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EFDA–JET–CP(03)01-55

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# Interaction between Core and Edge in JET ELMy H-Modes with Mixed Type I-II ELMs: Consequences for the Extension of this Regime towards High Plasma Current

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\*See Annex of J. Pamela et al., "Overview of Recent JET Results and Future Perspectives",  
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Preprint of Paper to be submitted for publication in Proceedings of the  
EPS Conference on Controlled Fusion and Plasma Physics,  
(St. Petersburg, Russia, 7-11 July 2003)

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## ABSTRACT.

In high triangularity ( $\delta \cong 0.5$ ) JET ELMy H-modes, with safety factor  $q_{95} \cong 3$  and plasma current  $I_p = 2.5\text{MA}$ , high confinement ( $H_{98y} \cong 1$ ) is achieved with line averaged density of  $\cong 0.95$  of the Greenwald density  $n_{GR}$ . In those high density plasmas the Type I ELM frequency  $f_{ELM}$  decreases and Type II ELMs are observed in the inter-ELM period: this regime is called the mixed Type I/II ELMy regime. New experiments have started in JET to study the scaling with plasma current (up to  $I_p = 3.5\text{MA}$ ) of the energy confinement of high  $\delta$  ELMy H-modes and the edge operational space of the mixed Type I/II regime. Since  $\delta \cong 0.5$  is not technically possible in JET at 3.5MA, a plasma configuration with reduced triangularity was selected for those experiments. The extension of high density ELMy H-modes to lower  $v^*$  also requires the higher input power  $P_{IN}$ , that will be soon available with the NBI upgrade, as well as the use of ICRH heating ( $P_{IN} = P_{NBI} + P_{ICRH} \geq 25\text{MW}$ ). This paper presents the results of experiments carried out to prepare for the extension to high  $I_p$  of the high density ELMy H-mode regime: it reports the study of the performance of the new plasma configuration at reduced  $I_p$  (2.5MA), and the results of the first experiment at 3MA.

## 1. DESCRIPTION OF THE EXPERIMENT

The good confinement at high density reported in [1] was obtained with a  $\delta \cong 0.47$  "ITER like" configuration at  $I_p = 2.5\text{MA}$  and toroidal field  $B_t = 2.7\text{T}$ , with  $q_{95} \cong 3$ . The new configuration (called HT3) designed for the experiment described here, has the maximum triangularity that is technically possible at 3.5MA:  $\delta \cong 0.42$ , and it is optimized to reduce the VDE forces. Compared with the ITER like configuration HT3 has larger minor radius ( $a \cong 0.93$ -HT3,  $a \cong 0.89$ -ITER like), similar elongation (HT3,  $\kappa \cong 1.72$ ) and lower X-point position. Three separate experiments in the HT3 configuration are reported in this paper: the main parameters for each are summarized below.

<i>Configuration</i>	$B_T(T)$	$I_P(MA)$	$q_{95}$	<i>Injected power: <math>P_{IN}(MW)</math></i>	<i>Density</i>
<i>HT3</i>	2.7	2.5	3.6	12.5 - 16.5	<i>density scan</i>
<i>HT3</i>	2.25	2.5	3	14	<i>only high density</i>
<i>HT3</i>	2.7	3	3	16.5-17-19	<i>mainly high density</i>

In all cases additional heating is a combination of NBI and 2-3.5MW of ICRH power (H minority central resonance with dipole phasing). The aim of first experiment was to assess if the regime of good confinement found in the ITER like configuration, with mixed Type I/II ELMs, is still achievable with the new HT3 configuration at the same  $I_p$  and  $B_t$ , despite the lower  $\delta$ . In the second experiment,  $q_{95}$  was reduced from 3.6 to 3 at 2.5MA ( $B_t = 2.25\text{T}$ ), to study the performance at high density with the same  $q_{95}$  of the 3MA/2.7T and 3.5MA/3.2T plasma. Finally, results will be presented from a first assessment of the confinement at high density at 3MA/2.7T, with  $q_{95} = 3$ . A significant density scan, achieved by increasing the gas fuelling pulse by pulse, was only carried out in the first

