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A New Assessment of Cross Section Data for Helium Excitation by Protons

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ABSTRACT.

We present a new assessment of cross section data for excitation of $He(l^1S)$ by proton impact at energies larger than 10 keV (in a few cases 6 keV). Data for excitation to n^1L states of He ($n \le 4$, L = S, P and D) are given in tables. The data at high energy (-1000 keV) are linked to the first Born approximation and at low energies in some cases to the close coupling AO results of Fritsch, who also made a critical evaluation of the cross sections for excitation of He by protons.

Introduction.

A new assessment of proton impact cross-section data for excitation of helium from its I¹S ground state is presented. Similarly as our studies of excitation of helium by electron impact [1] and electron capture from helium by protons [2], this work was initiated in connection with the installation and recent use of ³He and ⁴He neutral heating beams on the JET tokamak [3]. The impact energy range covered extends up to higher impact energies than the ones of direct interest for the present beam based diagnostics at JET (energies smaller 200 keV) in order to be also of relevance to possible future He beam based diagnostics. At ITER neutral heating beams of a few hundreds of keV/amu are foreseen to closer match the velocity of the fusion-produced alpha particles [4].

Schematically the processes studied here are given by

$$H^+ + He(l^1S) \rightarrow H^+ + He^*(n^1L)$$
 (1)

Except for the 2¹S state, the excited states decay under photon emission and so it is appropriate to study helium excitation by means of photon emission spectroscopy with absolute intensity measurement. Therefore this papers brings back into memory the often forgotten compilation by Thomas [5] about optical experiments in ion-atom collisions. Besides it includes more recent work of this kind of the group of Schartner, e.g. refs. [6,7]. As mentioned before the 2¹S state of He cannot be observed optically and for the excitation of that state we have considered the energy loss experiments of the group of Park [8,9], who also measured the 2¹P excitation, for which optical work has only been performed by Hippler and Schartner [7]. The most recent compilation of cross section data has been provided by Barnett and coworkers [10]. Our recommended cross sections sometimes deviate from those given in that work.

Simultaneously with our study Fritsch [11] has made a critical evaluation of existing data for helium excitation in heavy-particle collisions, including proton impact, along with new calculations within the close coupling framework with atomic basis sets applied to excitation of He 2¹L and 3 ¹L levels. Our analysis for proton impact has a lot of overlapping with the work of Fritsch and therefore we only give some complementary and relevant information, i.e. data tables of recommended excitation cross sections.

Results and discussion

Compared with the compilation of Thomas [5] the principal new contribution in this paper is the inclusion of more recent experimental [6,7,9] and theoretical [11,12] work. We shall generally confine ourselves to impact energies above 10 keV. In the article of Thomas we can see that the cross sections obtained by the various groups often show approximately the same energy dependence, but differ in absolute scale. Therefore Thomas often applied a normalization procedure at one impact energy and then plotted the data again in one graph.

In order to establish an absolute scale we give generally preference to the data of Schartner's group [6,7] (for 2¹P, 3¹P and 4¹L levels) between 150 and 1000 keV, because at high energies they merge very well into the Born approximation. The corresponding cross sections have been calculated by Bell et al [13]. At these relatively high energies we assume for the missing levels (2¹S, 3¹S, 3¹D and $n^{1}L$ ($n \ge 4$)) that the cross section ratio's in a term series are independent of energy and have the same value as in the Born approximation. It appears experimentally that this ratio is approximately constant down to about 40 keV. Deviations then are largest for the lower levels in a term series where the difference in excitation energy is relatively the largest. This can be seen in the low energy experimental data of van den Bos et al [14] and the calculations of Fritsch [11] and Slim et al [12].

As we have remarked before, data below 150 keV often differ on absolute scale and have to be normalized. In this work we have often given preference to the optical data of van den Bos et al [14] between 1 and 150 keV and fitted them to the high energy data of Hasselkamp et al [6] and Hippler and Schartner [7] for the following two reasons: The energy dependence of the cross sections of van den Bos et al [14] is close to the average of various data sets as shown by Thomas [5] and in several cases their absolute scale agrees within experimental errors with that of Schartner's group [6,7]. The scaling factors for van den Bos' data used in the fitting procedure are mostly close to those of Fritsch [11] but deviate in a few cases as indicated furtheron.

