# Deuterium Implantation in Actively Cooled Beryllium Monoblocks 

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## 1. INTRODUCTION

The amount of tritium buried in the first wall of a fusion reactor could turn out to be the decisive factor for the choice of wall material. In the case of copper, the common material for beam dumps, we have found a large discrepancy between data published from surface analysis groups and data from neutron generators on the amount of implanted hydrogen [1]. In neutron generators with high particle energies and fluxes, trapping appears to be dominant and implanted hydrogen concentrations of up to $40 \mathrm{at} \%$ are reported [2,3]. Well controlled low flux experiments done for surface analysis yield concentrations below 1 at $\%$ and the amount of implanted hydrogen is dominated by diffusion and surface recombination [4]. To investigate if similar discrepancies are found in Beryllium we analysed a test section with Nuclear Reaction Analysis after an exposure test with Deuterium beams. In this exposure surface temperature and power density are representative of a first wall component. However the particle energy is higher than for first wall components and the flux is correspondingly lower.

The test section is an actively cooled Beryllium Monoblock (Brush Wellman grade S65C) which had been used before for several power handling tests. Part of the surface had been above liquidus in these tests with a melt depth of less than 1 mm . Aim of this experiment was to determine the quantity if implanted deuterium with Nuclear Reaction Analysis. The experiment was done in two steps:

1. The test section was exposed to a power density of 4-5 MW/m ${ }^{2}$ and analysed some days after exposure.
2. The test section was exposed to high power densities of $20 \mathrm{MW} / \mathrm{m}^{2}$ which brings the surface into melting and was analysed after exposure.

## 2. EXPERIMENTAL SETUP

Fluxes, energies and temperatures of this section refer to the first part of the experiment in which the surface temperature of the beryllium was well below melting.

The experimental setup is the same as in the previous experiment [5]: The test section consists of 20 Beryllium monoblocks brazed onto a 12 mm id cooling pipe (Fig.1). A disc assembly consisting of two silicon wafers and one glass disc was installed above the test section at a distance of approximately 230 mm measured along the line connecting the centers of test section and disc holder. This line is at


Fig.1: Schematic of the test section: Beryllium monoblocks are brazed to copper cooling pipe. The full length of the test section is exposed.


Fig.2: Elevation of the test section. A disk holder is mounted vertically behind the top scraper at 230 mm from the test section. an angle of 45 degree to the surface of the test section (Fig.2). A photograph of the disc assembly is inserted in Fig. 2

The test section was exposed to a composite beam of Hydrogen and Deuterium ions and neutrals accelerated to 60 keV . The deuterium content in the beam increased steadily from initially $20 \%$ to $75 \%$ at the end of the panel charging. Deuterium content and beam composition where measured by $\mathrm{H}_{\alpha}$ spectroscopy. The composition of Hydrogen and Deuterium is derived from the unshifted peaks and shown in Fig.3, the species mix of the deuterium part of the beam is given in table 1 :

Table 1: Composition of the deuterium fraction (ions and neutrals) during charging

| Component | Full energy | Half energy | Third energy | Impurities |
| :---: | :---: | :---: | :---: | :---: |
| Composition (\% of power) | 43.4 | 10.5 | 34 | 12 |
| Commposition (\% of flux) | 22 | 10 | 50 | 18 |
| Energy per atom [keV] | 60 | 30 | 20 | 3.2 |

Fig. 4 shows a plan view of the experimental setup. The test section is partially shaded by a scraper. The inertial calorimeter is behind the test section.

The upper part of the test section can not be seen by the IR imaging system (Fig.2) and no measurement of the surface temperature is available for this upper part. By comparing the


Fig.4: Plan view of the test setup. The monoblock is partially shaded by a scraper. The calorimeter is mounted behind the test section.
vertical profiles from the surface temperature (IR), the bulk temperature (TC) and the power density (inertial calorimeter) (Fig. 5 and Fig.6) we can see that all three quantities are proportional to each other. This allows estimate a surface temperature from the bulk temperature by using the correlation $\mathrm{T}_{\text {surface }}=2.15 \times \mathrm{T}_{\text {bulk }}$ (temperature in ${ }^{\circ} \mathrm{C}$ ).


