectroscopy. Note that N=1 ³He minority is a more efficient scheme for ion heating depending on the He3 concentration used in the plasma [57].

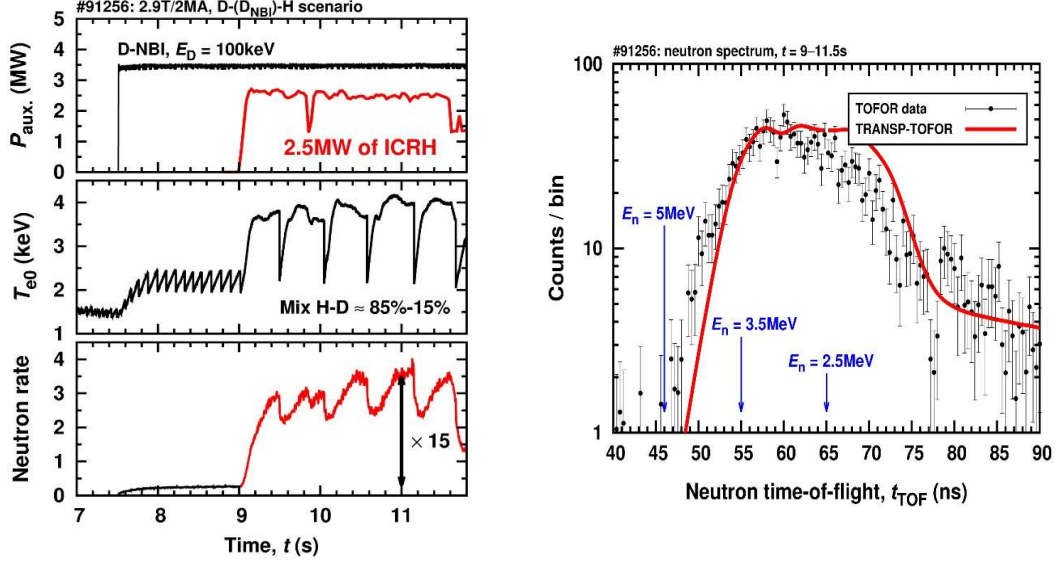


Figure 7a (left): Performance of the three-ion scheme D-(DNBI)-H scenario on JET-ILW. (for $H/D \sim 6$, $n_{eo} = 4 \times 10^{19} \text{ m}^{-3}$, $f_{ICRH} = 25 \text{ MHz}$, dipole phasing). Note the sharp increase of neutron during the ICRH pulse.

Figure 7b(right): For the pulse of figure 6a, comparison of the TRANSP synthetic neutron spectroscopy time-of-flight diagnostics with the actual data of the diagnostic showing good agreement between the two [67].