Experimental energy loss measurements of Kvale et al [9] giving cross

sections for 2¹S and 2¹P at 25, 50, 75 and 100 keV generally confirm the consistency of our procedure.

Table 1 summarises the primary choices of experimental data and theoretical data together with the assessments of the accuracy of the cross over various the energy intervals. Table 2 summarizes recommended cross sections for excitation to the different $He(n^{1}L)$ levels, $n \le 4$ and L = S, P and D from the He(l^1 S) ground state. For n > 4 we recommend to extrapolate the cross sections with the same ratios as in the Born approximation at 100 keV (see ref. [13] where these cross sections are tabulated up to 7^{1} S, 6^{1} P and 6^{1} D). For higher *n*-values one may extrapolate according to the n^{*-3} proportionality where n^{*} is the effective principal quantum number. All these extrapolated cross sections have approximately similar errors as the corresponding n = 4 levels at comparable impact energies.

Excitation to the n^1F (G, H,...) levels is relatively small and negligible.

Conclusions

As also pointed ou by Fritsch [11] the excitation cross sections for H^+ -He(I^1S) collisions are generally well established particularly in the energy range relevant for nuclear fusion applications, i.e. larger than 10 keV impact energy. The higher the impact energy, the more accurate cross sections have been obtained experimentally and theoretically. Near one thousand keV generally the theoretical Born values are expected to be within 5%. At lower impact energies optical experiments give data accurate to 10% down to about 150 keV and at still lower energies scaled experimental data are expected to be better than 10-30% from 100 keV down to 10 keV. Close coupling extended atomic orbital calculations have been used to predict missing experimental data for 2^1S and 2^1P at low energies [11] and appear to be consistent with experiment for 3^1S and 3^1P and 3^1D below about 30 keV.

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Table 1. $H^+ + He(I^1S) \rightarrow H^+ + He(n^1L)$

state	energy (keV)	method	reference	ассигасу
2 ¹ S	6 - 28	AO	Fritsch [11]	~20%
	40 - 150	11.6 $\sigma(4^{1}S)$	van den Bos [14]	20 - 10%
	150 - 1000	11.6 $\sigma(4^{1}S)$	Hasselkamp et al [6]	< 10%
_	> 1000	Born	Bell et al [13]	< 5%
31S	10 - 150	opt.exp. $\times 0.784$	van den Bos [14]	30 - 10%
	200 - 1000	2.61 $\sigma(4^{1}S)$	Hasselkamp et al [6]	< 10%
	> 1000	Born	Bell et al [13]	< 5%
41S	10 - 150	opt.exp.	van den Bos [14]	30 - 10%
	200 - 1000	opt.exp.	Hasselkamp et al [6]	< 10%
	> 1000	Born	Bell et al [13]	< 5%
2 ¹ P	6 - 28	AO	Fritsch [11]	~20%
	30 - 150	$4.09 \ \sigma(3^{1}P)$	van den Bos [14]	20 - 10%
	150 - 1000	opt.exp.	Hippler and Schartner [7]	< 10%
	> 500	Born	Bell et al [13]	< 5%
3 ¹ P	10 - 150	opt.exp. ×0.82	van den Bos [14]	30 - 10%
	150 - 1000	opt.exp.	Hippler and Schartner[7]	< 10%
	> 500	Born	Bell et al [13]	< 5%
4 ¹ P	10 - 150	opt.exp. ×1.25	van den Bos [14]	30 - 10%
	200 - 1000	opt.exp.	Hippler and Schartner [7]	< 10%
	> 1000	Born	Bell et al [13]	< 5%
3 ¹ D	10 - 150	opt.exp. \times 0.8	van den Bos [14]	30 - 10%
	200 - 1000	1.88 $\sigma(4^{1}D)$	Hasselkamp et al [6]	< 10%
	> 1000	Born	Bell et al [13]	< 5%
4 ¹ D	10 - 150	opt.exp. × 0.9	van den Bos [14]	30 - 10%
	200 - 1000	opt.exp.	Hasselkamp et al [6]	< 10%
	> 1000	Born	Bell et al [13]	< 5%

Note: Fritsch [11] used the following correction factors in the case of the data of van den Bos [14]: 0.74 for 3¹S, 0.75 for 3¹P, 1.14 for 4¹P, 1.47 for 5¹P, 0.82 for 3¹D, 0.9 for 4¹D and 1.12 for 5¹D.

