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ABSTRACT

ITER operation might be limited by the tritium inventory that will be retained in carbon tiles and re-deposited carbon layers. Means to determine in-situ the tritium content of these layers and also the deuterium content in the non activated phase of ITER and tritium recovering techniques have to be developed. This paper describes an approach in which the thickness and deuterium (tritium) content of re-deposited carbon layers is determined by local laser heating. This leads to the release of the retained fuel by evaporation and, depending on the laser power density even to the removal of the layer itself. An essential element of this method is the quantitative determination of the released particles by local spectroscopy in the edge plasma. Results from this method from experiments in the laboratory, and on TEXTOR and JET are described in this paper.

1. INTRODUCTION

The present choice of first wall materials in ITER foresees beryllium at the main chamber, tungsten at the baffles and graphite CFC at the divertor plates. In particular the use of graphite can lead to accumulation of tritium in hydrocarbon layers and is a serious concern for ITER operation that in total allows only an amount of 350g Tritium inside the vessel. These layers are formed by co-deposition of eroded wall materials with deuterium and tritium and develop in net deposition areas, such as at the inner divertor area or the slits of the castellated tiles [1], [2], [3]. An estimation of the amount of retained T in ITER is highly uncertain because the prediction relies upon the knowledge of many parameters, such as the main chamber material erosion, the flows in the plasma edge, and the local transport of material inside the divertor. A diagnostic is urgently needed that monitors in-situ the T retention at different representative locations and, in combination with other methods such as a global gas balance, delivers reliable data of the in vessel T content in ITER. This is of particular importance in the non activated phase of ITER which should provide the information on this topic for the activated T-phase.

It is well known that a hydrocarbon layer releases the stored fuel gas upon heating [4]. Thus, laser radiation is proposed to heat any target surface. The temperature increase is mainly determined by the heat capacity and adhesion of the layer to the bulk material. However, the heat absorption can also be influenced by plasma that is created by interaction of released particles with high power densities of the laser radiation.

Experiments with laser heating have been performed already at JET and TEXTOR. At JET the Edge LIDAR beam always hits a carbon tile at the inner divertor with power densities of about 10GW/cm^2 during 300ps pulses. This generates ablation of deposited amorphous hydrocarbon layers (a-C:H) but might also damage the bulk material. The released D and C atoms were observed spectroscopically in the edge plasma [5]. In order to avoid the removal of bulk material a ruby laser pulse with comparable energy but much lower power density (40kW/cm^2 for 300μs) was fired onto a plasma exposed C test limiter at TEXTOR between discharges. The desorbed deuterium was detected with a **Q**uadrupole **M**ass **A**nalyser (QMA) [6]. Because of the large volume of ITER the sensitivity for the detection of the released particles might not be sensitive enough.

Therefore the application of emission spectroscopy in combination with laser induced desorption might be another possibility for an in-situ diagnostic for the determination of the tritium content in co-deposited layers. The ìlongî laser heating leads to the release of D (T) atoms and molecules. The evaporated particles enter the plasma and their D_α (T_α) radiation is observed. With known plasma parameters the emission light can be converted to absolute particle fluxes using appropriate conversion factors [7], [8].

In this paper results are presented where "free running" ruby lasers are fired onto graphite targets with hydrocarbon layers. The process of laser induced desorption on welldefined hydrocarbon layers is investigated in a laboratory experiment. The spectroscopic detection of laser desorbed particles in the plasma edge and conversion into total fluxes is studied at carbon test limiters in TEXTOR, results of which are presented in the 2nd section. For the same purpose, the Edge LIDAR at JET system was changed to long pulse operation. The evaluation of the D_α emission signals of desorbed D in 3 consecutive laser pulses (1s delay) was measured during the limiter and divertor phase. This is described in the last section.

