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Persistent Density Perturbations at Rational q Surfaces Following Pellet Injection in the Joint European Torus

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Abstract - In JET, the ablation of injected pellets produces a striking resonance effect when the pellets reach surfaces with q-values 1 and $3/2$. Subsequently, structures with mode numbers $m=1, n=1$ and $m=3, n=2$ are observed with the soft X-ray cameras for more than 2 s as compact snake-like perturbations. These structures, which persist through several sawtooth collapses, give information on the radii of the $q=1$ and $q=\frac{3}{2}$ surfaces and the q-profile evolution. The observations can be explained by the formation of magnetic islands.

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Introduction - The effects of injection of solid D₂ pellets in JET discharges have been studied¹⁾ using the soft X-ray imaging system. Immediately after pellet injection a very localised density and temperature perturbation has been found with m=1, n=1 topology. This perturbation exists for a very long time (> 2 s), is clearly associated with the q=1 magnetic surface and acts as a probe for this surface allowing, for example, the study of the position of the q=1 surface during a sawtooth cycle. Similar structures are seen with m=3, n=2 on the q=3/2 surface. Effects of ablation at rational q-values have been seen previously^{2,3)} as irregularities of the H_α emission.

The characteristics of these perturbations, their relationship to the q=1 and 3/2 surfaces, and possible explanations for their origin are presented.

Observation of the snake modulation - Two soft X-ray cameras^{4,5)} containing 100 detectors view the plasma with a spatial resolution of 7 cm in orthogonal directions at the same toroidal position as the D₂ pellet injector. Pellets of 2.2×10^{21} or 4.5×10^{21} atoms, injected radially in the equatorial plane into ohmically heated JET plasmas ($B_{\phi} = 2-3$ T, $I = 3.0-3.6$ MA, elongation of 1.4) with velocities of $\sim 1 \text{ km.s}^{-1}$, are detected with good temporal resolution (up to 100kHz) by the vertically mounted soft X-ray camera. The intense initial emission (figure 1) is caused by bremsstrahlung from interactions between plasma electrons and pellet particles. Immediately after pellet ablation, the

density profile becomes very hollow and the temperature drops, leading to decreased X-ray emission. However, the most striking effect is the observation of a snake-like perturbation superimposed on a symmetric emission profile. The observations also show that the snake is due to the rotation of a small region with enhanced X-ray emission.

The poloidal dimension (FWHM) of the snake is typically $\lambda_{\theta} \approx 25\text{cm}$ and is calculated from the transit time across the field of view of either a central X-ray channel or the interferometer. The radial dimension, typically $\lambda_r \approx 17\text{cm}$, is calculated from the relative intensities of the snake viewed radially or poloidally. A set of typical parameters is given in table 1.

In addition to the snake, a transient $m=1, n=1$ sinusoidal MHD oscillation is usually seen just after pellet injection. This effect has also been observed on other machines ⁶⁾.

Plasma parameters in the snake region - The temperature and density in the snake region are determined with an ECE polychrometer, a multichannel far infrared interferometer and a 2 mm microwave transmission interferometer (figure 2) with typical values given in table 1. The very large density perturbations are calculated from the line integral measurements of density using the dimensions of the snake region determined by the X-ray and line density measurements. The density within the snake can be up to twice that of the surrounding plasma although the total number of particles in

the snake is only $\approx 1\%$ of the injected pellet particles. It is also observed that the snake can survive the substantial changes in density profiles, from hollow to peaked, which take place in the time ($\approx 100\text{ms}$) immediately after pellet ablation. In the snake region, the temperature drop, ΔT_e , is always much smaller than the increase in density, implying locally increased pressure. The ΔT_e gradually reduces after $\sim 100\text{ ms}$ to an undetectable level less than 100 eV , although the density increase, Δn_e , is unchanged.

The phase difference between the interferometer and the X-ray camera signals together with the X-ray camera measurements show that the topology of the snake is $m=1$, $n=1$.

Relation to rational q-values - The observation that the snake has $m=1$, $n=1$ provides strong evidence that it is on the $q=1$ surface. This conclusion is reinforced by the fact that it is at the sawtooth inversion radius observed on the tomographically reconstructed soft X-ray emission before pellet injection. Also in agreement with this interpretation are the facts that:

(i) The formation of the snake depends sensitively on the location of maximum pellet penetration which must be inside the $q=1$ surface.