Table 2. Recommended excitation cross sections (in units of 10^{-20} cm^2) for $H^+ + \text{He}(l^1\text{S}) \rightarrow H^+ + \text{He}(n^1\text{L})$

6 71.1 194.5 19.5 27.0 10 241 27.4 8.82 331.3 45.5 14.8 29.9 12.6 12.5 36.1 10.4 311.4 37.1 16.7 14 460 311.4 311.4 33.5 38.7 16.5 17.5 13.5 13.5 13.5 13.5 13.5 20 598 84.6 20.6 295.3 101 40.0 29.9 11.6 25 112 33.8 100 40.0 21.8 10.2 28 536 116 40.8 425 104 45.6 23.4 10.4 35 44.8 512 125 52.5 26.7 11.9 40.6 567 128 48.9 556 136 56.3 32.0 13.1 50.5 547 131 47.1 793 194 75.0 29.3 13.7 60.5 520 112 44.8	E(keV)	2 ¹ S	3 ¹ S	4 ¹ S	2 ¹ P	3 ¹ P	4 ¹ P	3 ¹ D	4 ¹ D
10	6	71.1			194.5				
14 460	10	241			331.3	45.5	14.8	29.9	
20 598 84.6 20.6 295.3 101 40.0 29.9 11.6 25 112 33.8 100 40.0 21.8 10.2 28 536 428 100 40.0 21.8 10.2 28 536 428 100 40.0 21.8 10.2 28 536 44.8 512 125 52.5 26.7 11.9 40 567 128 48.9 556 136 56.3 32.0 13.1 50 547 131 47.1 793 194 75.0 29.3 13.7 60 520 112 44.8 954 234 87.5 29.1 14.2 70 523 94.7 45.1 958 234 98.7 27.9 13.6 80 483 102.6 41.7 1152 282 117 25.8 13.3 90 451 89.3 38.9 1183 289 114 23.3 12.8 100 407 90.8 35.1 1273 311 119 21.6 11.4 110 372 83.0 32.1 1340 328 144 20.0 11.2 120 343 77.8 29.6 1306 320 131 18.9 10.3 130 317 67.6 27.3 1333 326 135 18.1 9.8 1 140 287 69.8 24.7 1350 330 129 17.1 8.9 1 150 256 63.7 22.1 1287 323 130 16.1 8.19 200 219 49.3 18.9 1211 299 117 12.8 6.0 300 145 32.6 12.5 1006 256 101 7.1 3.8 400 110 24.8 9.5 864 220 93 6.0 3.2 500 87 19.6 7.5 787 197 77 4.5 2.4 600 69.6 15.7 6.0 682 178 72 3.9 2.1 70 59.2 13.3 5.1 647 165 63 3.2 1.7 800 53.4 12.0 4.6 568 150 58 2.8 1.5 1000 19.78 4.46 1.71 304.2 75.1 30.1 1.037 0.552 4000 9.94 2.24 0.857 178.8 44.1 17.7 0.524	15	460	43.1	11.5	311.4	89.3	32.5	38.7	
28 536	20	598			295.3				11.6
35	28	536							
50 547 131 47.1 793 194 75.0 29.3 13.7 60 520 112 44.8 954 234 87.5 29.1 14.2 70 523 94.7 45.1 958 234 98.7 27.9 13.6 80 483 102.6 41.7 1152 282 117 25.8 13.3 90 451 89.3 38.9 1183 289 114 23.3 12.8 100 407 90.8 35.1 1273 311 119 21.6 11.4 110 372 83.0 32.1 1340 328 144 20.0 11.2 120 343 77.8 29.6 1306 320 131 18.9 10.3 130 317 67.6 27.3 1333 326 135 18.1 9.81 140 287 69.8 24.7 1350 330 129 17.1 8.91 150 256 63.7 22.1 1287 323 130 16.1 8.19 200 219 49.3 18.9 1211 299 117 12.8 6.0 300 145 32.6 12.5 1006 256 101 7.1 3.8 400 110 24.8 9.5 864 220 93 6.0 3.2 500 87 19.6 7.5 787 197 77 4.5 2.4 600 69.6 15.7 6.0 682 178 72 3.9 2.1 700 59.2 13.3 5.1 647 165 63 3.2 1.7 800 53.4 12.0 4.6 568 150 58 2.8 1.5 1000 41.8 9.4 3.6 505 129 51 2.1 1.1 Born cross section (Bell et al [13]) Born cross section (Bell et al [13])	35	567		44.8	512	125	52.5	26.7	11.9
