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Key Quantities for ITB Formation and Sustainment

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ABSTRACT.

An international ITB database has been constructed to address primarily power threshold scaling, confinement scaling and underlying physics in ITB plasmas. For the power threshold scaling, a precise definition of ITB is required. Either the temporal separation of temperatures in two neighboring measurement channels or the discontinuity in gradients in temperature profiles is often used to judge the ITB formation. However, both of these criteria become unclear for low heating power approaching the threshold for the ITB formation. Given this uncertainty in determining ITB existence, parameters are sought which will better define the time and spatial location of ITB. A parameter frequently employed is R_{geo}/L_T where R_{geo} is the plasma major radius and L_T is the local temperature gradient scale length. Another parameter is $\rho_T^* = \rho_S / L_T$ where ρ_S is the local ion gyroradius at the sound speed; $\rho_S = (m_i T_e)^{1/2} / (Z_i e B_t)$, m_i is the ion mass, Z_i is the ion charge number and B_t is the toroidal field [1]. This parameter has a strong B_t dependence and no R_{geo} dependence in contrast to R_{geo}/L_T . The existence of ITBs in JET plasmas can be inferred in regions of space-time where $\rho_T^* > \rho_{\text{ITB}}^*$, where ρ_{ITB}^* is some critical value. Interestingly, a constant value of ρ_{ITB}^* , 0.014, is valid for a large variety of JET discharges heated by NB and IC with $B_t = 1.8\text{-}4.0$ T when L_{T_e} is used as L_T . As a first step in finding a parameter or a key quantity that can be generally used to judge the ITB existence, the validity of ρ_T^* is investigated for various devices in this paper.

1. ANALYSIS METHOD

Profile data from the following devices are used in this study; Alcator C-MOD (C-MOD), ASDEX Upgrade (AUG), CHS, DIII-D, FTU, JET, JT-60U, LHD, MAST, T-10, TCV and Tore Supra (TS). Here CHS and LHD are heliotron-type helical devices and the others are tokamaks. The data cover a wide region of $R_{\text{geo}} = 0.69\text{-}3.75\text{m}$ (a factor of 5.4) and $B_t = 0.53\text{-}5.28\text{T}$ (a factor of 10). To fix critical values for the ITB formation precisely, it is most efficient to use the values at the location and time of ITB emergence determined from the time evolution of profiles. In this study, to find critical values using a limited number of time slices, values at the boundary of ITB in radial profiles are used. Values of R_{geo}/L_{T_e} (R_{geo}/L_{T_i}) and ρ_S/L_{T_e} (ρ_i/L_{T_i}) are evaluated at the foot of T_e -ITB (T_i -ITB) and just outside the ITB. Here ρ_i is a gyroradius of thermal ion; $\rho_i = (m_i T_i)^{1/2} / (Z_i e B_t)$. The temperature scale length L_T is evaluated using the temperature gradient along the major radius, namely, by $L_T = T / (dT/dR)$. When the temperature profiles are given as a function of normalized radius ρ , dT/dR was evaluated by $dT/dR = (1/a) dT/d\rho$, where a is the horizontal minor radius. The location of ITB foot is defined as the point where $d/d\rho (R_{\text{geo}}/L_T)$ has its maximum, though other definition is also used in some cases where the maximum of $d/d\rho (R_{\text{geo}}/L_T)$ cannot be clearly determined. An example is shown in Fig.1. Since $\rho_i/L_{T_i} \propto T_i^{1/2} R_{\text{geo}}/L_{T_i}$, ρ_i/L_{T_i} tends to have larger values inside the ITB than at the edge, which is a favorable nature to define the region of ITB in a radial space. Our purpose is to find critical values of R_{geo}/L_{T_i} and ρ_i/L_{T_i} , $(R_{\text{geo}}/L_{T_i})^{\text{crit}}$ and $(\rho_i/L_{T_i})^{\text{crit}}$, to define the ITB region. Since both of R_{geo}/L_{T_i} and ρ_i/L_{T_i} change rapidly near the ITB foot, it is not easy to determine critical values that can be applied to a whole discharge duration from a single time

slice data especially in a strong or “box-type”ITB. In this study,critical values are defined as the average of the value at the ITB foot and that outside the ITB with error bars covering these two values as shown in Figs.1(b) and (c).