2. LABORATORY EXPERIMENT

The target is located at the centre of an UHV chamber with 300mm diameter. The base pressure generated by a turbomolecular pump is below 10^{-6}mb. A target manipulator allows a rotation along the symmetry axis and a vertical and horizontal movement over the whole area of a 5x5 cm^2 target. The radiation of a ruby laser is focussed onto the target. The spot size and correspondingly the power density can be varied from 1 to 10mm diameter by retraction of the focussing lens. A further variation of the beam power without changing the spot size is performed by the introduction of neutral filters into the beam line. The maximum laser energy is about 1.5J for pulse duration of 350µs. The laser power is monitored by a fast diode.

A QMA and an IR array camera are attached at the equatorial plane of the vacuum chamber each under 45^o in respect to the laser beam. The distance of the QMA to the target is 0.39m and the time resolution is about 1µs. For the investigation of a-C:H layers deposited on carbon the atomic masses 2 (D), 4 (D_2), 24 (C_2) and 36 (C_3) were selected at the QMA. A 128 element linear array camera was operated with 100kHz frame rate and 36µm/pixel spatial resolution. Only IR light above 850nm was selected by an optical filter and the intensities were compared with Planck's law to determine the surface temperature. The lowest detectable temperature at the highest time resolution was 700^oC.

The a-C:H layer on a graphite target was produced in a DC glow discharge with a bias of 600V [9]. The deuterium content of this layer with a thickness of 141nm is 0.31D/C that corresponds to $1.7 \cdot 10^{17}$ D atoms/cm^2 and was used for calibration of the QMA signals. The ratio of released atomic and molecular D might depend on the surface temperature [10] that yet has not be considered for the calibration of the QMA signals.

The time traces of a single laser pulse for power density, absorbed energy and surface temperature rise at the centre of the exposed area ($1.1 \cdot 0.036mm^2$) are presented in Fig.1. A temperature of

about 700^{o}C is reached after 120μs that is in reasonable agreement with calculations made for carbon, assuming a spiking laser radiation with 100 kHz repetition rate and 1μs spike length. With the experimental data of 0.3J energy over 350μs focussed on 2mm^2 a power density of 430kW/cm^2 is found in a single spike. To a certain extent temperature rises are coupled to the spike structure of the ruby laser.

A nearly complete desorption of deuterium out of the a-C:H layer in a single laser shot is shown in Fig.2. Two consecutive laser shots with energy of 1.5J each were focussed onto the same spot of 3mm^2. The target was at room temperature. The QMA signal of the second laser shot is significantly lower than in the first shot. In further laser pulses with same energy the C_2 QMA signal was observed that remained nearly constant at values obtained from a pure carbon surface. This behaviour can be explained by carbon evaporation induced by the high power densities (about 1MW/cm^2) in single spikes.

With reduced laser energy of 0.2J about 10 laser shots were necessary to reach the same reduction of the deuterium signal as in the previous experiment. Even with doubled laser energies the release of D_2 (except in the first laser shot) could not increased significantly as seen in Fig.3. The strong signals obtained in laser shot 1 and 17 might be related to adsorption of water. The C_2 signal measured at the end of this series was more than a factor of 10 lower compared to the previous experiment. This behaviour supports the assumption that only carbon from the amorphous hydrocarbon layers is removed at this power densities while the power of a single spike is not sufficient to cause sublimation of the bulk carbon. The temperature integrated over 10μs never increased above 7500C in this case.

3. TEXTOR EXPERIMENTS

The experimental set-up of the laser desorption system at TEXTOR is shown in Fig. 4. A carbon limiter was inserted into a limiter lock system on top of the vessel and the tip was positioned at the Last Closed Flux Surface (LCFS) at a radial plasma position of 46cm. The laser beam with energies up to 5J entered the vessel from the bottom. Results presented in this paper were obtained from a laser spot 1cm in the shadow of the LCFS located 35mm from tip of the limiter on the electron drift side. A polychromator equipped with interference filters detected simultaneously the CD, D_α, D_β, D_γ and CII light in front of the limiter surface with a time resolution of 5μs. Additionally a spectrum of the plasma background was measured with a spectrometer. The initial temperature distribution at the limiter surface was monitored with a CCD camera collecting the IR light above 850nm. The fast temperature rise averaged over the laser-exposed area was measured by a pyrometer with 15μs time resolution. The edge plasma parameters needed for the determinations of the conversion factors were measured with He beam diagnostic [11]. Also the local density were deduced from the ratio of D_α/D_β. A fast photodiode recorded the laser power.