experiment. In the  $q_{95} = 3$  cases mainly data at high density exist. In order to maintain a similar ratio between  $P_{IN}$  and the L-H threshold power,  $P_{L-H}$ , when both  $I_p$  and  $B_t$  are increased at high density ( $n_e \propto n_{GR}$ ),  $P_{IN}$  has to be increased approximately as the product  $I_p * B_t$ , since  $P_{L-H} \propto n_e B_t$  and  $n_{GR} \propto I_p$ . As a consequence, the maximum input power available in JET limits the maximum  $I_p$  ( $B_t$ ) at which high density ELMy H- modes are achievable. Therefore, in order to minimise the power requirements at high  $I_p$ ,  $P_{IN}$  was reduced at high density as part of the experiments at 2.5MA, to evaluate the minimum power at which the good confinement at high density can be obtained. It was found that this minimum  $P_{IN}$  was not limited by the transition to Type III ELMs but by continuous peaking of the density profile, with similar phenomenology as described in [2]. All the discharges analysed in this paper, for global and pedestal data, have no continuous density peaking: core and edge density are steady. A (3,2)NTM mode was often triggered, during the density built up, by the first sawtooth crash. Methods to reduce the extent of the sawtooth crash and avoid the NTM onset were employed. The data presented here are from plasma with no (3,2) NTM. Some of the 3MA/2.7T have a (4,3) MHD mode (see par.2). Given the potentially high disruption forces and the need of scenario development, the heating pulse length of some of the 3MA plasma was shorter than the 6s achieved subsequently. In those plasma, the density required 6/8 energy confinement times ( $\cong 3s$ ) to reach its final value, and the ELM frequency  $f_{ELM}$  even longer. Although the evaluation of the plasma with short pulse length ( $\cong 3.5s$ ) is more uncertain, they are included in the analysis.

## 2. RESULTS

All the characteristic signatures [1,2,3,4] of the mixed Type I/II ELMy regime are observed in the gas scan at 2.5MA/2.7T ( $q_{95} \cong 3.6$ ). The first characteristic is the “anomalous” behavior of the Type I ELM frequency that, after the initial increase with increasing density (Type I ELMs,  $f_{ELM} \sim 10^1$ Hz), decreases in the density range where Type II ELMs are observed in the inter ELM period ( $f_{ELM} \sim 10^0$ Hz), and then increases again ( $f_{ELM} \sim 10^1$ Hz) as the collisionality increases, and mixed Type I/II and III ELMs can be present simultaneously. The transition to a steady Type III ELMs regime is observed at the highest fuelling rate. The second, and most notable characteristic of the presence of mixed Type I/II ELMs is the high density ( $n_e/n_{GR} \cong 1$ ) that is achieved with good confinement ( $H98y \cong 1$ ), as shown in Fig.1 (red triangles).

$H98y \cong 1$  at  $n_e/n_{GR} \cong 1$  is comparable to the previous results with the ITER like configuration [1], at the same  $I_p$  and  $B_t$ . The third observation is that the high global confinement at high density is associated with high pedestal pressure,  $p_{ped}$ , at high density. The evolution of the parameters  $n_e$ ,  $T_e$  (red circles) and  $T_i$  (red triangle) at the top of the H-mode pedestal for the gas scan at  $q_{95} \cong 3.6$  are shown in Fig.2. The pedestal  $T_i$  is measured with Charge Exchange Spectroscopy,  $T_e$  with ECE while  $n_e$  is the edge line averaged density from the FIR interferometer. At low density, in the Type I ELMy regime, the pedestal pressure,  $p_{ped}$ , decreases with density, as typical for JET Type I ELMy H-modes [5]. This corresponds to an initial limited decrease of the total thermal stored energy,  $W_{th}$ . At high density, with mixed Type I/II ELMs,  $p_{ped}$  increases:  $n_e$  increases at almost constant  $T_i$  (and

$W_{th}$  increases). This increase of  $p_{ped}$  was observed previously with both for  $T_e$  [1] and  $T_i$  [6] data. When the fuelling is increased further, both  $n_e$  and  $T_i$  (and  $W_{th}$ ) decrease as the pedestal collisionality increases. The inter ELM Type II behavior is maintained, but with reduced confinement, and Type III ELMs also start to appear.

The good confinement with mixed Type I/II ELMs was achieved in all the power range explored. Plasma with lower  $P_{IN}$  require lower external fuelling to enter this regime and are more difficult to obtain, due to the tendency towards density peaking. The transition from the Type I to the mixed Type I/II ELM regime is also characterized by a reduction of the ELM energy losses ( $f_{ELM} \Delta W$ , where  $\Delta W$  is the prompt energy loss per ELM averaged over, typically, 8-10 ELMs).