(ii) There is a characteristic dip in the H_α -light from the ablating pellet as it crosses the $q=1$ surface. The reduction in ablation occurs on this surface because only particles from a flux tube with $m=1$, $n=1$ are available, rather than the particles over the whole magnetic surface.

A similar effect is also seen associated with the $q=3/2$ surface. The X-ray signal patterns are more complicated, but an $m=3$ pattern with twice the frequency of the simultaneously observed $m=1$ modulation is clearly seen and is well correlated with the signals from the $n=2$ magnetic pick-up coil combination. The radius of the $m=3$ perturbation coincides with the calculated $q=3/2$ radius.

Lifetime of the snake and the position of the $q=1$ surface -

The snake can persist on the $q=1$ surface for times greater than 2 s and can survive several sawtooth crashes. In figure 3 the X-ray emission profiles for two 100 ms time intervals show the persistence of the snake over 800ms and the effect of a sawtooth collapse. The rotation observed in the upper part of the figure, unlike that of MHD oscillations in ohmic plasmas, is in the direction of the plasma current; in the lower part there is no rotation although it restarts at later times. The snake is sometimes destroyed by a soft disruption or a very large sawtooth collapse.

The long lifetime of the snake allows the determination of the position of the $q=1$ surface (and $q=3/2$) during the sawtooth cycle. In figure 4 a substantial inward shift of the snake is seen after a sawtooth collapse with a 40% change in radius. This is followed by a slow outward movement of the snake and, therefore, also the $q=1$ surface.

Discussion - The plasma behaviour in the snake is unexpected and difficult to explain. If the equilibrium were to remain axisymmetric after injection of the pellet, the temperature

and density perturbations would rapidly spread out along the magnetic field lines. Consequently, the perturbations would spread by collisional diffusion over the flux surfaces except for a very narrow region close to the rational surfaces. This spreading of density would take place within tens of milliseconds outside the region $|1-q| < 10^{-2}$. The persistence of the density perturbation for ~ 2 s therefore implies a change in the magnetic topology. Calculations show that the observed temperature perturbation causes a drop in current density, allowing the growth of a magnetic island to the required size during the pellet deposition around the $q=1$ surface. The question then arises as to how the perturbation persists.

If the persistence is assumed to be due to good confinement a limit would be placed by Coulomb collisions. The estimated confinement time in the banana regime would be ~ 0.3 s. This seems to be too short to be consistent with the observed decay rates, which can be greater than several seconds. A further difficulty arises from the toroidal precession of trapped particle orbits. This precession plays no significant role in an axisymmetric plasma but in the case of the snake would lead to loss of particles on a timescale of ~ 10 ms. It is difficult, therefore, to understand the behaviour in terms of an exceptionally good particle confinement. It may be that the observed state is a deformed stationary equilibrium to which the injection of the pellet has allowed access. There is then no need for individual particles to be confined and some form of recycling could be occurring.

Another question that arises is how the magnetic island itself is maintained. One possibility is that the observed depression of the electron temperature reduces the electrical conductivity along the field lines producing a local reduction in current density and forming a magnetic island. The estimated value of $\Delta T_e/T_e$ to maintain a steady structure of the observed size is 10^{-2} to 10^{-1} depending on the somewhat uncertain magnetic shear at $q=1$. The observed value of $\Delta T_e/T_e$ is initially ~ 0.2 but decays on the timescale of ~ 100 ms to a value too small to be detected, that is $\lesssim 10^{-1}$. Thus although the required temperature depression is observed, no firm conclusion can be drawn. An alternative possibility is that the required decrease in the local resistivity is due to a small enhancement of the impurity concentration. This could be due to the electric potential which arises to confine the local deuteron pressure.

Experiments of this sort can clearly give useful information about the q -profile by identifying rational surfaces. For example, an estimate of $q(0)$ can be obtained from the observed decrease in the snake radius, and therefore the $q=1$ radius (r_1), during a sawtooth crash for which typically $\Delta r_1/r_1 \approx \frac{-1}{3}$ (fig.4). Since the calculated change in the current profile due to sawteeth is quite small ($\Delta q \approx 0.02$) in JET⁷⁾, a smooth q -profile would have to be very flat in order to give the large shift in r_1 which the snake behaviour reveals. In fact, if a parabolic q -profile is assumed, then

$q(0) \approx 0.97$ before the sawtooth crash. This is particularly important for the discussion of sawtooth models.