Fig.5: Vertical profile of the surface temperature along a line through the hottest parts. Surface temperature and bulk temperature ( 10 mm below the exposed surface) have the same profile.


Fig.6: Vertical profile of the bulk temperature ( 10 mm below the exposed surface) and of the power density from the inertial calorimeter. The profiles are identical.

The time constants for the exponential temperature rise during exposure and the cooldown after exposure are 2.6 seconds for the heat up and 1 second for the cooldown (Fig.7). The test section was exposed to 6 s pulses with a typical duty cycle of 15 pulses per hour. The water cooling loop was at room temperature. The test section was exposed to a peak fluence of $210^{19}$ deuterium atoms $/ \mathrm{cm}^{2}$ during 450 s of exposure (Fig.8).

The main parameters of the first part of the experiment are shown in annex 1.


Fig.7: Time constants of the surface temperature.

### 2.1 Concentration measurement

After charging the test section was exposed to air for 6 days before the concentration was measured with Nuclear Reaction Analysis. The Monoblock was mounted in aspecial large vacuum chamber with Be-handling facilities attached to the University of Sussex 3 MeV accelerator. For D analysis a 2.5 MeV 3 He beam was used, and protons produced by the reaction

$$
\mathrm{D}\left({ }^{3} \mathrm{He}, \mathrm{p}\right)^{4} \mathrm{He}
$$

Simultaneously Be and C can be measured using the reactions

$$
{ }^{9} \mathrm{Be}\left({ }^{3} \mathrm{He}, \mathrm{p}\right){ }^{11} \mathrm{~B}, \quad{ }^{12} \mathrm{C}\left({ }^{3} \mathrm{He}, \mathrm{p}\right){ }^{14} \mathrm{~N}
$$

The exact energy of the proton peak from the reaction with D depends on the depth within the surface from which the proton was emitted, so that the peak shape gives a picture of the distribution of D within the surface: the maximum depth from which information can be derived with a 2.5 MeV beam is $8 \mu \mathrm{~m}$.

## 3. EXPERIMENTAL RESULTS.

### 3.1 Implanted concentration

Table 2 shows the measured concentration for tiles $1,3,5,7,9$, and 11 , measured in the center of the tiles, and the typical bulk temperature rise of the respective tile, measured with a thermocouple 10 mm below the exposed surface. The tile numbering is from top to bottom.

| Table 2: Tc temperature and deuterium concentration |  |  |  |  |
| :---: | :---: | :---: | :---: | :---: |
|  | run without surface melting <br> concentration <br> $10^{18} / \mathrm{cm}^{2}$ |  | temperature $\left({ }^{\circ} \mathrm{C}\right)^{1}$ | run with surface melting <br> concentration <br> $10^{18} / \mathrm{cm}^{2}$ |
| block temperature $\left({ }^{\circ} \mathrm{C}\right)^{2}$ |  |  |  |  |
| 1 | 1.7 | $51(383)$ | 0.36 | $169(1070)$ |
| 2 | 1.74 | $69(421)$ | 0.13 | $(1270)$ |
| 3 | 1.63 | $86(458)$ | 0.009 | $244(1470)$ |
| 5 | 1.51 | $109(508)$ | 0 | 320 |
| 7 | 1.66 | $120(531)$ | 0 | 400 |
| 9 | 1.67 | $121(533)$ | 0 | 449 |
| 11 |  |  |  | 443 |

[^0]The measured concentrations after the charging run at $5 \mathrm{MW} / \mathrm{m}^{2}$ are almost uniform. Using the TRIM code for calculating the penetration depth and assuming that the implanted particles come to rest at the end of the range, we get the concentration profile shown in Fig. 9 for the deuterium atoms in the beam (impurities are not taken into account).