70 523 94.7 45.1 958 234 98.7 27.9 13.6 80 483 102.6 41.7 1152 282 117 25.8 13.3 90 451 89.3 38.9 1183 289 114 23.3 12.8 100 407 90.8 35.1 1273 311 119 21.6 11.4 110 372 83.0 32.1 1340 328 144 20.0 11.2 120 343 77.8 29.6 1306 320 131 18.9 10.3 130 317 67.6 27.3 1333 326 135 18.1 9.81 140 287 69.8 24.7 1350 330 129 17.1 8.91 150 256 63.7 22.1 1287 323 130 16.1 8.19 200 219 49.3 18.9 1211 299 117 12.8 6.0 300 145 32.6 12.5 1006 256 101 7.1 3.8 400 110 24.8 9.5 864 220 93 6.0 3.2 500 87 19.6 7.5 787 197 77 4.5 2.4 600 69.6 15.7 6.0 682 178 72 3.9 2.1 700 59.2 13.3 5.1 647 165 63 3.2 1.7 800 53.4 12.0 4.6 568 150 58 2.8 1.5 1000 41.8 9.4 3.6 505 129 51 2.1 1.1 Born cross section (Bell et al [13])	50	547	131	47.1	793	194	75.0	29.3	13.7
90	70	523	94.7	45.1	958	234	98.7	27.9	13.6
110	90	451	89.3	38.9	1183	289	114	23.3	12.8
140 287 69.8 24.7 1350 330 129 17.1 8.91 150 256 63.7 22.1 1287 323 130 16.1 8.19 200 219 49.3 18.9 1211 299 117 12.8 6.0 300 145 32.6 12.5 1006 256 101 7.1 3.8 400 110 24.8 9.5 864 220 93 6.0 3.2 500 87 19.6 7.5 787 197 77 4.5 2.4 600 69.6 15.7 6.0 682 178 72 3.9 2.1 700 59.2 13.3 5.1 647 165 63 3.2 1.7 800 53.4 12.0 4.6 568 150 58 2.8 1.5 1000 41.8 9.4 3.6 505 129 51 2.1 1.1 Born cross section (Bell et al [13])	110			32.1	1340	320	131		10.3
200	140	287	69.8	24.7	1350	330	129	17.1	8.91
400	200	219	49.3	18.9	1211	299	117	12.8	6.0
600 69.6 15.7 6.0 682 178 72 3.9 2.1 700 59.2 13.3 5.1 647 165 63 3.2 1.7 800 53.4 12.0 4.6 568 150 58 2.8 1.5 1000 41.8 9.4 3.6 505 129 51 2.1 1.1 Born cross section (Bell et al [13]) 1000 39.2 8.83 3.38 502.6 124.2 49.9 2.035 1.08 1.374 0.732 2000 19.78 4.46 1.71 304.2 75.1 30.1 1.037 0.552 3000 13.23 2.98 1.14 223.6 55.1 22.1 0.696 4000 9.94 2.24 0.857 178.8 44.1 17.7 0.524	400	110	24.8	9.5	864	220	93	6.0	3.2
800 53.4 12.0 4.6 568 150 58 2.8 1.5 1000 41.8 9.4 3.6 505 129 51 2.1 1.1 Born cross section (Bell et al [13]) 1000 39.2 8.83 3.38 502.6 124.2 49.9 2.035 1.08 1.374 0.732 2000 19.78 4.46 1.71 304.2 75.1 30.1 1.037 0.552 3000 13.23 2.98 1.14 223.6 55.1 22.1 0.696 4000 9.94 2.24 0.857 178.8 44.1 17.7 0.524	600	69.6	15.7	6.0	682	178	72	3.9	2. 1
Born cross section (Bell et al [13]) 1000	800	53.4	12.0	4.6	568	150	58	2.8	1.5
1500 2000 19.78 4.46 1.71 304.2 75.1 30.1 1.037 0.552 3000 13.23 2.98 1.14 223.6 55.1 22.1 0.696 4000 9.94 2.24 0.857 178.8 44.1 17.7 0.524									
2000 19.78 4.46 1.71 304.2 75.1 30.1 1.037 0.552 3000 13.23 2.98 1.14 223.6 55.1 22.1 0.696 4000 9.94 2.24 0.857 178.8 44.1 17.7 0.524		39.2	8.83	3.38	502.6	124.2	49.9		
4000 9.94 2.24 0.857 178.8 44.1 17.7 0.524	2000							1.037	
3000 7.96 1.79 0.686 130.0 37.0 14.8 0.420					178.8				