In summary, the snake is an unexpectedly persistent local perturbation of the plasma arising from pellet injection. It appears to be due to the generation of a magnetic island at the $q=1$ surface within which ablated pellet atoms are deposited. It might involve a remarkably good level of particle confinement, but it is more likely that its long duration indicates a change to a new non-axisymmetric equilibrium. These experiments also illustrate the use of pellet injection as a powerful method of probing the q -profile.

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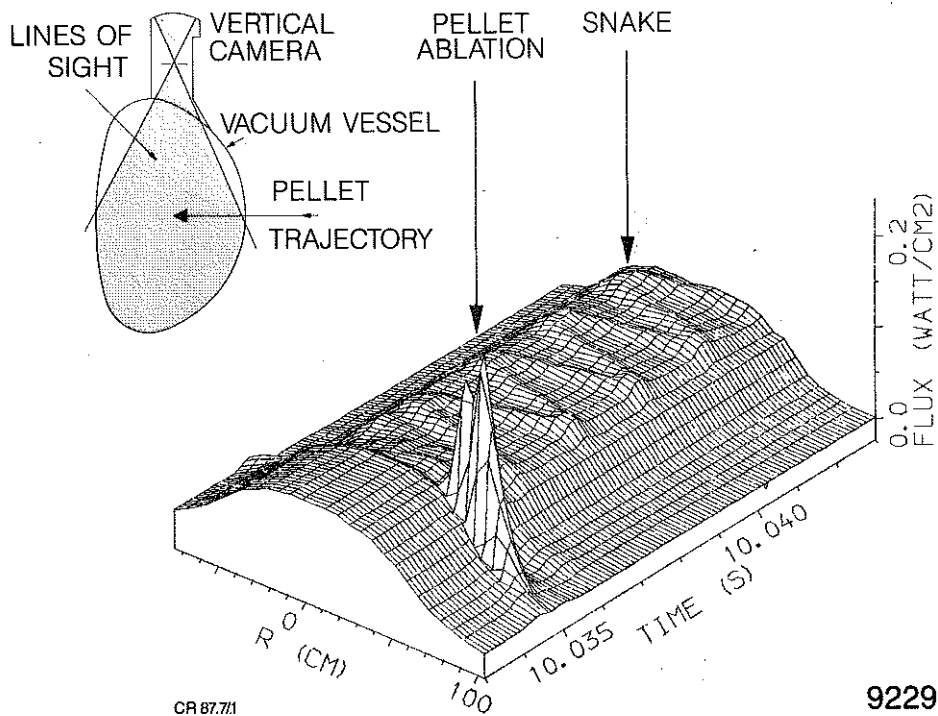
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TABLE 1 - Parameters of typical snakes

Shot number	Time after pellet ablation ms	$\Delta n_e \times 10^{19} \text{m}^{-3}$	$n_e \times 10^{19} \text{m}^{-3}$	$\Delta T_e \text{ eV}$	$T_e \text{ eV}$	$l_\theta \text{ (cm)}$	$l_r \text{ (cm)}$	Duration of snake (s)
9550	13	3.4	4.9	135	1200	24	14	0.85 (a)
9382	240	2.6	6.6	<100	1180	24	17	≥ 1.9 (b)
9378	220	1.5	6.2	<100	1150	31	16.5	≥ 1.9 (b)
9228	10	3.6	4.9	210	650	25	19.5	0.23 (a)

a) terminated by a soft disruption.

b) still present at end of data acquisition.



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9229

Fig.1 Pellet-ablation (10.035 s) and snake oscillation seen by the vertical soft X-ray camera (50 μm Be filter). The X-ray flux is shown as a function of time and detector chord radius (R) measured from the plasma centre.

