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## Fast Wave Current Drive experiments in JET divertor plasmas

F. Nguyen<sup>1</sup>, L.-G. Eriksson, B. Fischer, C. Gormezano,  
E. Righi, F. Rimini, G. Sadler, A. Sips, D. Start  
JET Joint Undertaking, Abingdon, Oxon, OX14 3EA, UK.

<sup>1</sup> Permanent address: CEA, Centre d'Etudes de Cadarache, France.

### Abstract:

As part of the JET experimental campaign in the pumped divertor configuration with the new ICRF A2 antennae [1], studies of Fast Wave Heating and Current Drive scenarios have been carried out [2]. For the experiments described in this paper, the scenario with 3rd harmonic D cyclotron layer at centre and both fundamental and 2nd harmonic H cyclotron layers at the edge has been chosen. The experiments have been carried out in 2 MA D plasmas and 2.2 T on axis with  $f_{\text{ICRF}}=51.4$  MHz and a target temperature of  $T_e(0)\sim 1.5\text{-}2$  keV. Significant Fast Wave Heating has been observed and studied for varying plasma parameters ( $B_T$ , ICRF power, antenna phasing, density). There is evidence of both Fast Wave Direct Electron Heating and heating through 3rd harmonic D cyclotron absorption. These two effects disappear when the 2nd harmonic H cyclotron layer enters in the plasma on the high field side. The power partition between these two heating channels is difficult to assess experimentally. The results are compared with numerical simulations of the power deposition (PION, SINGLE). Furthermore, a prediction of the current drive efficiency is given (ALCYON). This scenario provides efficient heating, since the H-mode is triggered when the electron density and the ICRF power in  $0\pi\pi 0$  phasing are increased and a RF only record DD reaction rate of  $1.8 \cdot 10^{16} \text{ s}^{-1}$  has been achieved.  $0\pi 0\pi$  phasing gives similar results.

### Fast Wave Electron Heating and Current Drive scenario:

The target plasma parameters are  $B_T=2.2$  T,  $I=2$  MA and  $T_e(0)\sim 1.5\text{-}2$  keV in Deuterium. At ICRF frequency of 51.4 MHz, 3rd harmonic D cyclotron layer is at centre while both fundamental H cyclotron resonance (also 2D and 2  $^4\text{He}$ ) and 2nd harmonic H (also 4D and 4  $^4\text{He}$ ) cyclotron layers are marginally present at the edge (e.g. # 35320, fig. 1). Unless otherwise specified, the phasing is  $0\pi\pi 0$ . A similar scenario has been experimented on Tore Supra [3].

### Experiment:

A heating with sawteeth is observed but is sensitive to H resonance layer position. A scan of the magnetic toroidal field has been done (# 35520). When the H cyclotron 2nd

harmonic layer enters plasma on the low field side (fig. 1), the density, the radiated power and neutral H flux (Neutral Particle Analyser horizontal line) increase while the DD reaction rate, high energy neutral D flux (NPA horizontal line) and  $T_e$  drop. There is evidence of Fast Wave Electron Heating. The analysis of the fast switch off of  $P_{ICRF}$  and of the sawteeth slope [4] yields peaked deposition profiles (# 35319, fig. 2a, b). One finds  $P_{FWEH}=0.9$  MW for an ICRF launched power of 6.5 MW which means 14% of damping on electrons in  $0\pi\pi0$  phasing (fig. 2b). The power directly coupled to electrons in  $+90^\circ$  and  $-90^\circ$  phasing is estimated to be less than 50% of  $0\pi\pi0$  heating [5]. The pre-heating with LH power produces no visible difference. Two pulses in  $^4\text{He}$  have been performed. The second one (# 35328) is approximately 50%  $^4\text{He}$  and 50% D and the results are similar to D plasma (# 35325). In particular, the  $T_e$  profiles are identical. Part of the ICRF power is clearly absorbed at the 3rd harmonic D cyclotron resonance as indicated by a high DD rate and  $\gamma$  emission from the plasma centre (3.1 MeV from  $^{12}\text{C}(d, p)^{13}\text{C}$  with  $E_d > 1.8$  MeV). Fast D is detected with the vertical high energy NPA up to the highest energy channel available ( $E=1.1$  MeV) (fig. 3). This contrasts with Tore Supra results that show dominant electron heating and no fast ions in  $^4\text{He}$  plasma [3]. This is likely to be due to the better ion confinement in JET that allows the build up of a hot ion tail but models have to be developed. A power and density scans have been done. When the density and ICRF power increase, the plasma evolves towards better performance (fig. 4): the diamagnetic energy content and the DD reaction rate increase and one finally obtains an elmy H-mode. The H factor was 1.6 (ITER89-P) and 0.8 for H93 (H-mode scaling with subtraction of the energy of fast particles) for pulse # 35525 with  $0\pi\pi0$  phasing (fig. 3) (similar results were obtained for # 35526 with  $0\pi0\pi$  phasing). For this pulse, the record DD reaction rate for ICRF only ( $1.8 \cdot 10^{16} \text{ s}^{-1}$ ) was achieved with  $P_{ICRF}=12.9$  MW. The other parameters are  $W_{DIA}=3.2$  MJ (with  $W_{\text{fast particles}}=0.6$  MJ),  $\langle n_e \rangle = 3.9 \cdot 10^{19} \text{ m}^{-3}$ ,  $T_e(0)=5.7$  keV,  $P_{\text{total}}=13.5$  MW. This performance is comparable to hot ion mode (DD reaction rate of  $2 \cdot 10^{16} \text{ s}^{-1}$  for  $P_{NBI} \sim 14.6$  MW at 1.75 MA, # 34489).