ANNEX

P.-H. REBUT, A. GIBSON, M. HUGUET, J.M. ADAMS¹, B. ALPER, H. ALTMANN, A. ANDERSEN², P. ANDREW³, M. ANGELONE⁴, S. ALI-ARSHAD, P. BAIGGER, W. BAILEY, B. BALET, P. BARABASCHI, P. BARKER, R. BARNSLEY⁵, M. BARONIAN, D.V. BARTLETT, L. BAYLOR⁶, A.C. BELL, G. BENALI, P. BERTOLDI, E. BERTOLINI, V. BHATNAGAR, A.J. BICKLEY, D. BINDER, H. BINDSLEV², T. BONICELLI, S.J. BOOTH, G. BOSIA, M. BOTMAN, D. BOUCHER, P. BOUCQUEY, P. BREGER, H. BRELEN, H. BRINKSCHULTE, D. BROOKS, A. BROWN, T. BROWN, M. BRUSATI, S. BRYAN, J. BRZOZOWSKI⁷, R. BUCHSE²², T. BUDD, M. BURES, T. BUSINARO, P. BUTCHER, H. BUTTGEREIT, C. CALDWELL-NICHOLS, D.J. CAMPBELL, P. CARD, G. CELENTANO, C.D. CHALLIS, A.V. CHANKIN⁸, A. CHERUBINI, D. CHIRON, J. CHRISTIANSEN, P. CHUILON, R. CLAESEN, S. CLEMENT, E. CLIPSHAM, J.P. COAD, I.H. COFFEY⁹, A. COLTON, M. COMISKEY 10, S. CONROY, M. COOKE, D. COOPER, S. COOPER, J.G. CORDEY, W. CORE, G. CORRIGAN, S. CORTI, A.E. COSTLEY, G. COTTRELL, M. COX¹¹, P. CRIPWELL¹², O. Da COSTA, J. DAVIES, N. DAVIES, H. de BLANK, H. de ESCH, L. de KOCK, E. DEKSNIS, F. DELVART, G.B. DENNE-HINNOV, G. DESCHAMPS, W.J. DICKSON 13, K.J. DIETZ, S.L. DMITRENKO, M. DMITRIEVA 14, J. DOBBING, A. DOGLIO, N. DOLGETTA, S.E. DORLING. P.G. DOYLE, D.F. DÜCHS, H. DUQUENOY, A. EDWARDS, J. EHRENBERG, A. EKEDAHL, T. ELEVANT⁷, S.K. ERENTS¹¹, L.G. ERIKSSON, H. FAJEMIROKUN¹², H. FALTER, J. FREILING¹⁵, F. FREVILLE, C. FROGER, P. FROISSARD, K. FULLARD, M. GADEBERG, A. GALETSAS, T. GALLAGHER, D. GAMBIER, M. GARRIBBA, P. GAZE, R. GIANNELLA, R.D. GILL, A GIRARD, A. GONDHALEKAR, D. GOODALL¹¹, C. GORMEZANO, N.A. GOTTARDI, C. GOWERS, B.J. GREEN, B. GRIEVSON, R. HAANGE, A. HAIGH, C.J. HANCOCK, P.J. HARBOUR, T. HARTRAMPF, N.C. HAWKES¹¹, P. HAYNES¹¹, J.L. HEMMERICH, T. HENDER¹¹, J. HOEKZEMA, D. HOLLAND, M. HONE, L. HORTON, J. HOW, M. HUART, I. HUGHES, T.P. HUGHES¹⁰, M. HUGON, Y. HUO¹⁶, K. IDA¹⁷, B. INGRAM, M. IRVING, J. JACQUINOT, H. JAECKEL, J.F. JAEGER, G. JANESCHITZ, Z. JANKOVICZ¹⁸, O.N. JARVIS, F. JENSEN, E.M. JONES, H.D. JONES, L.P.D.F. JONES, S. JONES¹⁹, T.T.C. JONES, J.