In Fig. 5 time traces of $H_\beta(D_\beta)$ and C*II* lines are presented which were obtained by laser induced desorption in Ohmic heated TEXTOR discharges (I_p = 350kA, $n_e = 2 \cdot 10^{19}/m^3$; B_t = 2.25T) in deuterium. The total fluxes for H/D and C given in Fig. 6 were deduced from H_β and C*II* line

intensities using the appropriate conversion factors (S/XB) of 67.5 and 6.97, respectively as given in the literature [7], [8] for the measured electron density and temperature. About (6-10) $\cdot 10^{15}$ /cm^2 H(D) atoms and 1 $\cdot 10^{15}$ /cm^2 C atoms are released from a spot of 8mm diameter with an initial temperature of 450^{o}C (Fig.6(a)). The fast pyrometer shows a temperature rise of about 300K at the surface during laser radiation. The amount of released carbon atoms was strongly dependent on the initial temperature of the test limiter before laser exposure. The carbon signals rose by a factor of 10 when the initial surface temperature was increased from 450^{o}C to 770^{o}C (Fig.6(b)). This effect indicates that the spiking power intensity of the ruby laser radiation surpasses the sublimation temperature of carbon for the later case.

4. JET EXPERIMENTS

The edge LIDAR system [12] without Q-switches delivered an average power density of about 10 kW/cm^2 for about 350µs. The beam hitted the upper inner divertor tile above the strike point of the plasma (see [5]). A series of 3 laser pulses with 1Hz repetition rate was fired onto the same spot during the limiter and divertor phases. The D_α signal of the desorbed particles was measured with a fast wide-angle observation system (250kHz sampling rate). Laser induced desorption measurements were applied for a variety of different plasma conditions. The D_α signals as shown in Fig.7(a) showed always a significant drop after the first laser shot and a weaker decay for the next 2 laser shots. In series of consecutive discharges the amplitude of D_α signals varied within about a factor of 3. This might be explained mainly by changes of plasma parameter but more analysis is needed here. This behaviour although less pronounced was also found if the laser was fired during limiter mode for which the D_α signals were much lower (Fig.7(b)). However, to determine the total amount of desorbed D, the higher S/XB factor for lower electron densities and temperatures and the lower amount of particles exited in the observation volume must be taken into account.

The retention of deuterium on these areas occurs mainly during the divertor phase since only a very weak signal was observed in the first shot during the limiter phase if the laser was fired once before the discharge. However laser induced desorption during the divertor phase did not show a clear influence on any change of discharge conditions. An explanation for this behaviour can be the high particle flux to the divertor tile that leads to saturation within a few seconds or that the laser power is too weak to desorb all the deuterium form deeper zones. The measured deuterium signal would thus only correspond to a shallow surface layer.

CONCLUSION

The content of deuterium and carbon in co-deposited hydrocarbon layers has been measured in a laboratory experiment by means of laser desorption using “free running” ruby lasers with pulse durations of 350µs. Released hydrogen and carbon was measured using a line of sight QMA system. In TEXTOR and JET the desorbed particles were analysed by insitu emission spectroscopy. At

lower average power densities (80kW/cm^2) sublimation of carbon during the laser spikes (about 1μs) is avoided. The threshold for the onset of carbon sublimation decreases with increasing initial temperature of the target. Laser desorption on a-C:H layers produced in the laboratory shows that more than 80% of the deuterium was desorbed in a single laser shot. In TEXTOR, the amount of deuterium and carbon released by the laser pulse from limiter was measured by quantitative emission spectroscopy on a shot by shot basis. This amount is understood in terms of the retention of fuel in a shallow surface layer due to direct implantation from the edge plasma, while a fraction of fuel is desorbed also in form of hydrocarbons, causing the carbon light emission. In JET, the deuterium release from an inner divertor tile has been measured in a similar way, but without absolute calibration so far. The fuel retention at this location occurs mainly during the divertor operation phase. No clear dependence of the amount of desorbed fuel on the number of accumulated discharges before the laser pulse could be found. This indicates that the laser power in JET is insufficient to desorb the fuel from thicker co-deposits. For application at ITER further improvements of this diagnostic are necessary. A constant power density distribution over the beam cross-section as described already in the literature [13] could reduce the uncertainty of the measurements. A more challenging task is the development of a ruby laser with constant temporal output to minimize the heating by the laser spikes during the laser pulse [14].