In Fig.3, the variation with density of the ELM energy losses normalized to the power to the separatrix ( $f_{ELM} \Delta W/P_{sep}$ ) of the  $q_{95} = 3.6, 2.5MA/2.7T$ , H-modes is compared with the  $q_{95} = 3$  data at 2.5 and 3MA. For the  $q_{95} = 3$  data, it is not possible to distinguish the ELM character by the variation of  $f_{ELM}$  with density, since only high density data are available at 2.5MA and only partial gas scans, at different  $P_{IN}$ , were carried out at 3MA. Fig 3 shows that, at  $q_{95} = 3.6$  (red triangles), a large reduction of the energy loss by ELMs is observed at the transition from the Type I to the Type I/II ELM regimes, in agreement with previous results [1]. The figure also shows that the normalized ELM energy losses at high density are reduced (or low) also in the  $q_{95} = 3$  H-modes, indicating that their pedestal is in the Type I/II ELM regime.

The mixed Type I/II ELM regime is observed above a pedestal  $n_e$  of  $\cong 80\% n_{GR}$  (and pedestal collisionality  $0.4 \leq \nu_{ped} \leq 0.65$ ). Although the scatter of the data is large, there seem to be a correlation between  $f_{ELM} \Delta W/P_{sep}$  of mixed Type I/II ELMs and their  $p_{ped}$  ( $p_{ped}$  taken here as  $\propto 2n_e T_i$ ), with higher  $p_{ped}$  for decreasing  $f_{ELM} \Delta W/P_{sep}$ . Another characteristic signature of the mixed Type I/II ELM regime comes from the Fourier spectra of magnetic fluctuations [1, 5, 6]. The comparison, in Fig.4, of an H-mode with Type I ELMs and one with mixed Type I/II ELMs, at  $q_{95} = 3.6$ , shows an enhancement of the magnetic fluctuations in the low frequency ( $f < 40kHz$ ) range coincident with Type I/II ELMs, and a reduction in the high frequency range, as typically observed [1,3]. In Fig.5, the MHD fluctuation spectra of a  $q_{95} = 3$  Type I ELM plasma, and of two plasma at the same  $q_{95}$  and reduced  $f_{ELM} \Delta W/P_{sep}$ , are compared. Although this comparison at  $q_{95} = 3$  does show the characteristic enhancement at lower frequencies, this is less pronounced than at higher  $q_{95}$ . Both MHD and ELM loss measurements demonstrate that the mixed Type I/II ELM regime is achieved at  $q_{95} = 3$ , but weaker Type II signature is seen in the MHD data. A weaker Type I/II MHD signature and decreased inter-ELM losses are also seen, at  $q_{95} = 3.6$ , in Type I/II ELM plasmas with increased gas fuelling and collisionality (with reduced  $p_{ped}$ ), indicating that an optimum density or collisionality range for good performance exist.

The global confinement and pedestal  $n_e$ - $T_i$  of the  $q_{95} = 3.6$  data with clear mixed Type I/II behavior are compared in Fig.1 and 2 with the results at  $q_{95} = 3$ : the 2.5MA/2.25T (black diamonds) and 3MA/2.7T (green circles). The comparison of the 2.5MA data at  $q_{95} = 3.6$  and  $q_{95} = 3$  show that with reduced  $q_{95}$  it was possible to achieve similar density but with lower energy confinement: H98y

decreased from  $\cong 1$  to  $\cong 0.9$ . The pedestal  $T_i$  of the 2.5MA at  $q_{95} = 3$  is also lower than at  $q_{95} = 3.6$ . A similar reduction in global confinement at high density is observed when comparing the  $q_{95} = 3.6$  with the 3MA data. The higher pedestal pressure at higher  $I_p$  results from higher pedestal  $n_e$ , with lower or similar  $T_i$ . While the difference in H factor between the  $q_{95} = 3.6$   $q_{95} = 3$  H-modes is not large ( $\cong 10\%$ ), the consistency between pedestal and global confinement data indicate that this is a real trend. For example, Pulse No's: 56146 (H98y  $\cong 1$ ,  $q_{95} = 3.6$ ) and 58012 (H98y  $\cong 0.9$ ,  $q_{95} = 3$ ) with the same  $I_p$  and similar input power and external fuelling, reach similar core ( $n_e/n_{eGR} = 1$ ) and pedestal density. The main difference between the two shot is in the pedestal (and core)  $T_i$ , which is lower at  $q_{95} = 3$ . Consistently, the contribution of the pedestal stored energy,  $W_{ped}$ , ( $W_{ped} \propto p_{ped}$ ) to the total stored energy is  $\cong 50\%$  for both discharges.