-F. JUNGER, F. JUNIQUE, A. KAYE, B.E. KEEN, M. KEILHACKER, G.J. KELLY, W. KERNER, A. KHUDOLEEV²¹, R. KONIG, A. KONSTANTELLOS, M. KOVANEN²⁰, G. KRAMER¹⁵, P. KUPSCHUS, R. LÄSSER, J.R. LAST, B. LAUNDY, L. LAURO-TARONI, M. LAVEYRY, K. LAWSON¹¹, M. LENNHOLM, J. LINGERTAT²², R.N. LITUNOVSKI, A. LOARTE, R. LOBEL, P. LOMAS, M. LOUGHLIN, C. LOWRY, J. LUPO, A.C. MAAS 15, J. MACHUZAK 19, B. MACKLIN, G. MADDISON 11, C.F. MAGGI 23, G. MAGYAR, W. MANDL 22, V. MARCHESE, G. MARCON, F. MARCUS, J. MART, D. MARTIN, E. MARTIN, R. MARTIN-SOLIS²⁴, P. MASSMANN, G. MATTHEWS, H. McBRYAN, G. McCRACKEN¹¹, J. McKIVITT, P. MERIGUET, P. MIELE, A. MILLER, J. MILLS, S.F. MILLS, P. MILLWARD, P. MILVERTON, E. MINARDI⁴, R. MOHANTI²⁵, P.L. MONDINO, D. MONTGOMERY²⁶, A. MONTVAI²⁷, P. MORGAN, H. MORSI, D. MUIR, G. MURPHY, R. MYRNÄS²⁸, F. NAVE²⁹, G. NEWBERT, M. NEWMAN, P. NIELSEN, P. NOLL, W. OBERT, D. O'BRIEN, J. ORCHARD, J. O'ROURKE, R. OSTROM, M. OTTAVIANI, M. PAIN, F. PAOLETTI, S. PAPASTERGIOU, W. PARSONS, D. PASINI, D. PATEL, A. PEACOCK, N. PEACOCK¹¹, R.J.M. PEARCE, D. PEARSON¹², J.F. PENG 16, R. PEPE DE SILVA, G. PERINIC, C. PERRY, M. PETROV 21, M.A. PICK, J. PLANCOULAINE, J.-P POFFÉ, R. PÖHLCHEN, F. PORCELLI, L. PORTE¹³, R. PRENTICE, S. PUPPIN, S. PUTVINSKII⁸, G. RADFORD³⁰, T. RAIMONDI, M.C. RAMOS DE ANDRADE, R. REICHLE, J. REID, S. RICHARDS, E. RIGHI, F. RIMINI, D. ROBINSON¹¹, A. ROLFE, R.T. ROSS, L. ROSSI, R. RUSS, P. RUTTER, H.C. SACK, G. SADLER, G. SAIBENE, J.L. SALANAVE, G. SANAZZARO, A. SANTAGIUSTINA, R. SARTORI, C. SBORCHIA, P. SCHILD, M. SCHMID, G. SCHMIDT³¹, B. SCHUNKE, S.M. SCOTT, L. SERIO, A SIBLEY, R. SIMONINI, A.C.C. SIPS, P. SMEULDERS, R. SMITH, R. STAGG, M. STAMP, P. STANGEBY³, R. STANKIEWICZ³², D.F. START, C.A. STEED, D. STORK, P.E. STOTT, P. STUBBERFIELD, D. SUMMERS, H. SUMMERS¹³, L. SVENSSON, J.A. TAGLE³³, M. TALBOT, A. TANGA, A. TARONI, C. TERELLA, A TERRINGTON, A TESINI, P.R. THOMAS, E. THOMPSON, K. THOMSEN, F. TIBONE, A. TISCORNIA, P. TREVALION, B. TUBBING, P. VAN BELLE, H. VAN DER BEKEN, G. VLASES, M. VON HELLERMANN, T. WADE, C. WALKER, R. WALTON³¹, D. WARD, M.L. WATKINS, N. WATKINS, M.J. WATSON, S. WEBER 34 , J. WESSON, T.J. WIJNANDS, J. WILKS, D. WILSON, T. WINKEL, R. WOLF, D. WONG, C. WOODWARD, Y. WU 35 , M. WYKES, D. YOUNG, I.D. YOUNG, L. ZANNELLI, A. ZOLFAGHARI 19 , W. ZWINGMANN

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