3. RESULTS

Critical values of $R_{\text{geo}}/L_{\text{Te}}$ and ρ_s/L_{Te} for T_e ITBs are shown in Fig.2. Here,EC heating was employed in TCV, CHS, T-10, JT-60U and LHD, while LH+EC heating in FTU, IC heating in TS and LH or NB+IC or NB+IC+LH in JET. The electron densities are distributed in a wide range; $n_e(0) < 2 \times 10^{19} \text{ m}^{-3}$ in EC heated discharges and JET data in LH phase while $n_e(0) = (5-11) \times 10^{19} \text{ m}^{-3}$ in FTU and TS. The critical values of $R_{\text{geo}}/L_{\text{Te}}$ lie in 5 to 20. In DIII-D, very strong ITB data is used, which resulted in a large critical value in the present analysis; the values ear the lower boundary seem to be valid for the ITB criterion.The dispersion is larger in $(\rho_s/L_{\text{Te}})^{\text{crit}}$ than in $(R_{\text{geo}}/L_{\text{Te}})^{\text{crit}}$. In T-10, FTU and JET (and DIII-D),critical values of ρ_s/L_{Te} are close to the JET criterion, 0.014, while TCV and CHS have larger values and TS, JT-60U and LHD have lower values; ρ_s/L_{Te} reaches 0.03 outside the ITB in TCV.

Figure 3 shows critical values of $R_{\text{geo}}/L_{\text{Ti}}$ and ρ_i/L_{Ti} for T_i ITBs and n_e ITB (C-MOD).I C-MOD,where off-axis IC heating was employed,the peaking of density was remarkable with small changes in the temperature gradient and hence the scale length of pressure L_p was used instead of L_T . All other data are from NB heated plasmas,some of which were also heated by LH,IC and EC in addition to NB.In DIII-D,QDB discharges with positive magnetic shear have low values of $R_{\text{geo}}/L_{\text{Ti}}$ and ρ_i/L_{Ti} , $(R_{\text{geo}}/L_{\text{Ti}})^{\text{crit}} < 5$,than weak/negative shear discharges. In JET and JT-60U, no large differences are observed between reversed shear discharges and weak positive shear discharges (optimized shear in JET and high β_p H- mode in JT-60U). In JT-60U plasmas with low beta ($\beta_N < 1$), low values ρ_i/L_{Ti} , $(\rho_i/L_{\text{Ti}})^{\text{crit}} < 0.005$, are observed. Recently ITBs has been observed in MAST [2] and NSTX [3], in which large values of $(\rho_i/L_{\text{Ti}})^{\text{crit}} > 0.03$ are obtained.

DISCUSSION

Both of ρ_i/L_{Te} (T_e ITB)and ρ_i/L_{Ti} (T_i ITB)have been distributed in a wide range exceeding one order; for T_e -ITB $(\rho_s/L_{\text{Te}})^{\text{crit}} > 0.03$ in TCV and CHS while $(\rho_s/L_{\text{Te}})^{\text{crit}} \sim 0.003$ in TS and for T_i -ITB $(\rho_i/L_{\text{Ti}})^{\text{crit}} > 0.03$ in MAST while $(\rho_i/L_{\text{Ti}})^{\text{crit}} \sim 0.003$ in JT-60U. In Fig.4, ratios of ρ_s and ρ_i to R_{geo} are plotted.It is found that ρ_s/R_{geo} or ρ_i/R_{geo} is large in devices with small R_{geo} , TCV, CHS and MAST, except for high field devices, FTU and C-MOD, while JT-60U and TS plasmas with low foot temperatures have low values due to large radii and high magnetic fields.