The ratio between pedestal and thermal stored energy is 0.4 to 0.5 for all the data, independent of the ELM type (for Type I and Type I/II ELMs) and of  $q_{95}$ . This suggests that the improved global confinement at high density is associated with the improved pedestal resulting from the change in the ELM behavior. In addition, Fig.4 shows that  $W_{ped}$  of the  $q_{95} = 3$  plasmas at 2.5 and 3MA scales as  $\propto I_p (T_i)^{0.5}$  [6]. The scaling of the pedestal pressure with  $I_p$  (or stronger) results in similar global H factor. The pedestal data for a gas scan with the HT3 configuration at  $q_{95} = 4.6$  (2.5MA/3.4T) are also shown in Fig.2 to demonstrate that the differences between the  $q_{95} = 3.6$  and  $q_{95} = 3$  cannot simply be attributed to a trend of decreasing local and global confinement with decreasing  $q_{95}$ , since both pedestal and global [4] confinement are lower also at  $q_{95} = 4.6$ . At  $q_{95} = 4.6$ , the transition to Type III ELMs occurs at lower density. Those data are described in more detail in [4]. All those observations indicate that the main difference between the  $q_{95} = 3.6$  and  $q_{95} = 3$  data derive from differences in the pedestal. The lower  $p_{ped}$  at lower  $q_{95}$  might be attributed to the weaker Type I/II ELM signature and weaker inter-ELM transport. It has to be noted however, that additional elements can, when combined, contribute to some further reduction in the H factor of the 3MA H-modes: the positive density scaling of the energy confinement (generally not observed in the data), some reduction of the central beam deposition at higher density, and sometimes the presence of a (4, 3) MHD mode (for comparison, a continuous (3, 2) mode reduces the confinement by 5-8%).

## CONCLUSIONS

The good confinement at high density, previously reported with the ITER like configuration, has been obtained also in the configuration designed for high  $I_p$  (i.e., up to 3.5MA). The improved pedestal confinement at high density is associated with the transition from the Type I to the mixed Type I/II regime. The achievement of the Type I/II ELM regime at the reduced  $\delta$  of this configuration is probably facilitated by increased  $q_{95}$ . At  $q_{95} = 3.6$ , 2.5MA/2.7T, H98y  $\delta 1$  at  $n_e/n_{eGR} = 1$  were obtained. The reduced pedestal and core confinement (H98y  $\cong 0.9$  at  $n_e/n_{eGR} = 1$ ) at reduced  $q_{95}$  ( $q_{95} = 3$ ,  $I_p = 2.5$  and 3MA) could be linked to the weaker Type II ELM character of the H-mode pedestal and reduced inter-ELM transport, since no trend of reduced confinement with  $q_{95}$  is seen for the Type I ELM plasmas. Further experiments both at 3MA and 2.5MA, in particular more extensive density variation



and the exploitation of higher  $P_{IN}$ , are planned. Those experiments should clarify if the reduced Type I/II signature is due to the reduced  $q_{95}$  of those plasmas, or if the optimum conditions for Type I/II have not been achieved yet (or both). The experiments at higher  $I_p$  (3.5MA/3.2T) and at higher  $q_{95}$  (2.5MA/3.45T and 3MA/3.45T), to confirm and clarify the trend of the data with  $I_p$  and  $q_{95}$ , will require, for operation at high density, the higher PIN provided by the NB upgrade.

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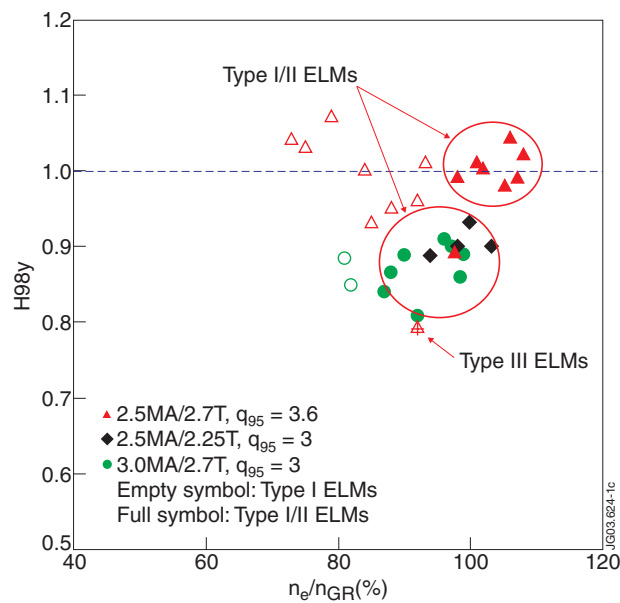