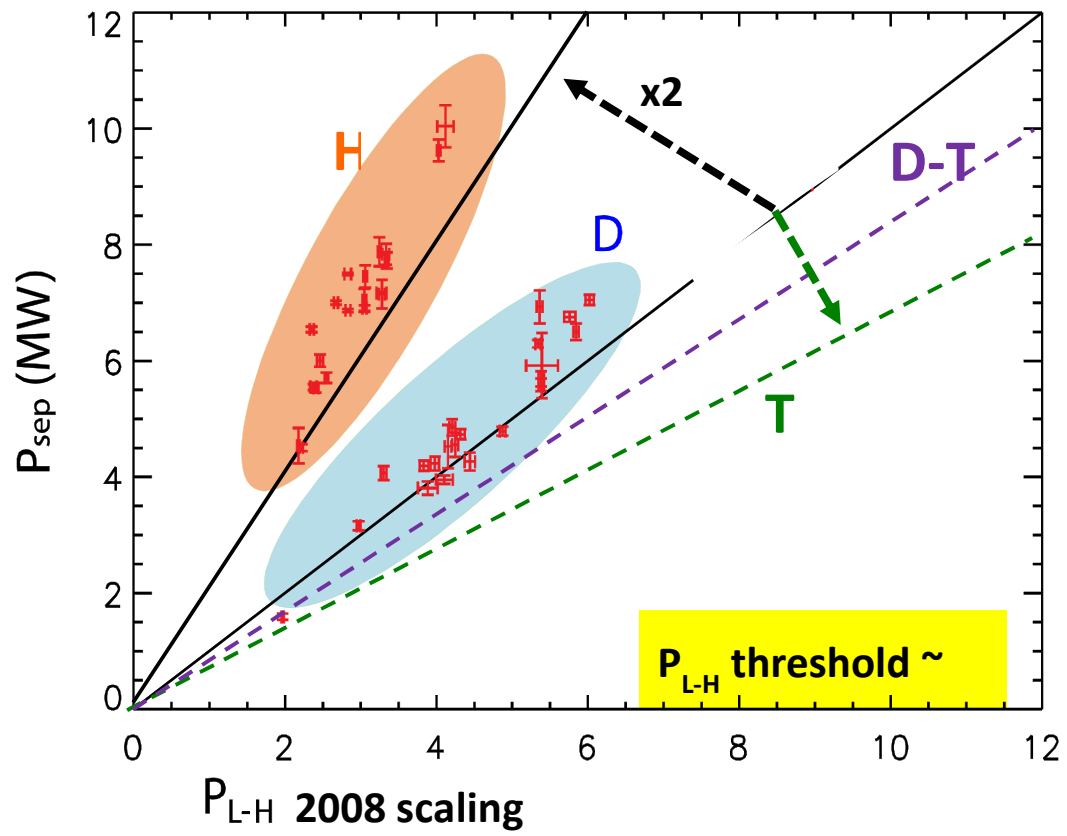


Figure 8: L-H Power threshold for hydrogen and deuterium against the scaling of the power threshold [73] for vertical and corner divertor geometries. The L-H threshold in T-T is expected to decrease by 2/3 and by 1/3 in D-T.

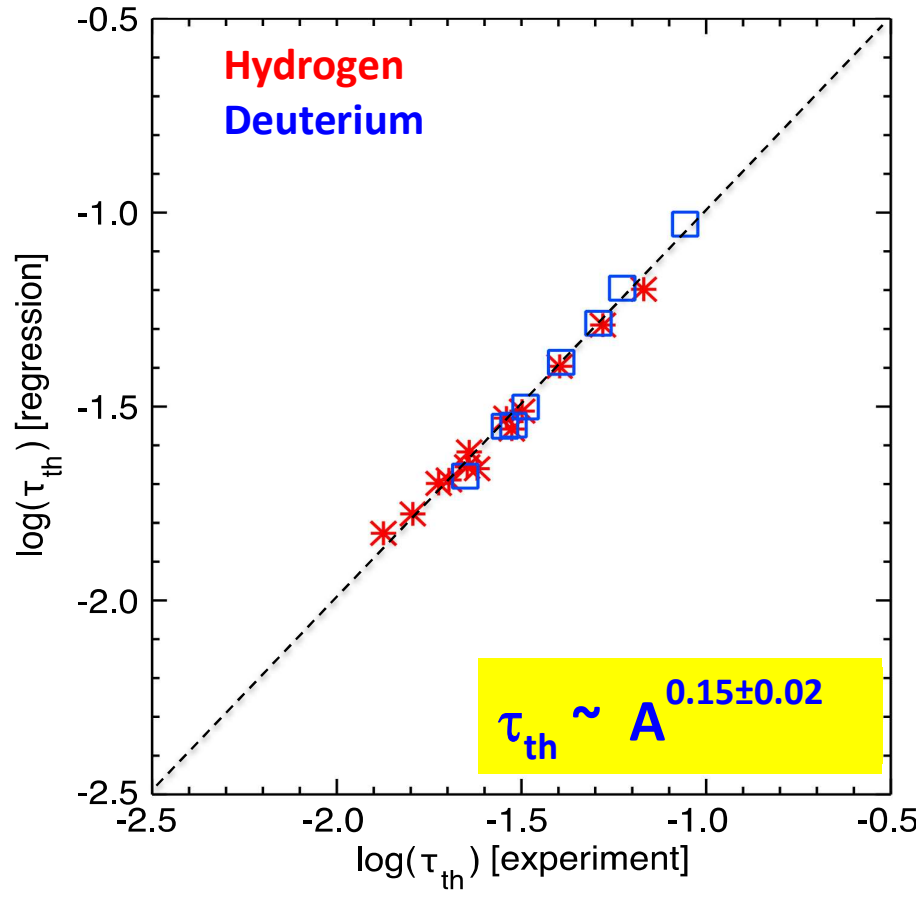


Figure 9: thermal confinement regression with the isotope mass using deuterium and hydrogen discharges in L-mode at constant plasma average density, plasma current of 2.5MA and toroidal field strength of 2.9T [70].