#### Numerical simulations (based on the ICRF and plasma parameters of pulse # 35320):

The power partition between Fast Wave Direct Electrons Heating and heating through D cyclotron absorption at the 3rd harmonic is difficult to estimate. Single Pass Absorption (SPA) estimate (SINGLE code [6], (fig. 5)) is larger for electron (4.5 keV) compared to D absorption of the bulk (2.7 keV) but the SPA for a population of fast D (1% of the D bulk) becomes predominant above 300 keV at the most favourable  $k_{\parallel}$  for electron damping ( $N=33$ ,  $k_{\parallel}=11 \text{ m}^{-1}$  on axis). The PION code simulation (self consistent treatment of the D tail formation and the Fast Wave absorption [7]) with hot D confined up to an energy of 2.5 MeV indicates equipartition of the coupled ICRF power between electron and deuterium. The ALCYON (2D full wave code [8]) simulation provides almost identical power deposition profile on electrons

for  $90^\circ$ ,  $0\pi\pi 0$  and  $0\pi 0\pi$  phasing and predicts a current drive efficiency of  $0.018 \cdot 10^{20} \text{ A/W/m}^2$  in  $90^\circ$  phasing.

### Conclusion:

Significant Fast Wave Direct Electron heating ( $\sim 14\%$ ) and heating through D absorption at cyclotron 3rd harmonic have been observed. This contrasts with the Tore Supra results that show dominant electron heating and no fast ions with the same scenario [3]. The power partition is difficult to assess. The competition with lower cyclotron harmonic (H or D) should be avoided (edge absorption). The Fast Wave Current Drive efficiency is estimated to be  $0.018 \cdot 10^{20} \text{ A/W/m}^2$  for  $90^\circ$  phasing with the plasma parameters of pulse # 35320 which is comparable to the Tore Supra results [3]. Elmy H-mode is triggered when density and ICRF power in  $0\pi\pi 0$  phasing are increased and a RF only record DD rate of  $1.8 \cdot 10^{16} \text{ s}^{-1}$  has been achieved. Similar effect is observed in  $0\pi 0\pi$  phasing. An efficient heating scenario at 2 MA, 2.2 T with no minority ion population has been found, which can start on low  $T_e$  (1-2 keV) target plasma and might provide suitable conditions for very long pulse (60 s) operation for advanced tokamak scenarii on JET.

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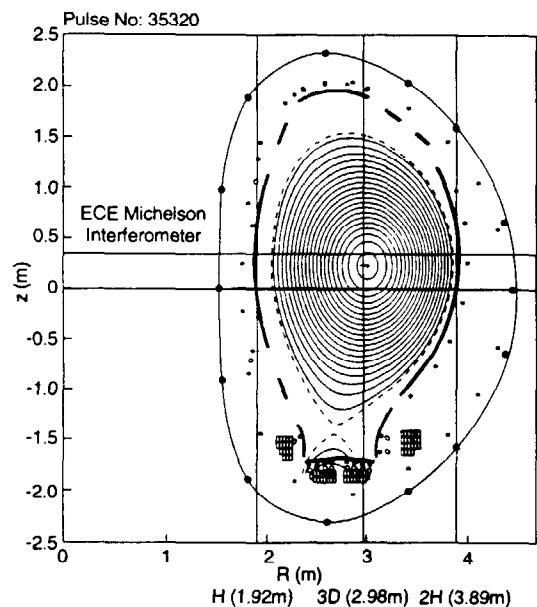


fig. 1: location of the different ion cyclotron layers on a poloidal view of the plasma.

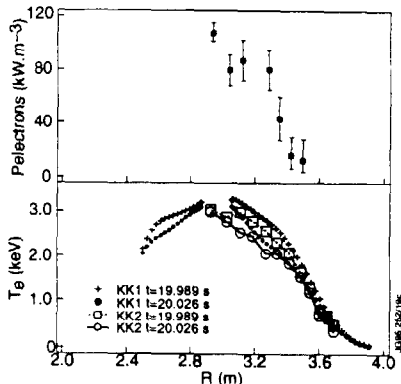
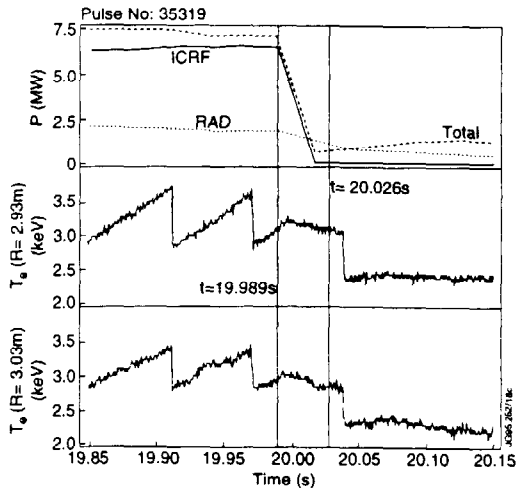