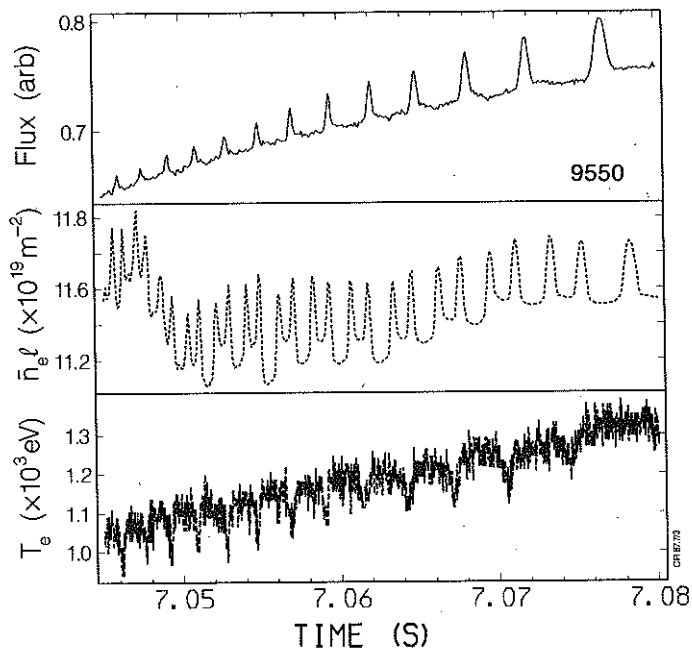


Fig.2 Time correlation of X-ray flux, line-density and T_e at the snake radius. (The soft X-ray channel views close to the snake radius and therefore sees only 1 peak per turn). The different phases of the signals are due to different measuring locations.

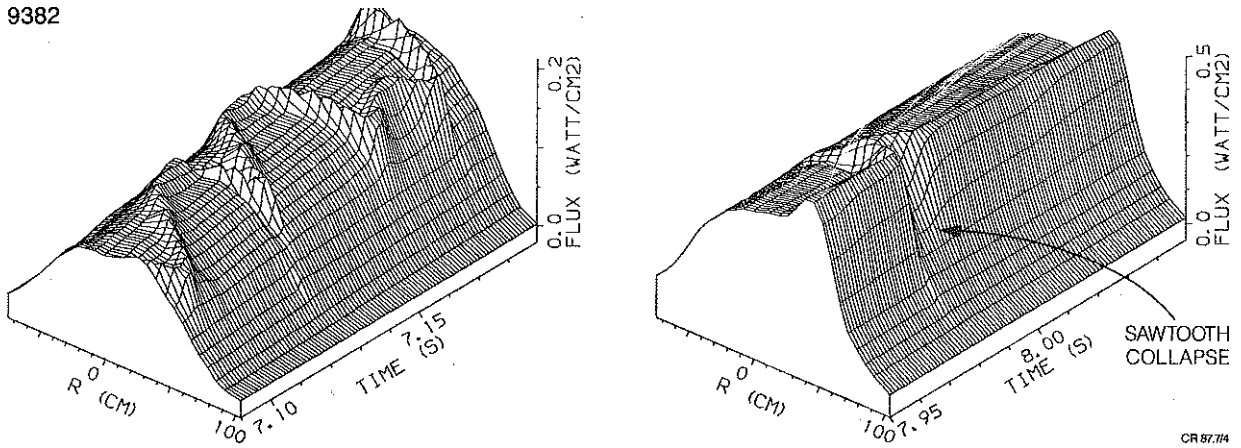


Fig.3 X-ray signals from the vertical camera showing a long lasting snake oscillation, which is locked in the lower figure. The relative amplitude of the snake is unchanged by the sawtooth collapse.

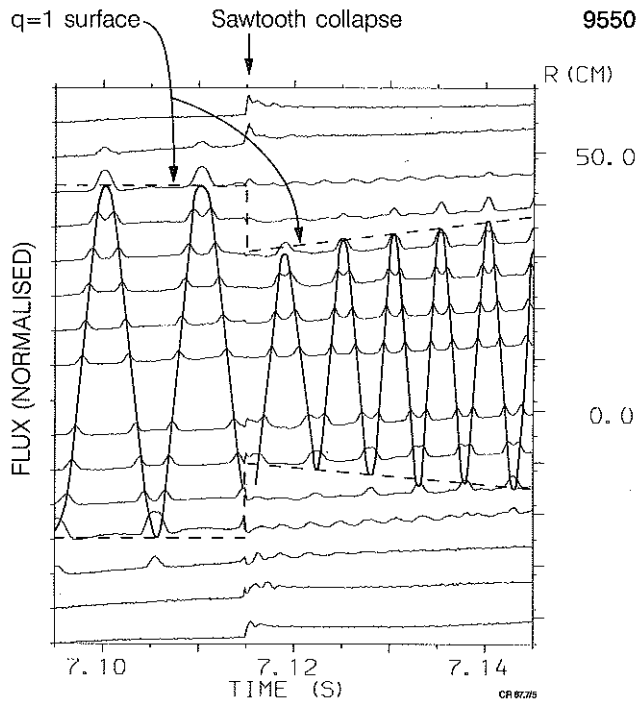


Fig.4 X-ray flux plot for the vertical camera signals showing the inward shift of the snake during a sawtooth collapse. The full line follows the point of maximum emission and the dotted line shows the inferred radius of the q=1 surface.