In the present study, it is found that the critical values of ρ_s/L_{Ti} or ρ_i/L_{Ti} are not constant but are distributed in a wide range. The variation seems to be mainly caused by changes in ρ_s/R_{geo} or ρ_i/R_{geo} , and the critical values of R_{geo}/L_T vary within a rather smaller range than the critical values of ρ_s/L_{Te} or ρ_i/L_{Ti} . Dependence of $R_{\text{geo}}/L_{\text{Ti}}$ on the ion temperature at the ITB foot, T_i^{foot} , in JT-60U plasmas is shown in Fig.5. $R_{\text{geo}}/L_{\text{Ti}}$ seems to increase with T_i^{foot} for fixed B_i ,which means ρ_i/L_{Ti} changes

more widely than R_{geo}/L_{Ti} since ρ_i at the ITB foot is proportional to $(T_i^{\text{foot}})^{0.5}$. On the other hand, in the JET Bt scan data with q_{95} constant, variation in ρ_s/L_{Te} was smaller than that in R_{geo}/L_{Te} . More work is required to find what parameters determine the critical values of L_T for the ITB formation.

SUMMARY

Critical values of temperature scale length L_T for the ITB formation have been investigated in various devices using the ITPA international ITB Database. Large variations in the ratio of ion gyroradius ρ_s or ρ_i to L_T are found in some cases. It is suggested that the critical values of ρ_s/L_{Te} and ρ_i/L_{Ti} depend on plasma parameters or other quantities that ρ_s (ρ_i) should be introduced to normalize L_T .

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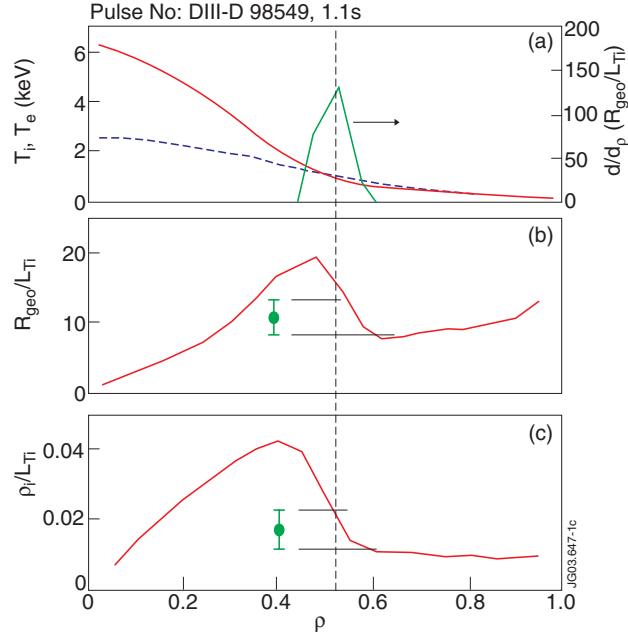


Figure 1: Radial profiles of (a) $T_i, T_e, d/d\rho (R_{geo}/L_{Ti})$, (b) R_{geo}/L_{Ti} and (c) ρ/L_{Ti} in a DIII-D weak shear ELMy H-mode plasma with T_i ITB. The location of ITB "foot" is denoted by a vertical dotted line. The circles with error bars are "critical" values.