As only $75 \%$ of the beam were deuterium atoms we estimate that the concentration in a pure deuterium beam would have been 1.33 times the measured concentration, say $2.210^{18}$ atoms $/ \mathrm{cm}^{2}$. Using an implanted range of 0.2-0.6 micron (Fig.9) for $80 \%$ of the flux (impurities neglected) we get an implanted density of $410^{22}$ atoms $/ \mathrm{cm}^{3}$ ( $30 \mathrm{at} \%$ )

After the second charging scan, in which the surface temperature on most tiles was well above liquidus, the deuterium is released on all measured tiles apart from the uppermost three tiles, which stayed below liquidus. That the upper 3 tiles stayed below liquidus can be seen from


Fig.9: Deuterium deposition profile from TRIM.
Fig. 10 which shows the test section before the first charging scan at $5 \mathrm{MW} / \mathrm{m}^{2}$ (left) and after the second charging scan with surface melting (right). The actual surface temperature of the three unmelted tiles at the top could not be measured with the IR imaging system. If we assume that tile 3 was just below liquidus - say $1200^{\circ} \mathrm{C}$ - and that the surface temperature is proportional to the thermocouple temperature we estimate a temperature of $800^{\circ} \mathrm{C}$ for tile 1 and $1000^{\circ} \mathrm{C}$ for tile 2.


Fig.10: Photographs of the test section before the charging run at $540 K$ (left) and after the second charging run at high temperatures (right). No additional melting is observed on the upper three tiles.

### 3.2 Distribution of the implanted Deuterium

Figure 11 shows a comparison of the NRA spectrum from the center of the Monoblock (a) with a spectrum from a point on one of the Mk I Be divertor tiles after use in JET during 1995 (b). The spectrum (c) is recorded from a graphite sample implanted with $6 \times 10^{17}$ atoms $\mathrm{cm}^{-2}$ of D at 5 keV , which TRIM calculations suggest should give a film $0.15 \mu \mathrm{~m}$ thick saturated with D (i.e. to a D:C ratio of 0.4:1).

The D peaks in (a) and (c) are at the same channel number (which is proportional to energy) and of the same width. Furthermore spectrum (b) shows that for a film with D present throughout the analysable depth the peak would be much wider, to higher channel numbers. (Reactions occurring at greater depth lead to protons of greater energy due to the kinematics of the reaction.) The peak from the implanted standard has a half-width perhaps a factor of ten less than from the thick ( $8 \mu \mathrm{~m}$ or more) film, but the D is actually within $0.15 \mu \mathrm{~m}$, so this peak width represents the resolution of the detector. The peak from the monoblock is the same width, so we can conclude that the D is certainly all within the first micron, but we cannot say how much thinner than that the layer is.


Fig.11: comparison of the NRA spectrum from the center of the Monoblock (a) with a spectrum from a point on one of the Mk I Be divertor tiles after use in JET during 1995 (b). The spectrum (c) is recorded from a graphite sample implanted with $6 \times 10^{17}$ atoms $\mathrm{cm}^{-2}$ of $D$ at 5 keV , which TRIM calculations suggest should give a film $0.15 \mu \mathrm{~m}$ thick saturated with $D$ (i.e. to a $D: C$ ratio of 0.4:1).

### 3.3 Implantation on the disc holder assembly.

After the first exposure with surface temperatures of $500-550 \mathrm{~K}$ we find an average Beryllium layer of $310^{16}$ atoms $/ \mathrm{cm}^{2}$ corresponding to a depth of $310^{-3} \mu \mathrm{~m}$. The Deuterium/Beryllium
ratio is rather high but this might be wrong because of the contribution from the base material. After the second exposure the Beryllium deposit is up by two order of magnitude to $510^{18}$ atoms $/ \mathrm{cm}^{2}$ corresponding to $0.4 \mu \mathrm{~m}$. The beam fluence in both experiments was almost the same and the thicker layer must be a consequence of the much higher evaporation rate in the second exposure. The Deuterium content increased roughly by a factor of 6 and the Deuterium Beryllium ratio is now only $3-4 \%$.