Figure 1: Confinement enhancement factor versus normalized density for ELMy H-modes with the HT3 plasma magnetic configuration.

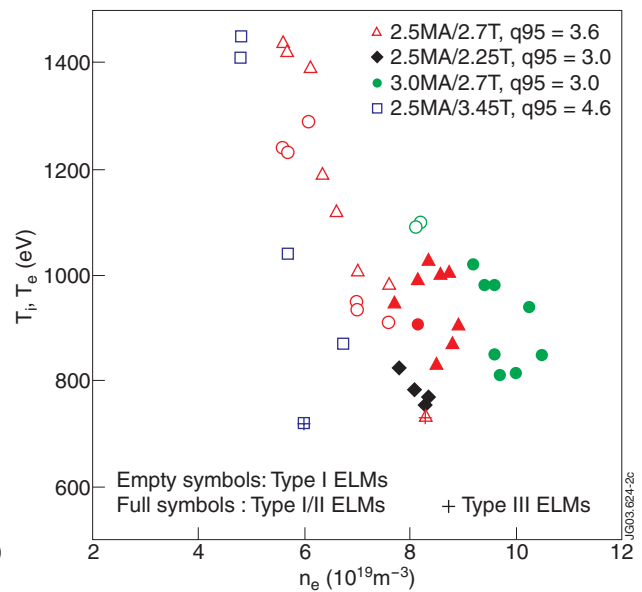


Figure 2: Pedestal temperature ( $T_i$ ,  $T_e$ ) versus pedestal density  $n_e$ .  $n_e$  from the outermost channel of the FIR interferometer (at  $R = 3.75m$ ),  $T_i$  from Charge Exchange Spectroscopy (at a fixed position,  $R = 3.76-3.78m$ ) and  $T_e$  from ECE (at the pedestal top). Due to ECE cut-off, the  $T_e$  data are only available at low density at 2.5MA/2.7T: while at the lowest density  $T_e > T_i$ , at intermediate density  $T_e \cong T_i$ . The  $T_e$  value at the highest density (Type I/II ELMs) is underestimated: it is taken before the end of the ELM cycle (just before the ECE cuts-off).

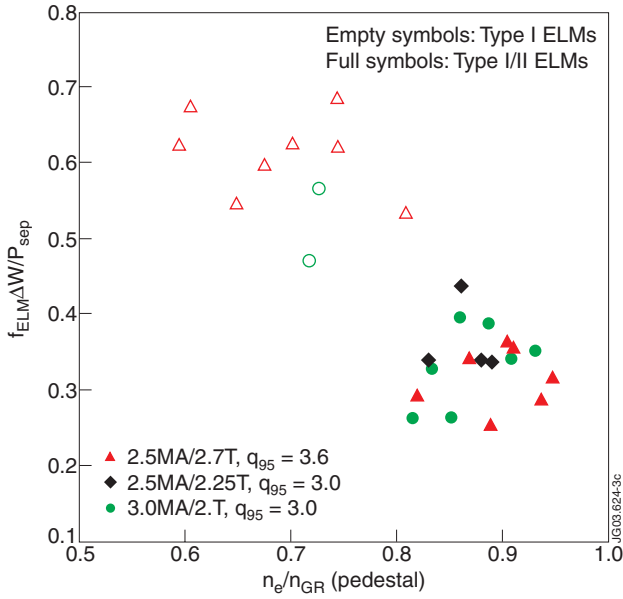


Figure 3: ELM ene gy losses normalized to the power to the separatrix versus normalized pedestal density for the gas scans at  $q_{95}=3.6$  (2.5MA) and  $q_{95}=3$  (2.5MA and 3MA)