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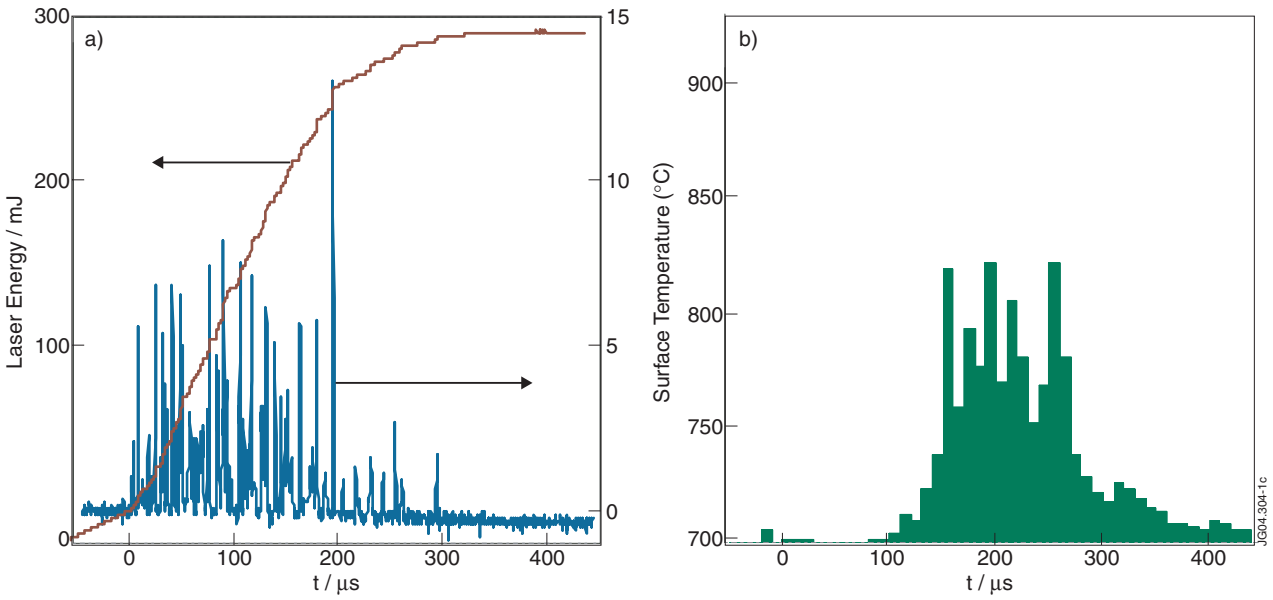

Figure 1: temporal development of the a) laser power and laser energy , b) surface temperature at the spot centre.

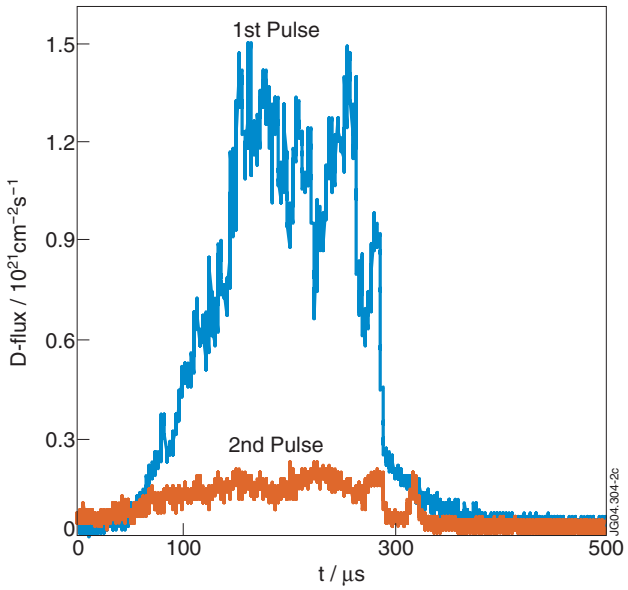

Figure 2: time traces of D_2 QMA signals for two consecutive laser pulses (average power density 150kW/cm^2).