| Table 3:Deposition and implantation on the disc holder assembly in atoms $/ \mathrm{cm}^{2}$ |  |  |  |  |
| :---: | :---: | :---: | :---: | :---: |
|  | First exposure (520K) |  |  |  |
| Position | C | BE | D | D/BE |
| upper disc cover | $8 \mathrm{E}+16$ | $3 \mathrm{E}+16$ | $3 \mathrm{E}+15$ | 9\% |
| upper disc | $4 \mathrm{E}+16$ | 4E++16 | $3 \mathrm{E}+16$ | 74\% |
| middle glass | $3 \mathrm{E}+16$ | $5 \mathrm{E}+16$ | $3 \mathrm{E}+16$ | 57\% |
| middle glass cover | 7E+16 | $4 \mathrm{E}+16$ | $3 \mathrm{E}+16$ | 61\% |
| copper base | $3 \mathrm{E}+16$ | 5E++16 | $3 \mathrm{E}+16$ | 60\% |
| lower disc | $3 \mathrm{E}+16$ | $4 \mathrm{E}+16$ | $3 \mathrm{E}+16$ | 76\% |
| lower disc cover | 7E+16 | $3 \mathrm{E}+16$ | $1 \mathrm{E}+16$ | 34\% |
|  | Second exposure (>2000K) |  |  |  |
|  | C | BE | D | D/BE |
| upper disc | $1 \mathrm{E}+18$ | $4 \mathrm{E}+18$ | $1 \mathrm{E}+17$ | 3\% |
| middle glass | $1 \mathrm{E}+18$ | $5 \mathrm{E}+18$ | $2 \mathrm{E}+17$ | 3\% |
| middle glass cover | $1 \mathrm{E}+18$ | $5 \mathrm{E}+18$ | $2 \mathrm{E}+17$ | 4\% |
| lower disc | $9 \mathrm{E}+17$ | $4 \mathrm{E}+18$ | $2 \mathrm{E}+17$ | 4\% |
| lower disc cover | $5 \mathrm{E}+17$ | 4E+18 | $1 \mathrm{E}+17$ | 4\% |

## 4. DISCUSSION

Our results agree perfectly well with previous measurements carried out by Moeller on Beryllium supplied by JET (S65) [6]. Fig. 12 shows a summary of these results together with the


Fig.12: Comparison of the retained Deuterium in various experiments.
concentration from our first charging scan. More recent retention measurements in Beryllium with lower particle fluxes and lower energy largely agree with our results if we use the scaling of Fig. 12.

1. Yoshida observes a retention of $10^{17} / \mathrm{cm}^{2}$ after implantation at $673^{\circ} \mathrm{K}$ with $8 \mathrm{keV} \mathrm{D}_{3}{ }^{+}$ions [7]. This is $5 \%$ of the implantation we find. The particle energy is roughly $10 \%$ of our energy and the particle flux is with $310^{14}$ atoms $/ \mathrm{cm}^{2}$ roughly $1 \%$ of our flux. The data point is added in Fig. 12.
2. Causey [8] finds in a 100 eV plasma exposure a retention of $310^{16} / \mathrm{cm}^{2}$ at a temperature of $573^{\circ} \mathrm{K}$ an a flux in the range of $1-310^{18}$ atoms $/ \mathrm{cm}^{2}$. This implantation is roughly $1 \%$ of our measurement, but the penetration depth is also much smaller.
3. Anderl [9] reports an implantation of $0.1 \mathrm{at} \%$ from a $3 \mathrm{keV} \mathrm{D}_{3}{ }^{+}$beam with a flux of $510^{15}$ atoms $/ \mathrm{cm}^{2}$ on a target of 623 and $703{ }^{\circ} \mathrm{K}$. The peak concentration was at a depth of 0.2 $0.4 \mu \mathrm{~m}$ while the implantation depth is quoted as $0.016 \mu \mathrm{~m}$. We find an implanted concentration of $30 \mathrm{at} \%$ if we believe in the deposition profile derived from TRIM.

The hydrogen density trapped in Beryllium is very high, but as the range over which hydrogen is trapped is very narrow, the total quantity of trapped hydrogen is modest, provided the hydrogen does not spread with time and fluence. Experiments to measure the spreading of the trapped hydrogen are in preparation.