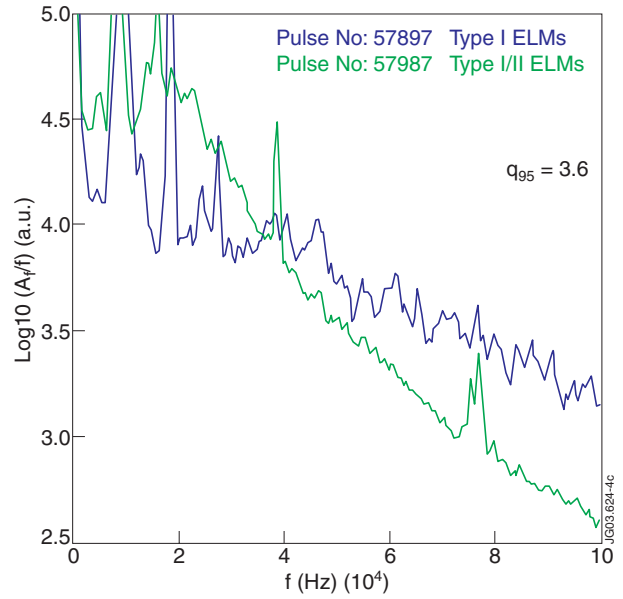


Figure 4: Spectra of the magnetic fluctuation in the inter-ELM period for Pulse No: 57987 (green) and Pulse No: 57897 (blue). An enhancement of the fluctuations in the low frequency range is observed with Type I/II ELMs.

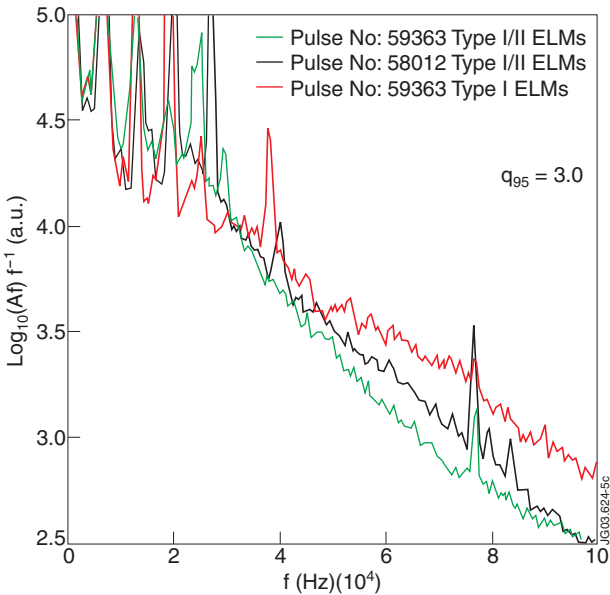


Figure 5: Spectra of the magnetic fluctuation in the inter-ELM period for Pulse No's: 59368 (red, 3MA/2.7T, Type I ELMs), Pulse No: 58012 (2.5MA/2.25T, Type I/II ELMs) and Pulse No: 59363 (3MA/2.7T, Type I/II ELMs)

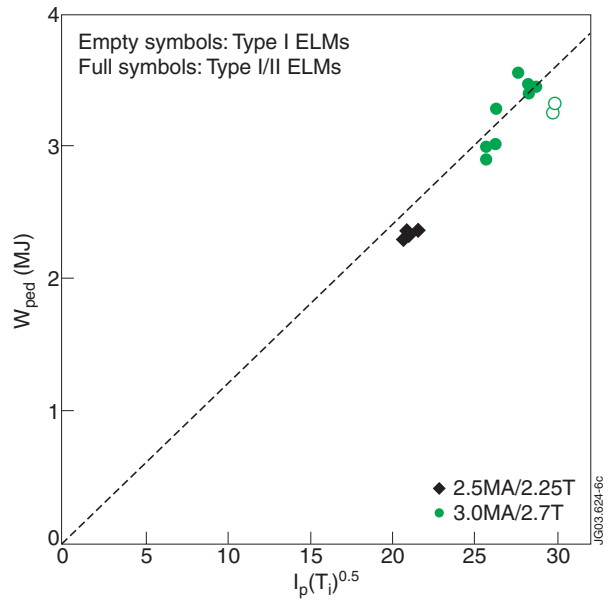


Figure 6: Pedestal stored energy ( $\propto n_e T_i$ ) versus  $I_p(T_i)^{0.5}$ , for the  $q_{95} = 3$  ELMy H-modes at 2.5 and 3MA.