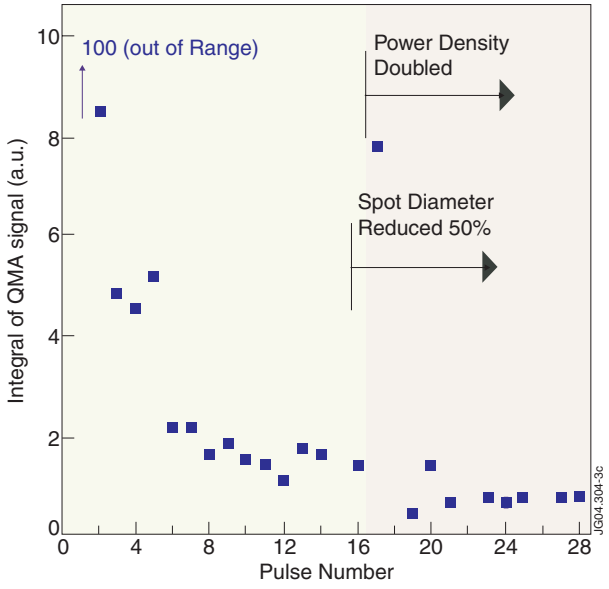

Figure 3: Decay of desorbed D_2 (integrated QMA signal) for consecutive laser pulses (average power density 20kW/cm^2).

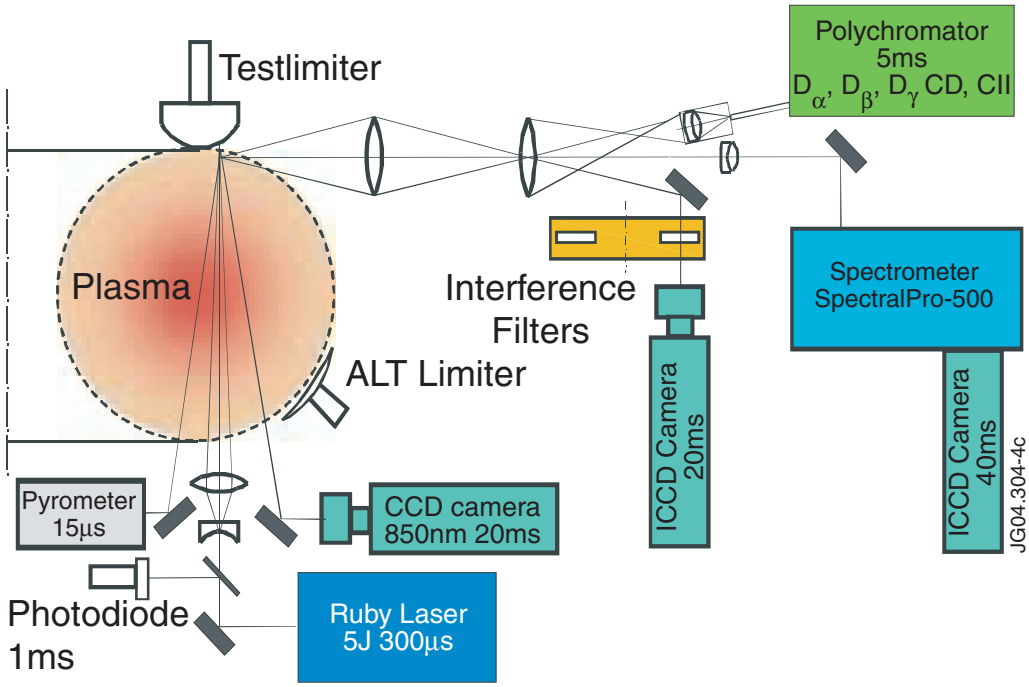

Figure 4: Experimental set-up for laser desorption on test limiters at TEXTOR.