The Deuterium/Beryllium ratio on the disc holder is with 3-4\% one order of magnitude lower than that reported by Mayer [10] in an experiment where the Beryllium and Deuterium on the target plate originate from the sputtering/reflection of a $4.5 \mathrm{keV} \mathrm{D}_{3}{ }^{+}$beam on Beryllium. In our case the Beryllium originates from evaporation the codeposited deuterium can originate from reflected ions and neutrals, from gas collisions (the pressure is 2-3 $10^{-3} \mathrm{mbar}$ ) or from interactions of the target with the beam plasma formed by collisions between beam particles and
background gas. The difference between our result and that of Mayer indicates that the quantity of codeposition is likely to depend on both fluxes - hydrogen and beryllium. More experiments are required to meaure the influence of energy, temperature and flux ratio.

We also conclude, that in the case of Beryllium an almost identical hydrogen retention is found in experiments with low and high fluxes. This gives confidence that the measurement is not sensitive to flux and energy.

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| Pulse Number | Notes | Beam Voltage [kV] |  | Extractedpower[MW] | Peak PD [MW/m²] | Peak PD [MW/m²] | Perveance[uperv] | Beam On Time [s] |  | $\begin{array}{\|c\|} \hline \text { MAX TC } \\ \hline \text { TEMP } \\ \hline \end{array}$ | $\begin{array}{\|c\|} \hline \text { average } \\ \hline \text { flux } \\ \hline \end{array}$ | Cumulativ On time | measured <br> ratio | flux <br> deut. | cumm. <br> flux |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  |  | set | measured |  |  |  |  | set | measured |  |  |  |  |  |  |
|  | ITER Be monoblock - 2nd test |  |  |  |  |  |  |  |  | [ ${ }^{\text {C] }}$ | [cm-2] |  | D/(H+D) | CM-2 | D2/cm ${ }^{2}$ |
|  | 02-Aug-96 |  |  |  | Water Cal | Inertial Cal |  |  |  |  |  |  |  |  |  |
| FORMULAS |  |  |  | 0.00 |  |  |  |  |  |  |  |  |  |  |  |
| 88908 |  | 60 | 58.70 | 0.17 | 5.27 | 5.46 | 0.203 | 6.041 | 6.000 |  | $4.20 \mathrm{E}+17$ | 203.030 | $6.68 \mathrm{E}-01$ | $2.80 \mathrm{E}+17$ | $8.52 \mathrm{E}+18$ |
| 88909 |  | 60 | 58.20 | 0.17 | 5.27 | 5.46 | 0.204 | 6.041 | 6.000 |  | $4.20 \mathrm{E}+17$ | 209.030 | $6.60 \mathrm{E}-01$ | $2.81 \mathrm{E}+17$ | $8.80 \mathrm{E}+18$ |
| 88910 |  | 60 | 58.10 | 0.17 | 4.91 | 5.51 | 0.207 | 6.041 | 6.000 |  | $3.95 \mathrm{E}+17$ | 215.030 | $6.71 \mathrm{E}-01$ | $2.66 \mathrm{E}+17$ | $9.07 \mathrm{E}+18$ |