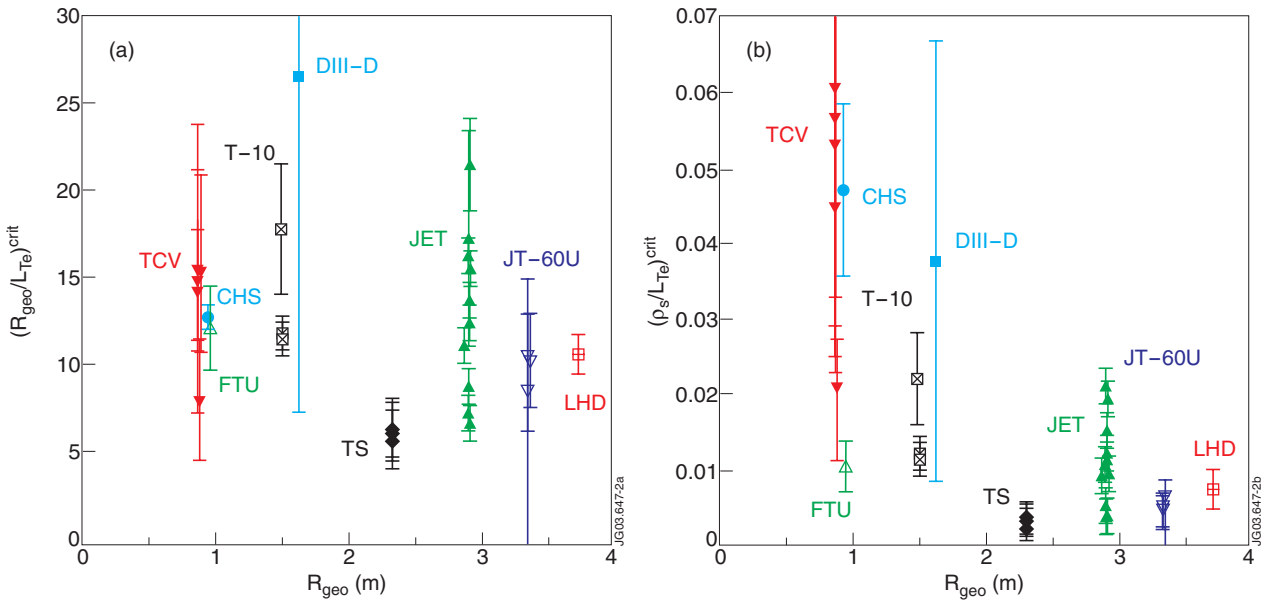


Figure 2: Critical values of (a) R_{geo}/L_{Te} and (b) ρ_s/L_{Te} for T_e ITB.

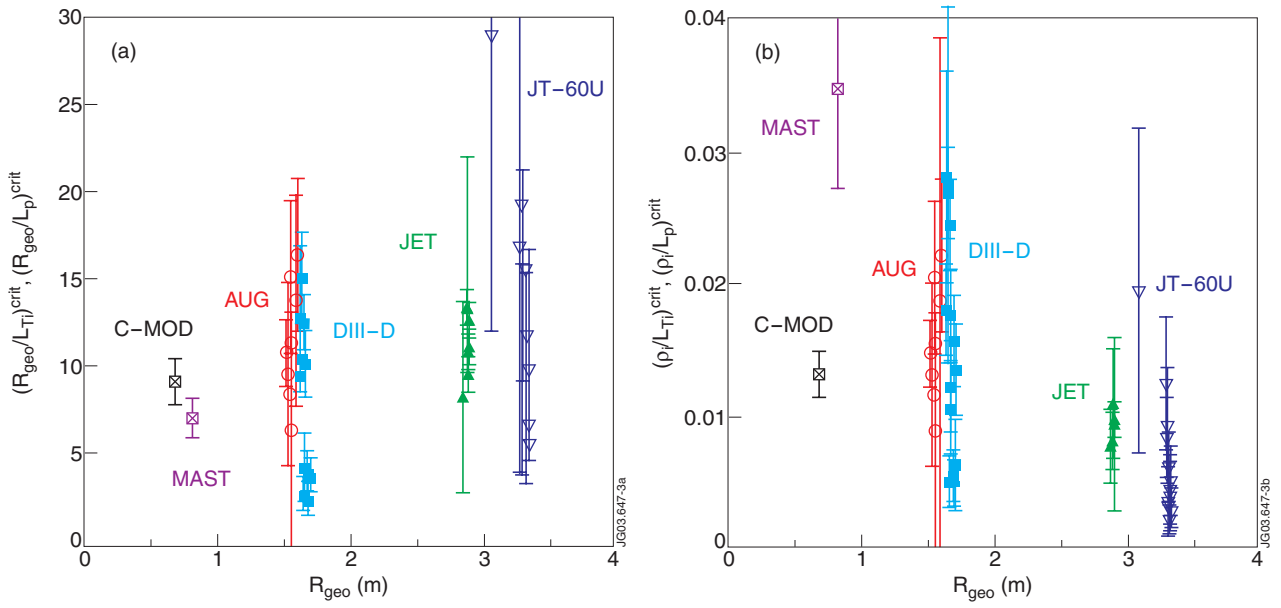


Figure 3: Critical values of (a) R_{geo}/L_{Ti} and (b) ρ_i/L_{Ti} for T_i ITB. For Alcator C-MOD, scale length of pressure, L_p is used instead of L_{Ti} .