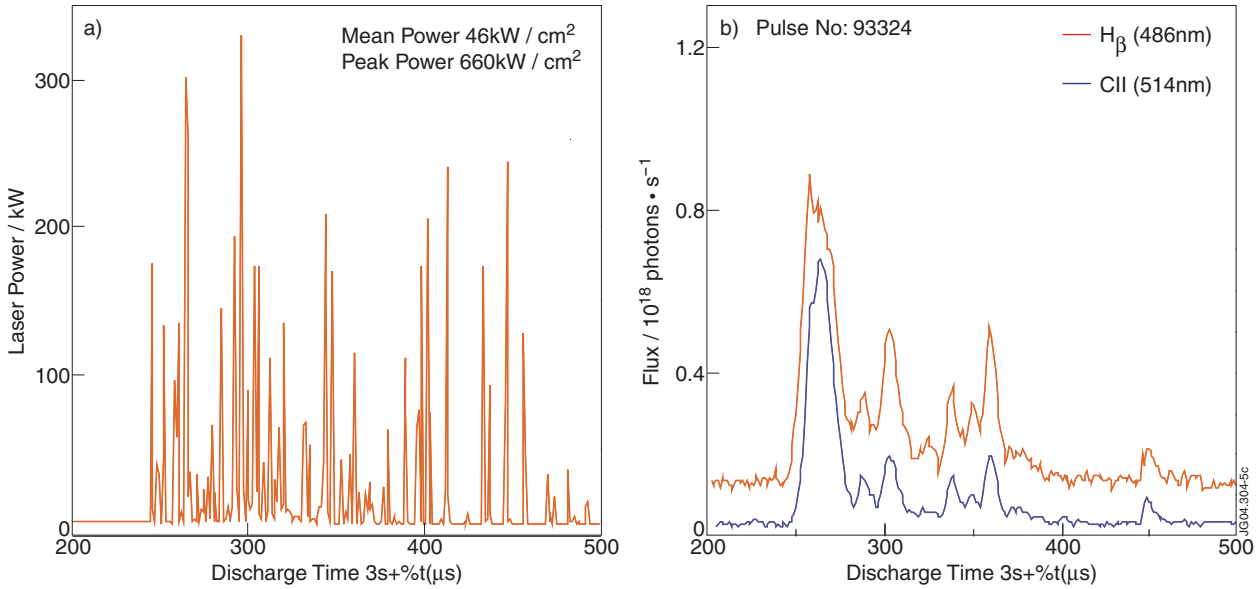

Figure 5: a) temporal behaviour of the ruby laser radiation and b) the corresponding line intensities (H_β and CII) of desorbed particles in front of a carbon test limiter during a TEXTOR discharge.