Fig. 2: pulse n° 35319,  $0\pi\pi 0$  phasing, 51.4 MHz, 2.18 T, Deuterium, 1.8 MA, a: PICRF sharp switch off, effect on the sawteeth, ECE Grating Polychromator ( $z=25\text{cm}$ ), b: direct electron heating profile c: ECE  $T_e$  profiles in the plane  $z=z(\text{axe})$ .

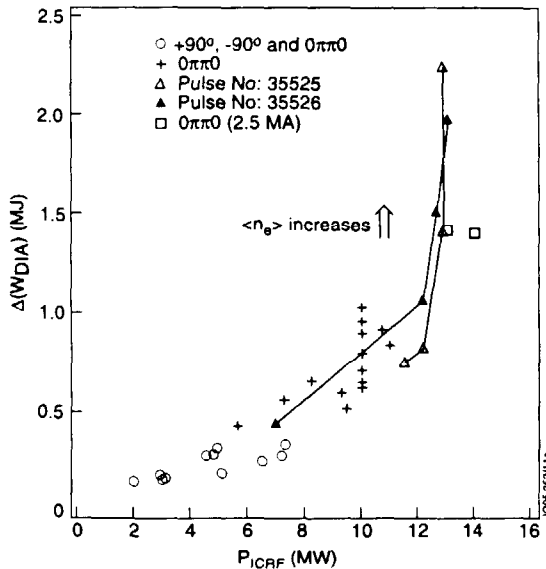


Fig. 4 diamagnetic energy increase in function of the ICRF power, 51.4 MHz, 2 MA, 2.18 T.

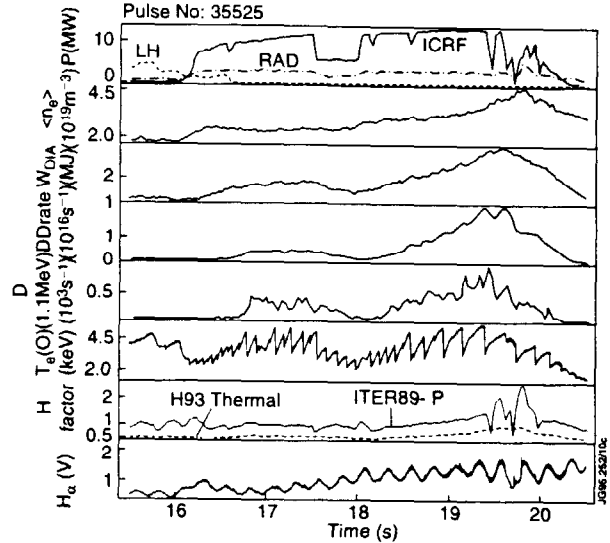


Fig. 3: pulse n° 35525  $0\pi\pi 0$  phasing, 51.4 MHz, 2 MA, 2.18 T, Deuterium.

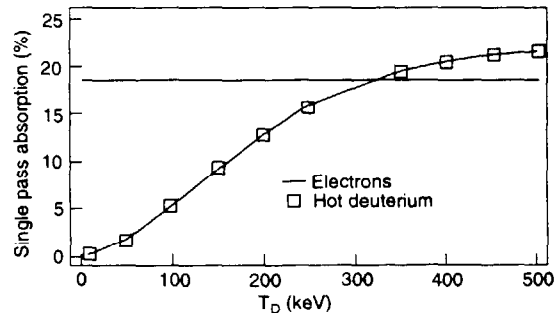
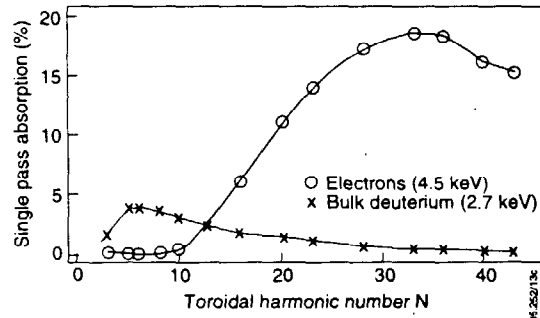


Fig. 5: SINGLE simulation of the pulse n° 35320 at  $t=16.2\text{ s}$ , a: SPA on electrons (TMP and ELD) and on bulk D (3rd harmonic) in function of the toroidal harmonic number  $N$ , b: SPA for the toroidal harmonic number  $N=33$  ( $k_{||}=11\text{ m}^{-1}$  on axis) on electrons and hot D (1% of the bulk in density) in function of  $T_D$ .