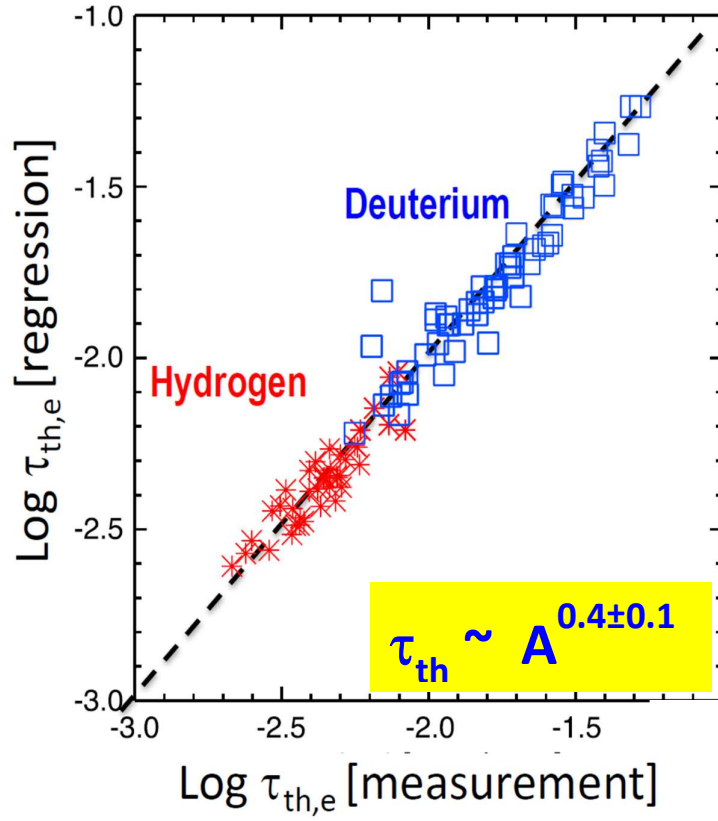


Figure 10: Thermal confinement regression with the isotope mass using deuterium and hydrogen discharges in H-mode [88] with plasma current and toroidal field strength of 1.0MA/1.0T and 1.4MA/1.7T in hydrogen and 1.0MA/1.0T, 1.4MA/1.7T and 1.7MA/1.7T in deuterium.