| 88911 |  | 60 | 58.20 | 0.17 | 4.96 | 5.51 | 0.207 | 6.041 | 6.000 |  | $3.95 \mathrm{E}+17$ | 221.030 | $6.82 \mathrm{E}-01$ | $2.67 \mathrm{E}+17$ | $9.33 \mathrm{E}+18$ |
| 88912 |  | 60 | 58.20 | 0.17 | 4.96 | 5.51 | 0.206 | 6.041 | 6.000 |  | $3.95 \mathrm{E}+17$ | 227.030 | $6.80 \mathrm{E}-01$ | $2.68 \mathrm{E}+17$ | 9.60E+18 |
| 88913 |  | 60 | 58.00 | 0.17 | 4.96 | 5.51 | 0.207 | 6.041 | 6.000 |  | $3.95 \mathrm{E}+17$ | 233.030 | $6.85 \mathrm{E}-01$ | $2.69 \mathrm{E}+17$ | 9.87E+18 |
| 88914 |  | 60 | 58.00 | 0.17 | 4.96 | 5.51 | 0.207 | 6.041 | 6.000 |  | $3.95 \mathrm{E}+17$ | 239.030 | $6.84 \mathrm{E}-01$ | $2.70 \mathrm{E}+17$ | $1.01 \mathrm{E}+19$ |
| 88915 |  | 60 | 57.80 | 0.17 | 4.96 | 5.51 | 0.206 | 6.041 | 6.000 |  | $3.95 \mathrm{E}+17$ | 245.030 | $6.85 \mathrm{E}-01$ | $2.72 \mathrm{E}+17$ | $1.04 \mathrm{E}+19$ |
| 88916 | g3 ex. error | 60 |  | 0.00 | 4.96 | 5.51 |  | 6.041 |  |  |  | 245.030 | ERR | $0.00 \mathrm{E}+17$ | $1.04 \mathrm{E}+19$ |
| 88917 |  | 60 | 57.40 | 0.16 | 5.37 | 5.63 | 0.204 | 6.041 | 6.000 | 137 | $4.28 \mathrm{E}+17$ | 251.030 | $6.59 \mathrm{E}-01$ | $2.96 \mathrm{E}+17$ | $1.07 \mathrm{E}+19$ |
| 88918 |  | 60 | 57.30 | 0.16 | 5.37 | 5.63 | 0.206 | 6.041 | 6.000 |  | $4.28 \mathrm{E}+17$ | 257.030 | $6.85 \mathrm{E}-01$ | $2.97 \mathrm{E}+17$ | 1.10E+19 |
| 88919 |  | 60 | 57.10 | 0.17 | 5.37 | 5.63 | 0.210 | 6.041 | 6.000 |  | $4.28 \mathrm{E}+17$ | 263.030 | $6.94 \mathrm{E}-01$ | $2.98 \mathrm{E}+17$ | 1.13E+19 |
| 88920 |  | 60 | 56.80 | 0.16 | 5.37 | 5.63 | 0.208 | 6.041 | 6.000 | 136 | $4.28 \mathrm{E}+17$ | 269.030 | 7.03E-01 | $2.99 \mathrm{E}+17$ | 1.16E+19 |
| 88921 |  | 60 | 56.60 | 0.16 | 5.37 | 5.63 | 0.209 | 6.041 | 6.000 |  | $4.28 \mathrm{E}+17$ | 275.030 | $6.99 \mathrm{E}-01$ | $3.00 \mathrm{E}+17$ | 1.19E+19 |
| 88922 |  | 60 | 58.00 | 0.17 | 5.37 | 5.63 | 0.207 | 6.041 | 6.000 |  | $4.28 \mathrm{E}+17$ | 281.030 | $7.12 \mathrm{E}-01$ | $3.01 \mathrm{E}+17$ | $1.22 \mathrm{E}+19$ |
| 88923 |  | 60 | 58.30 | 0.17 | 5.37 | 5.63 | 0.203 | 6.041 | 6.000 |  | $4.28 \mathrm{E}+17$ | 287.030 | $7.08 \mathrm{E}-01$ | $3.02 \mathrm{E}+17$ | $1.25 \mathrm{E}+19$ |
| 88924 |  | 60 | 58.00 | 0.17 | 5.37 | 5.63 | 0.205 | 6.041 | 6.000 |  | $4.28 \mathrm{E}+17$ | 293.030 | 7.05E-01 | $3.03 \mathrm{E}+17$ | $1.28 \mathrm{E}+19$ |