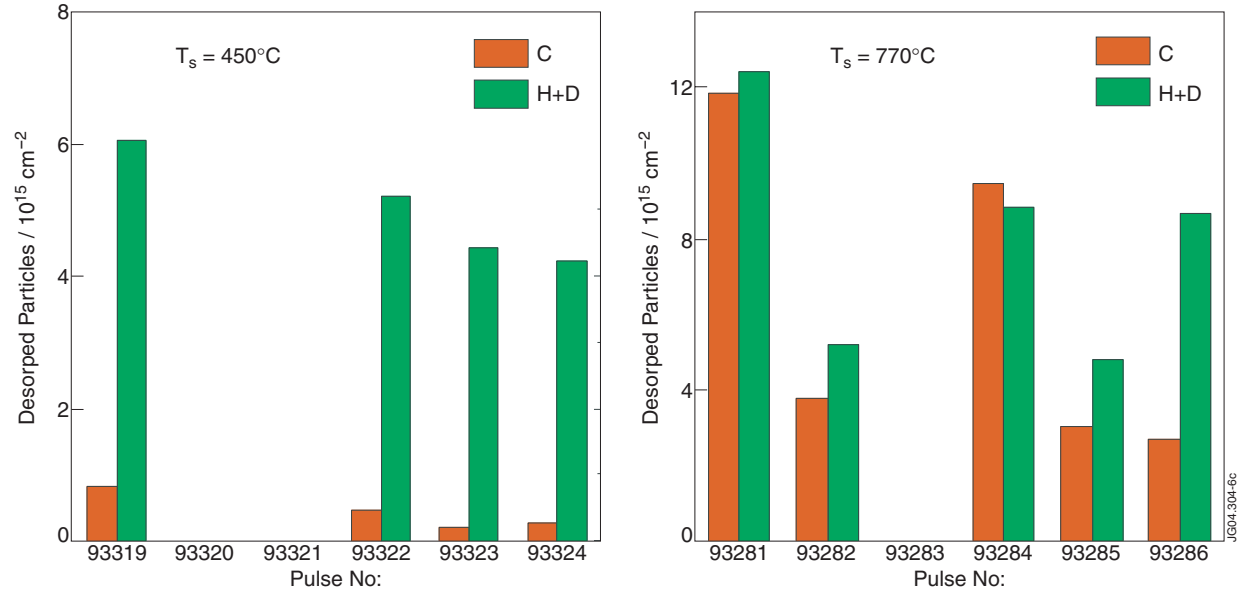

Figure 6: total amount of D and C released from a carbon test limiter at TEXTOR at different limiter temperatures (plasma parameter $n_e = 2 \cdot 10^{18}/cm^3$, Te = 50eV).

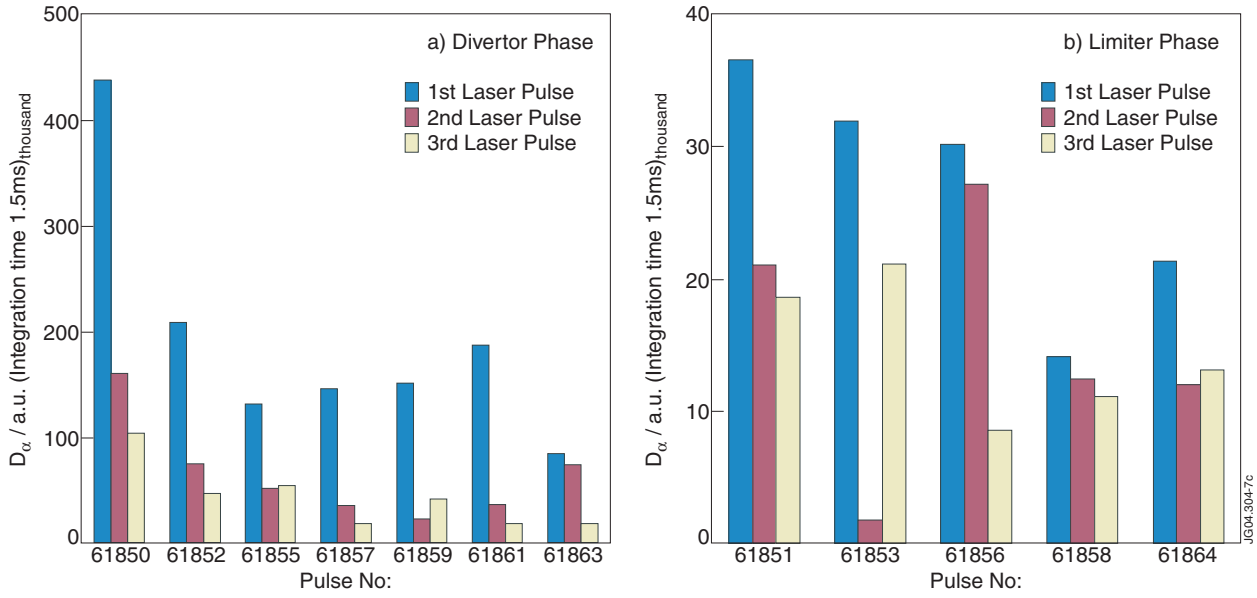

Figure 7: Integrated D_α line intensity for 3 consecutive laser pulses during the a) divertor and b) limiter phase.