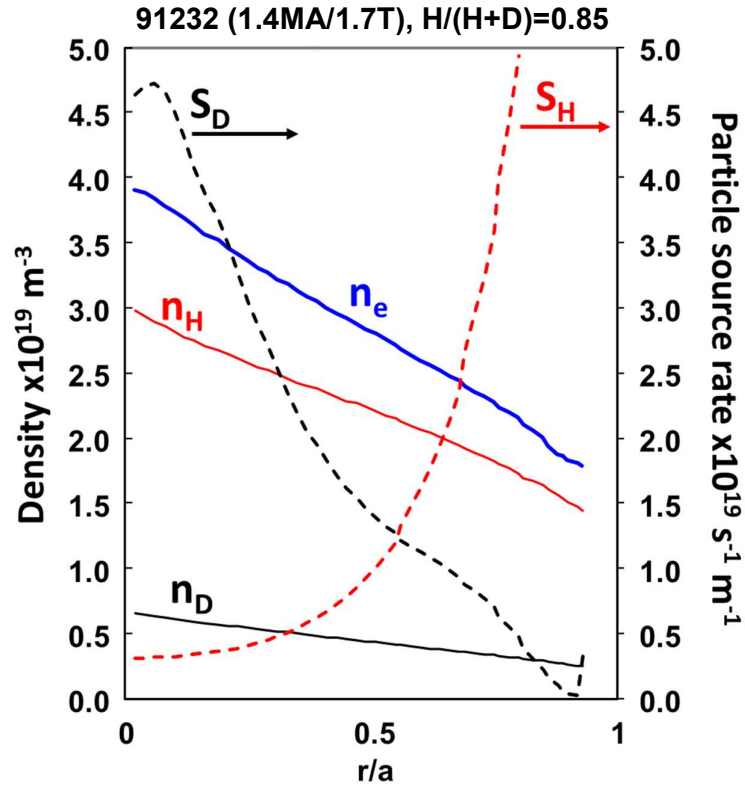


Figure 11: Comparison of electron hydrogen and deuterium density profiles in a mixed plasma with $H/(H+D)=0.85$ and different particle fuel sources (deuterium core fueling by neutral beam and hydrogen edge fueling by gas dosing). Note that the normalized gradients for both deuterium and hydrogen are comparable: $R/L_{nH} \sim R/L_{nD} \sim 2.5$ [90].

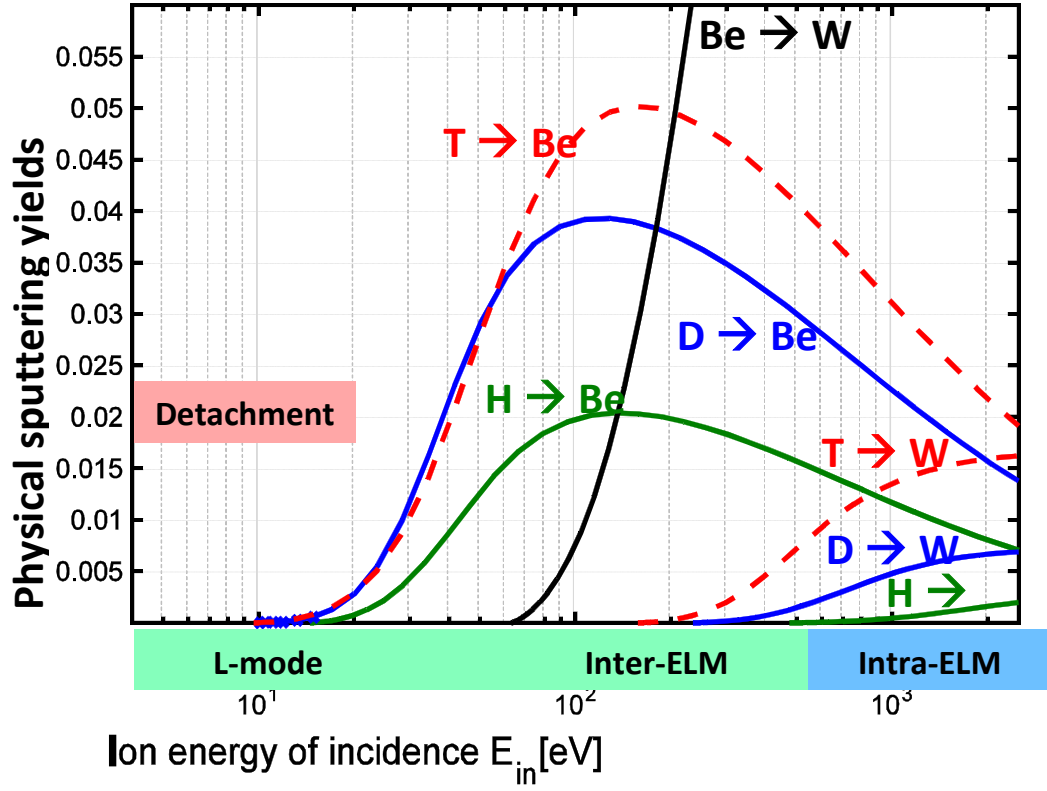


Figure 12: Isotope (H: green; D:blue; T: dotted red line) effect on Be and tungsten sputtering yields computed by the code SDTrimSP [111] for various phases of plasma scenarios. Note the higher tungsten sputtering by tritium in the intra-ELM phase as well as the increased Be sputtering by tritium leading to an increased tungsten sputtering by Be (in black).