| 88925 | G3 ex error | 60 |  | 0.00 | 5.37 | 5.63 |  | 6.041 | 6.000 |  |  | 293.030 | ERR | $0.00 \mathrm{E}+00$ | $1.28 \mathrm{E}+19$ |
| 88926 |  | 60 | 58.60 | 0.17 | 5.37 | 5.63 | 0.199 | 6.041 | 6.000 |  | $4.28 \mathrm{E}+17$ | 299.030 | $6.95 \mathrm{E}-01$ | $3.04 \mathrm{E}+17$ | $1.31 \mathrm{E}+19$ |
| 88927 |  | 60 | 58.60 | 0.17 | 5.37 | 5.63 | 0.202 | 6.041 | 6.000 |  | $4.28 \mathrm{E}+17$ | 305.030 | $7.11 \mathrm{E}-01$ | $3.04 \mathrm{E}+17$ | 1.34E+19 |
| 88928 |  | 60 | 58.70 | 0.17 | 5.37 | 5.63 | 0.202 | 6.041 | 6.000 |  | $4.28 \mathrm{E}+17$ | 311.030 | $7.10 \mathrm{E}-01$ | $3.05 \mathrm{E}+17$ | $1.37 \mathrm{E}+19$ |
| 88929 |  | 60 | 58.60 | 0.17 | 5.37 | 5.63 | 0.202 | 6.041 | 6.000 |  | $4.28 \mathrm{E}+17$ | 317.030 | 7.12E-01 | $3.06 \mathrm{E}+17$ | $1.40 \mathrm{E}+19$ |
| 88930 | trigger IR | 60 | 58.50 | 0.17 | 4.62 | 5.34 | 0.201 | 6.041 | 6.000 | 130 | $3.68 \mathrm{E}+17$ | 323.030 | $7.19 \mathrm{E}-01$ | $2.64 \mathrm{E}+17$ | $1.43 \mathrm{E}+19$ |
| 88931 |  | 60 | 58.70 | 0.17 | 4.62 | 5.34 | 0.201 | 6.041 | 6.000 |  | $3.68 \mathrm{E}+17$ | 329.030 | $7.21 \mathrm{E}-01$ | $2.65 \mathrm{E}+17$ | $1.46 \mathrm{E}+19$ |
| 88932 |  | 60 | 58.70 | 0.17 | 4.62 | 5.34 | 0.202 | 6.041 | 6.000 |  | $3.68 \mathrm{E}+17$ | 335.030 | $7.20 \mathrm{E}-01$ | $2.65 \mathrm{E}+17$ | $1.48 \mathrm{E}+19$ |
| 88933 |  | 60 | 58.40 | 0.17 | 4.62 | 5.34 | 0.203 | 6.041 | 6.000 |  | $3.68 \mathrm{E}+17$ | 341.030 | $7.19 \mathrm{E}-01$ | $2.66 \mathrm{E}+17$ | $1.51 \mathrm{E}+19$ |
| 88934 |  | 60 | 58.10 | 0.17 | 4.62 | 5.34 | 0.203 | 6.041 | 6.000 |  | $3.68 \mathrm{E}+17$ | 347.030 | $7.38 \mathrm{E}-01$ | $2.66 \mathrm{E}+17$ | $1.54 \mathrm{E}+19$ |
| 88935 |  | 60 | 58.30 | 0.17 | 4.62 | 5.34 | 0.205 | 6.041 | 6.000 |  | $3.68 \mathrm{E}+17$ | 353.030 | $7.28 \mathrm{E}-01$ | $2.67 \mathrm{E}+17$ | $1.56 \mathrm{E}+19$ |
| 88936 |  | 60 | 58.10 | 0.17 | 4.62 | 5.34 | 0.204 | 6.041 | 6.000 |  | $3.68 \mathrm{E}+17$ | 359.030 | 7.19E-01 | $2.67 \mathrm{E}+17$ | $1.59 \mathrm{E}+19$ |
| 88937 |  | 60 | 58.40 | 0.17 | 4.62 | 5.34 | 0.201 | 6.041 | 6.000 |  | $3.68 \mathrm{E}+17$ | 365.030 | $7.27 \mathrm{E}-01$ | $2.68 \mathrm{E}+17$ | $1.62 \mathrm{E}+19$ |
| 88938 |  | 60 | 58.60 | 0.17 | 4.62 | 5.34 | 0.200 | 6.041 | 6.000 |  | $3.68 \mathrm{E}+17$ | 371.030 | $7.30 \mathrm{E}-01$ | $2.68 \mathrm{E}+17$ | $1.64 \mathrm{E}+19$ |