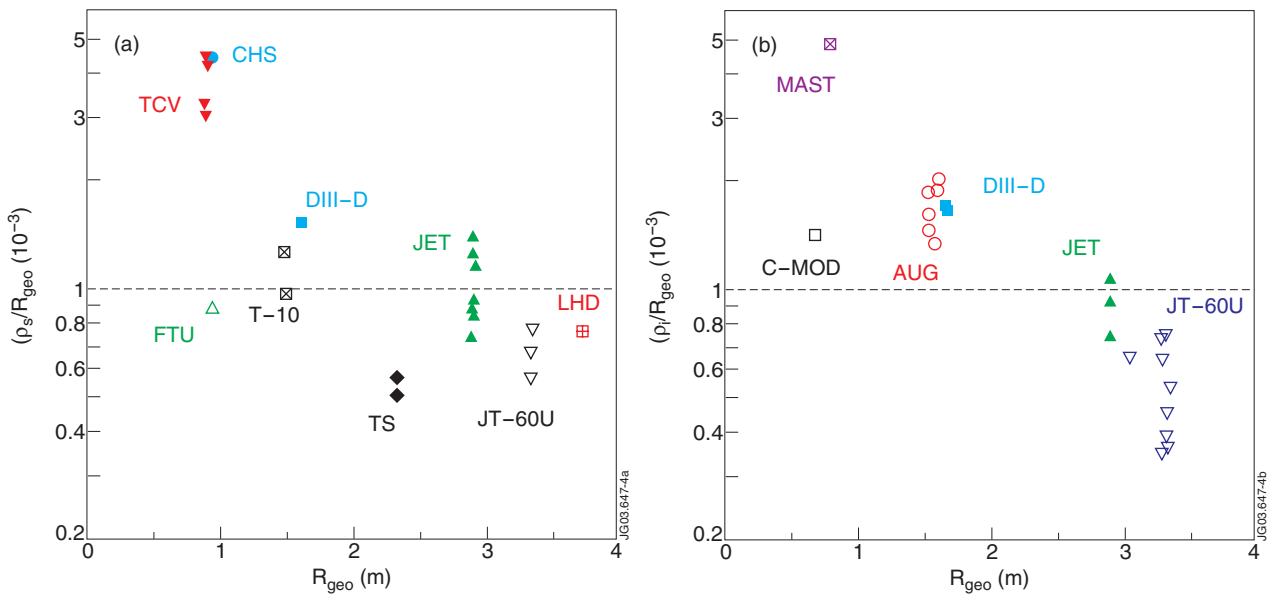


Figure 4: (a) ρ_s/R_{geo} and (b) ρ_i/R_{geo} at the ITB foot as functions of R_{geo} .

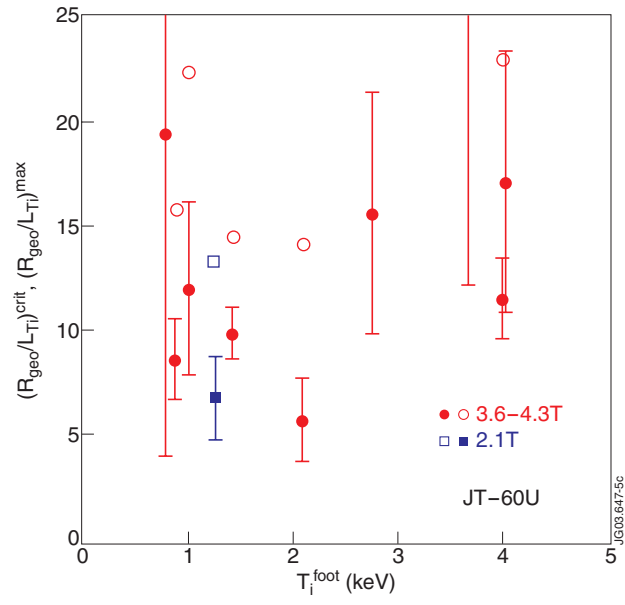


Figure 5: Dependence of R_{geo}/L_{Ti} on the ion temperature at the ITB foot in JT-60U plasmas. Closed symbols denote the critical values while open ones denote the maximum values.