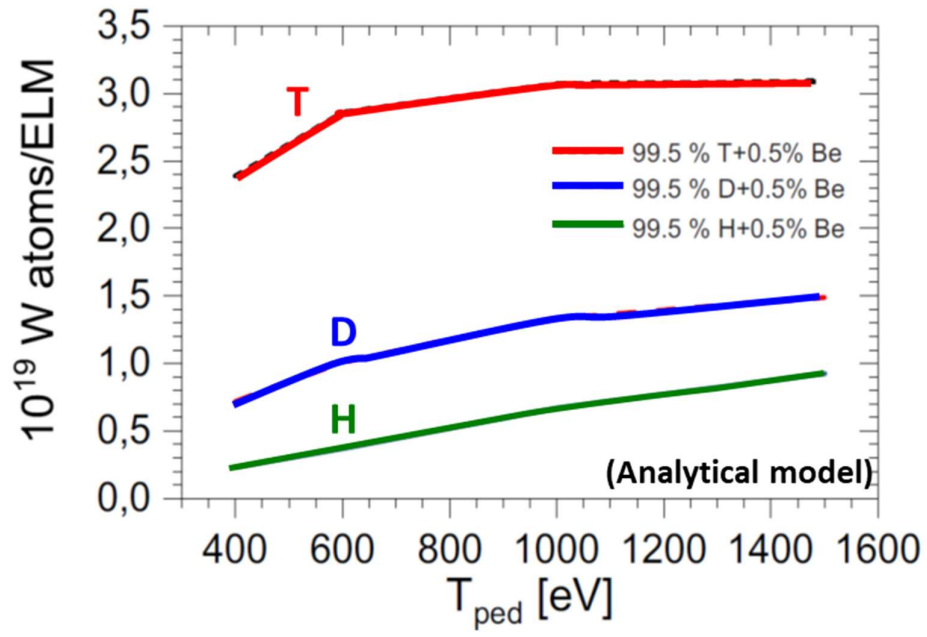


Figure 14: calculation of the amount of tungsten atoms per ELM for each isotope assuming free streaming transport during the ELM crash and typical plasma conditions with 0.5% of Beryllium in the plasma with $B_t = 3.0T$ and $n_{ePED} = 1 \cdot 10^{20}m^{-3}$

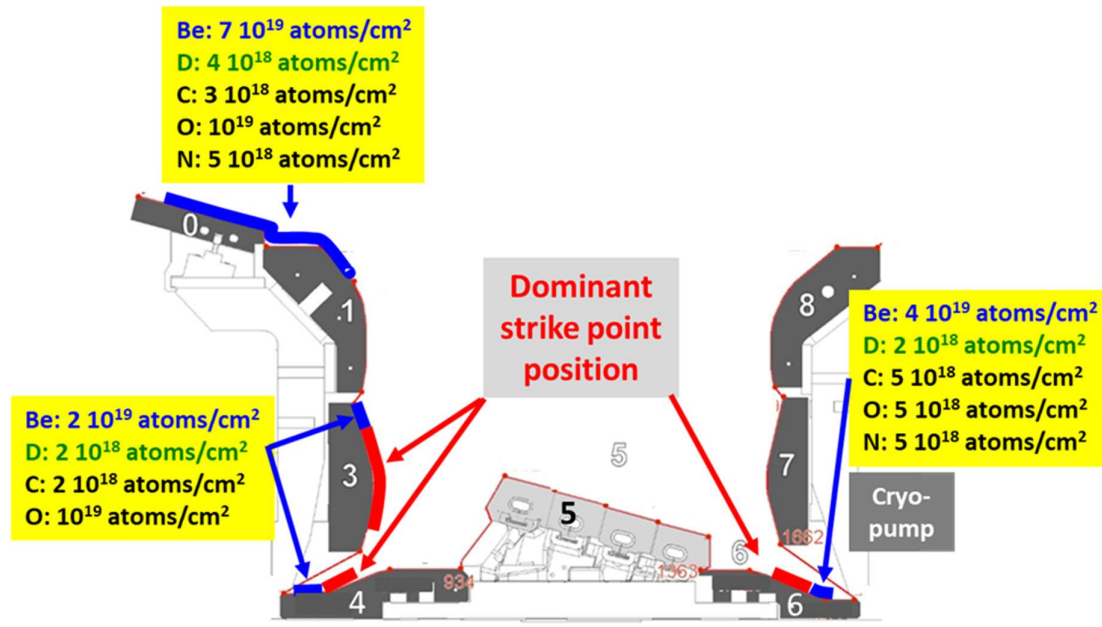


Figure 13: material deposition patterns in the divertor after three JET campaigns [118] with the ITER-like wall and their relation with the dominant strike point locations. Note that the maximum deposition occurs on the inboard side on the top of the divertor baffle (tile 0 and 1).