| 88939 |  | 60 | 58.80 | 0.17 | 4.62 | 5.34 | 0.197 | 6.041 | 6.000 |  | $3.68 \mathrm{E}+17$ | 377.030 | $7.25 \mathrm{E}-01$ | $2.69 \mathrm{E}+17$ | $1.67 \mathrm{E}+19$ |
| 88940 | trigger IR | 60 | 58.90 | 0.17 | 5.06 | 5.06 | 0.199 | 6.041 | 6.000 | 126 | $4.03 \mathrm{E}+17$ | 383.030 | $7.29 \mathrm{E}-01$ | $2.95 \mathrm{E}+17$ | $1.70 \mathrm{E}+19$ |
| 88941 |  | 60 | 58.90 | 0.17 | 5.06 | 5.06 | 0.201 | 6.041 | 6.000 |  | $4.03 \mathrm{E}+17$ | 389.030 | $7.31 \mathrm{E}-01$ | $2.95 \mathrm{E}+17$ | $1.73 \mathrm{E}+19$ |
| 88942 |  | 60 | 59.90 | 0.17 | 5.06 | 5.06 | 0.196 | 6.041 | 6.000 |  | $4.03 \mathrm{E}+17$ | 395.030 | 7.31E-01 | $2.96 \mathrm{E}+17$ | $1.76 \mathrm{E}+19$ |
| 88943 |  | 60 | 59.60 | 0.17 | 5.06 | 5.06 | 0.197 | 6.041 | 6.000 |  | $4.03 \mathrm{E}+17$ | 401.030 | $7.30 \mathrm{E}-01$ | $2.96 \mathrm{E}+17$ | $1.79 \mathrm{E}+19$ |
| 88944 |  | 60 | 59.90 | 0.17 | 5.06 | 5.06 | 0.194 | 6.041 | 6.000 |  | $4.03 \mathrm{E}+17$ | 407.030 | 7.28E-01 | $2.97 \mathrm{E}+17$ | $1.82 \mathrm{E}+19$ |
| 88945 |  | 60 | 59.80 | 0.17 | 5.06 | 5.06 | 0.196 | 6.041 | 6.000 |  | $4.03 \mathrm{E}+17$ | 413.030 | $7.31 \mathrm{E}-01$ | $2.97 \mathrm{E}+17$ | $1.85 \mathrm{E}+19$ |
| 88946 |  | 60 | 59.20 | 0.17 | 5.06 | 5.06 | 0.201 | 6.041 | 6.000 |  | $4.03 \mathrm{E}+17$ | 419.030 | $7.42 \mathrm{E}-01$ | $2.98 \mathrm{E}+17$ | $1.88 \mathrm{E}+19$ |
| 88947 |  | 60 | 59.80 | 0.18 | 5.06 | 5.06 | 0.198 | 6.041 | 6.000 |  | $4.03 \mathrm{E}+17$ | 425.030 | $7.28 \mathrm{E}-01$ | $2.98 \mathrm{E}+17$ | $1.91 \mathrm{E}+19$ |
| 88948 |  | 60 | 59.80 | 0.17 | 5.06 | 5.06 | 0.196 | 6.041 | 6.000 |  | $4.03 \mathrm{E}+17$ | 431.030 | $7.41 \mathrm{E}-01$ | $2.98 \mathrm{E}+17$ | $1.94 \mathrm{E}+19$ |
| 88949 |  | 60 | 59.80 | 0.17 | 5.06 | 5.06 | 0.106 | 6.041 | 6.000 |  | $4.03 \mathrm{E}+17$ | 437.030 | $7.38 \mathrm{E}-01$ | $2.99 \mathrm{E}+17$ | $1.97 \mathrm{E}+19$ |
| 88950 | Trigger IR | 60 | 58.70 | 0.17 | 4.93 | 4.89 | 0.203 | 6.041 | 6.000 | 121 | $3.93 \mathrm{E}+17$ | 443.030 | 7.48E-01 | $2.92 \mathrm{E}+17$ | $2.00 \mathrm{E}+19$ |


[^0]:    ${ }^{1}$ This is the peak TC temperature measured 10 mm below the exposed surface. In brackets the estimated surface temperature is given in ${ }^{\circ} \mathrm{K}$
    ${ }^{2}$ In brackets given is a rough estimate of the surface temperature as explained in the text.