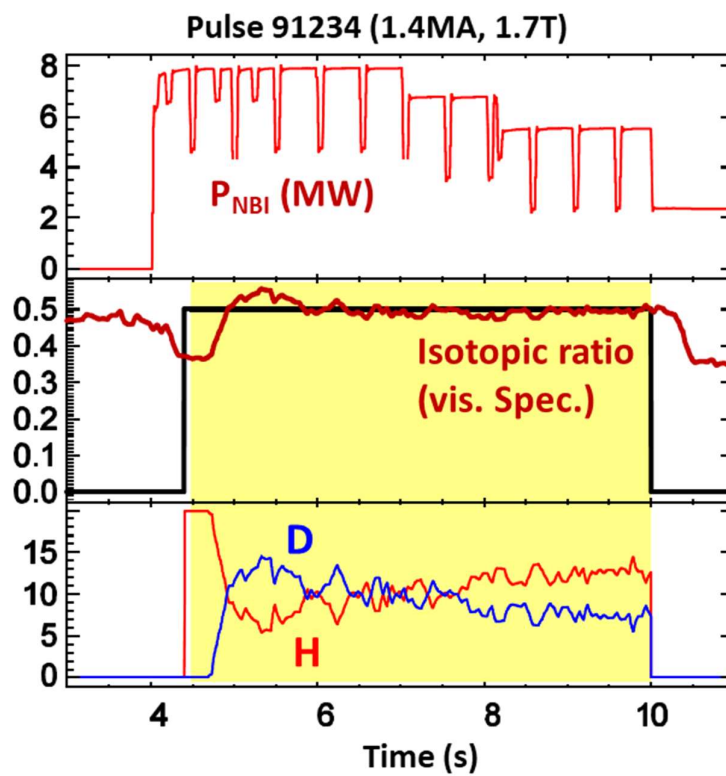


Figure 15: typical example the real time control of the isotope ratio at $H/(H+D)=0.5$ using edge gas dosing of the H and D isotopes.

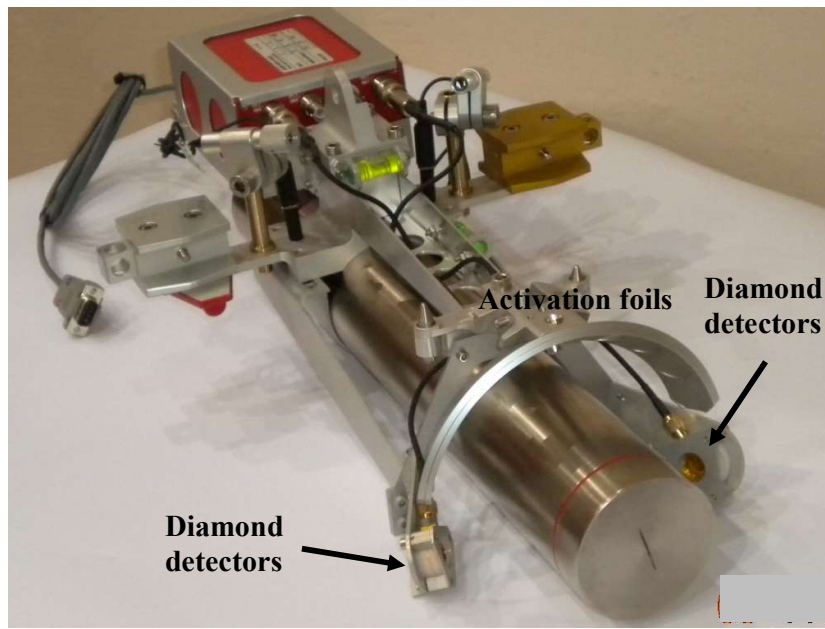


Figure 16: Neutron generator deployed inside the JET vacuum vessel on the remote handling arm. The two single crystal diamond detector and the activation foils are symmetrically hosted in the support around the neutron generator, at the same distance from the target centre [141].